Poincaré series and linking of Legendrian knots Joint work with Nguyen Viet Dang (Université Lyon 1)

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November, 12th 2020

#### Introduction

In all this talk, (X, G) will be a smooth  $(\mathcal{C}^{\infty})$ , compact, connected, oriented and Riemannian **surface** which has no boundary and which has **negative curvature** (a priori non constant).



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$$\mathcal{N}_{T}(c_{1},c_{2}):=|\{\gamma\in\mathcal{P}_{c_{1},c_{2}}:\ 0<\ell(\gamma)\leq T\}|\sim \mathcal{A}_{c_{1},c_{2}}e^{Th_{top}},$$

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In constant curvature  $K \equiv -1$ ,  $h_{top} = 1$ . Otherwise take this Theorem as a definition of  $h_{top}$ !

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 In constant curvature, related to the spectral decomposition of the Laplacian (Delsarte, Huber).  In constant curvature, related to the spectral decomposition of the Laplacian (Delsarte, Huber).

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► In *variable curvature*, related to the **mixing properties** of the Bowen-Margulis measure (Margulis).

► For more informations on the behaviour of N<sub>T</sub>(c<sub>1</sub>, c<sub>2</sub>) and its applications, see the survey of Parkkonen and Paulin (LMS Lecture notes **425**, 2016).

Poincaré series as zeta renormalization of  $\mathcal{N}_T(c_1, c_2)$ 

Let  $s \in \mathbb{C}$  and set

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Thanks to Margulis Theorem, the Poincaré series

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defines a holomorphic function in the half plane

$$\{w \in \mathbb{C} : \operatorname{\mathsf{Re}}(w) > h_{\operatorname{top}}\}.$$

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#### Comments and related results

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- ▶ In the case of **constant curvature**  $K \equiv -1$ , Paternain proved (2000) :

$$\lim_{T \to +\infty} \int_{X \times X} \mathcal{N}_T(c_1, c_2, s) d\operatorname{vol}_G(c_1) d\operatorname{vol}_G(c_2) = \frac{4\pi^2 \chi(X)}{1 - s^2}.$$

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 Similar results for Ruelle zeta functions which "count periodic orbits" (Selberg, Smale, Ruelle, Rugh, Fried, Kitaev, Baladi-Tsujii, Giulietti-Liverani-Pollicott, Dyatlov-Zworski, Faure-Tsujii, Dyatlov-Guillarmou, Jezequel, etc.)

## Meromorphic continuation of Poincaré series

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For  $\gamma \in \mathcal{P}_{c_1,c_2}$ , one has

 $\varphi^{\ell(\gamma)}(S^*_{c_1}X)\cap S^*_{c_2}X\neq\emptyset.$ 

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The current

$$\mathsf{R}_{\mathsf{T}} := -\int_0^{\mathsf{T}} \iota_V \varphi^{-t*}[S_c^*X] dt$$

represents the integration on the surface

$$\left\{ arphi^t(x) : x \in S^*_c X \text{ and } 0 \leq t \leq T 
ight\} \subset S^* X.$$



Let  $c_1$  and  $c_2$  be two points in X and let T > 0. Then, one has

$$\mathcal{N}_{T}(c_{1},c_{2},s)=-\int_{S^{*}X}[S^{*}_{c_{2}}X]\wedge\int_{0}^{T}e^{-st}\iota_{V}\varphi^{-t*}[S^{*}_{c_{1}}X]dt.$$

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Formally, we can then let  $\, \mathcal{T} 
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Hence, formally, we need to understand the meromorphic continuation of

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Part 1 of our Theorem follows if we have the meromorphic continuation of the  $\ensuremath{\textbf{spectral resolvent}}$ 

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 $\rightsquigarrow$  Related to the sudy of transfer operators in dynamical systems.

Proving the meromorphic continuation of  $(\mathcal{L}_V + s)^{-1}$  requires to introduce an appropriate functional framework made of distributions (or currents) with **anisotropic Hölder/Sobolev regularity** :
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- Using Fourier and/or microlocal analysis. Baladi, Baladi-Tsujii, Faure-Roy-Sjöstrand, Faure-Sjöstrand, Tsujii, Faure-Tsujii, Dyatlov-Zworski, Dyatlov-Guillarmou, Bonthonneau-Weich, etc.

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 $\rightsquigarrow$  Non exhaustive list...

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A key ingredient is the **precise description of the spectral projector**  $\pi_0$  on the eigenvalue 0 made by Dyatlov-Zworski (2017) :

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Moreover, they describe  $Ran(\pi_0)$  in terms of generators of the De Rham cohomology.

Using this result and the fact that  $[S_c^*X]$  is exact <sup>2</sup>, one finds

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Thus, one has

$$\mathcal{N}_{\infty}(c_1, c_2, 0) = -\int_{S^*X} [S^*_{c_2}X] \wedge \iota_V \mathcal{L}_V^{-1}[S^*_{c_1}X].$$

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Write now

$$[S_c^*X] = (d\iota_V + \iota_V d)\mathcal{L}_V^{-1}[S_c^*X] = d\left(\iota_V \mathcal{L}_V^{-1}[S_c^*X]\right) = dR_c.$$

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In summary, the value at 0 is given by :

$$\mathcal{N}_{\infty}(c_1, c_2, 0) = -\int_{S^*X} [S^*_{c_2}X] \wedge R_{c_1} =: -\mathbf{L}(c_1, c_2).$$

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### The value at 0 as the linking of two knots

One has

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We say that  $S_c^*X$  is a **Legendrian knot** in  $S^*X$ .

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$$L(c_1, c_2) := \int_{S^*X} [S^*_{c_2}X] \wedge R_{c_1} = -\frac{1}{\chi(X)}.$$

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Implicitely, this shows that  $[S_a^*X]$  is exact.

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$$= -\sum_{a \in Crit(f)} (-1)^{ind(a)} L(c_1, a)$$
  

$$= -\chi(X) L(c_1, c_2).$$

#### A more general picture

Let  $\mathbf{c}_1$  and  $\mathbf{c}_2$  in  $\pi_1(X)$ . If  $\mathbf{c}_i$  is nontrivial, one can find an unique geodesic  $c_i$  in the conjugacy class of  $\mathbf{c}_i \in \pi_1(X)$ . We say that it is a **geodesic representative** of  $\mathbf{c}_i$ .





Let  $c_1$  and  $c_2$  be two elements in  $\pi_1(X)$  and let  $c_1$  and  $c_2$  be two of their geodesic representatives in X. Set

 $\mathcal{P}_{c_1,c_2} := \{ \text{geodesic arcs joining } c_1 \text{ and } c_2 \text{ and directly } \perp \text{ to } c_1 \text{ and } c_2 \}.$ 

Let  $c_1$  and  $c_2$  be two elements in  $\pi_1(X)$  and let  $c_1$  and  $c_2$  be two of their geodesic representatives in X. Then,

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$$s\mapsto \mathcal{N}_\infty(c_1,c_2,s):=\sum_{\gamma\in\mathcal{P}_{\mathbf{c_1},\mathbf{c_2}}:\;\ell(\gamma)>0}e^{-s\ell(\gamma)}$$

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▶ in that case, N<sub>∞</sub>(c<sub>1</sub>, c<sub>2</sub>, 0) is the linking number of the unit (direct) conormal bundles of c<sub>1</sub> and c<sub>2</sub>.

#### Comments.

When c₁ and c₂ are both nontrivial in π₁(X), the linking number we obtain is also the linking number of the closed orbits lifting c₁ and c₂ in S\*X.
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There is an explicit expression for the linking number of c<sub>1</sub> and c<sub>2</sub> in terms of Euler characteristics :

$$\pm \mathcal{N}_{\infty}(c_1,c_2,0) = rac{\chi(X(c_1))\chi(X(c_2))}{\chi(X)} - \chi(X(c_1)\cap X(c_2)) + rac{1}{2}\chi(c_1\cap c_2).$$

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- ► In the arithmetic framework, similar phenomena occur for Dirichlet series (Duke et al. 2017) and *L*-functions (Bergeron et al. 2018). Relation with these results?
- What about higher dimensions? Probably a pole when the dimension is odd.
- What can be extracted using the spectral decomposition of the Laplacian ?

Thank you for your attention.