# Poincaré series and linking of Legendrian knots Joint work with Nguyen Viet Dang (Université Lyon 1) 

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## Introduction

In all this talk, $(X, G)$ will be a smooth $\left(\mathcal{C}^{\infty}\right)$, compact, connected, oriented and Riemannian surface which has no boundary and which has negative curvature (a priori non constant).


Let $c_{1}$ and $c_{2}$ be two points in $X$ and set ${ }^{1}$

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\mathcal{P}_{c_{1}, c_{2}}:=\left\{\gamma: \text { geodesic arcs joining } c_{1} \text { and } c_{2}\right\} .
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1. Geodesic arcs are parametrized by arc-length.

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Theorem (Delsarte (1942), Huber (1959), Margulis (1969)) There exists $A_{c_{1}, c_{2}}>0$ such that, as $T \rightarrow+\infty$,

$$
\mathcal{N}_{T}\left(c_{1}, c_{2}\right):=\left|\left\{\gamma \in \mathcal{P}_{c_{1}, c_{2}}: 0<\ell(\gamma) \leq T\right\}\right| \sim A_{c_{1}, c_{2}} e^{T h_{\text {top }}}
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where $\ell(\gamma)$ is the length of $\gamma$ and $h_{\text {top }}>0$ is the topological entropy of the geodesic flow.

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where $\ell(\gamma)$ is the length of $\gamma$ and $h_{\text {top }}>0$ is the topological entropy of the geodesic flow.

In constant curvature $K \equiv-1, h_{\text {top }}=1$. Otherwise take this Theorem as a definition of $h_{\text {top }}$ !

- In constant curvature, related to the spectral decomposition of the Laplacian (Delsarte, Huber).
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- For more informations on the behaviour of $\mathcal{N}_{T}\left(c_{1}, c_{2}\right)$ and its applications, see the survey of Parkkonen and Paulin (LMS Lecture notes 425, 2016).


## Poincaré series as zeta renormalization of $\mathcal{N}_{T}\left(c_{1}, c_{2}\right)$

Let $s \in \mathbb{C}$ and set

$$
\mathcal{N}_{T}\left(c_{1}, c_{2}, s\right):=\sum_{\gamma \in \mathcal{P}_{c_{1}, c_{2}}: 0<\ell(\gamma) \leq T} e^{-s \ell(\gamma)} .
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Thanks to Margulis Theorem, the Poincaré series

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\mathcal{N}_{\infty}\left(c_{1}, c_{2}, s\right):=\lim _{T \rightarrow+\infty} \mathcal{N}_{T}\left(c_{1}, c_{2}, s\right)=\sum_{\gamma \in \mathcal{P}_{c_{1}}, c_{2}: \ell(\gamma)>0} e^{-s \ell(\gamma)}
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defines a holomorphic function in the half plane

$$
\left\{w \in \mathbb{C}: \operatorname{Re}(w)>h_{\text {top }}\right\}
$$

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Let $c_{1}$ and $c_{2}$ be two points in $X$. Denote by $\chi(X)=2-2 g(X)$ the Euler characteristic of $X$. Then,

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- if $c_{1}=c_{2}$, one has

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\lim _{T \rightarrow+\infty} \int_{X \times X} \mathcal{N}_{T}\left(c_{1}, c_{2}, s\right) d \operatorname{vol}_{G}\left(c_{1}\right) d \operatorname{vol}_{G}\left(c_{2}\right)=\frac{4 \pi^{2} \chi(X)}{1-s^{2}}
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- Similar results for Ruelle zeta functions which "count periodic orbits" (Selberg, Smale, Ruelle, Rugh, Fried, Kitaev, Baladi-Tsujii, Giulietti-Liverani-Pollicott, Dyatlov-Zworski, Faure-Tsujii, Dyatlov-Guillarmou, Jezequel, etc.)


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For $\gamma \in \mathcal{P}_{c_{1}, c_{2}}$, one has

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\varphi^{\ell(\gamma)}\left(S_{c_{1}}^{*} X\right) \cap S_{c_{2}}^{*} X \neq \emptyset
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For $\psi \in \Omega^{1}\left(S^{*} X\right)$, we set

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& =\left[S_{c}^{*} X\right]-d\left(\int_{0}^{T}{ }^{\iota} \varphi^{-t *}\left[S_{c}^{*} X\right] d t\right)
\end{aligned}
$$

The current

$$
R_{T}:=-\int_{0}^{T} \iota v \varphi^{-t *}\left[S_{c}^{*} X\right] d t
$$

represents the integration on the surface

$$
\left\{\varphi^{t}(x): x \in S_{c}^{*} X \text { and } 0 \leq t \leq T\right\} \subset S^{*} X .
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## Proposition

Let $c_{1}$ and $c_{2}$ be two points in $X$ and let $T>0$. Then, one has

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\mathcal{N}_{T}\left(c_{1}, c_{2}, s\right)=-\int_{S^{*} X}\left[S_{c_{2}}^{*} X\right] \wedge \int_{0}^{T} e^{-s t} \iota_{V} \varphi^{-t *}\left[S_{c_{1}}^{*} X\right] d t
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\mathcal{N}_{\infty}\left(c_{1}, c_{2}, s\right) & =-\int_{S^{*} X}\left[S_{c_{2}}^{*} X\right] \wedge \iota \nu\left(\mathcal{L}_{V}+s\right)^{-1}\left[S_{c_{1}}^{*} X\right] .
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Part 1 of our Theorem follows if we have the meromorphic continuation of the spectral resolvent

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$\rightsquigarrow$ Related to the sudy of transfer operators in dynamical systems.

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- Using Fourier and/or microlocal analysis. Baladi, Baladi-Tsujii, Faure-Roy-Sjöstrand, Faure-Sjöstrand, Tsujii, Faure-Tsujii, Dyatlov-Zworski, Dyatlov-Guillarmou, Bonthonneau-Weich, etc.

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$\rightsquigarrow$ Non exhaustive list...


## Behaviour at $s=0$

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A key ingredient is the precise description of the spectral projector $\pi_{0}$ on the eigenvalue 0 made by Dyatlov-Zworski (2017) :

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Moreover, they describe $\operatorname{Ran}\left(\pi_{0}\right)$ in terms of generators of the De Rham cohomology.

Using this result and the fact that $\left[S_{c}^{*} X\right]$ is exact ${ }^{2}$, one finds

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\pi_{0}\left(\left[S_{c}^{*} X\right]\right)=0
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Write now

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In summary, the value at 0 is given by :

$$
\mathcal{N}_{\infty}\left(c_{1}, c_{2}, 0\right)=-\int_{S^{*} X}\left[S_{c_{2}}^{*} X\right] \wedge R_{c_{1}}=:-\mathbf{L}\left(c_{1}, c_{2}\right) .
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## The value at 0 as the linking of two knots

One has

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\left[S_{c}^{*} X\right]=d R_{c} \quad \text { and } \quad \alpha \wedge\left[S_{c}^{*} X\right]=0
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where $\alpha$ is the contact 1-form on $S^{*} X$.

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We say that $S_{c}^{*} X$ is a Legendrian knot in $S^{*} X$.

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Implicitely, this shows that $\left[S_{a}^{*} X\right]$ is exact.

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& =-\sum_{a \in \operatorname{Crit}(f)}(-1)^{\operatorname{ind}(a)} \mathbf{L}\left(c_{1}, a\right) \\
& =-\chi(X) \mathbf{L}\left(c_{1}, c_{2}\right) .
\end{aligned}
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## A more general picture

Let $\mathbf{c}_{1}$ and $\mathbf{c}_{2}$ in $\pi_{1}(X)$. If $\mathbf{c}_{i}$ is nontrivial, one can find an unique geodesic $c_{i}$ in the conjugacy class of $\mathbf{c}_{i} \in \pi_{1}(X)$. We say that it is a geodesic representative of $\mathbf{c}_{i}$.



Let $\mathbf{c}_{1}$ and $\mathbf{c}_{2}$ be two elements in $\pi_{1}(X)$ and let $c_{1}$ and $c_{2}$ be two of their geodesic representatives in $X$. Set
$\mathcal{P}_{c_{1}, c_{2}}:=\left\{\right.$ geodesic arcs joining $c_{1}$ and $c_{2}$ and directly $\perp$ to $c_{1}$ and $\left.c_{2}\right\}$.

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- the map

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s \mapsto \mathcal{N}_{\infty}\left(c_{1}, c_{2}, s\right):=\sum_{\gamma \in \mathcal{P}_{c_{1}}, c_{2}: \ell(\gamma)>0} e^{-s \ell(\gamma)}
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- in that case, $\mathcal{N}_{\infty}\left(c_{1}, c_{2}, 0\right)$ is the linking number of the unit (direct) conormal bundles of $c_{1}$ and $c_{2}$.


## Comments.

- When $c_{1}$ and $c_{2}$ are both nontrivial in $\pi_{1}(X)$, the linking number we obtain is also the linking number of the closed orbits lifting $c_{1}$ and $c_{2}$ in $S^{*} X$.


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$\rightsquigarrow$ Related results on the modular surface (Ghys, Duke-Immamoglu-Toth) and for suspension of toral automorphisms (Bergeron-Charollois-Garcia-Venkatesh).
- There is an explicit expression for the linking number of $c_{1}$ and $c_{2}$ in terms of Euler characteristics:

$$
\pm \mathcal{N}_{\infty}\left(c_{1}, c_{2}, 0\right)=\frac{\chi\left(X\left(c_{1}\right)\right) \chi\left(X\left(c_{2}\right)\right)}{\chi(X)}-\chi\left(X\left(c_{1}\right) \cap X\left(c_{2}\right)\right)+\frac{1}{2} \chi\left(c_{1} \cap c_{2}\right) .
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- What about higher dimensions ? Probably a pole when the dimension is odd.
- What can be extracted using the spectral decomposition of the Laplacian?

Thank you for your attention.

