

Coupling of two Discretization Schemes for the Lattice Boltzmann Equation

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Content

- 1 Context
- 2 Theory
- 3 Coupling
- 4 Conclusion

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PhD thesis is financed within the **CLIMB** project.

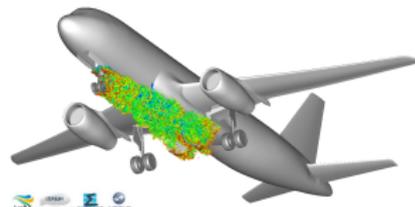
Computational methods with **Intensive Multiphysics Boltzmann** solver

Partners

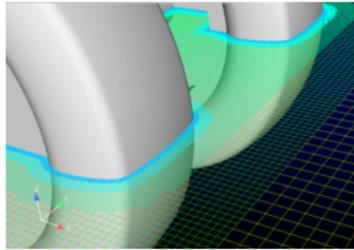
Industrial	Academic
Airbus	AMU
CS	ECL
Renault	UPS
...	...



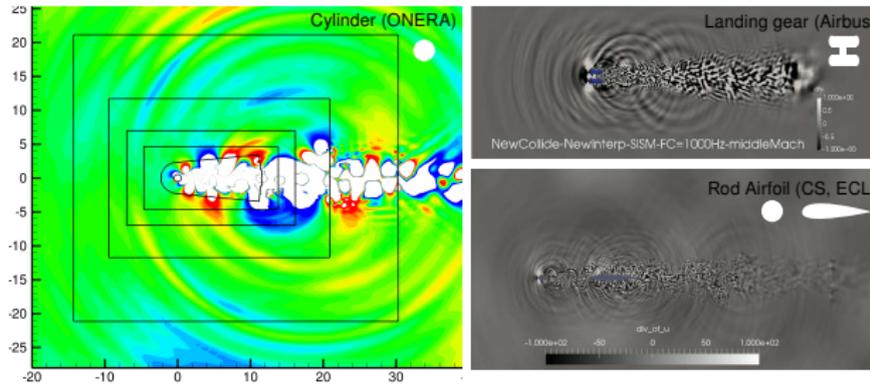
- **ONERA** is involved with a number of departments
- Thesis: Mesh refinement



Numerical noise generation at the interface



Mesh refinement in classical LB algorithm by **factor 2** (or multiple of 2)
 → aeroacoustic solutions may be contaminated



Solution approach: Coupling of classical algorithm with finite volume formulation

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Discretization in space and time of the DVBE

$$\frac{\partial f_\alpha}{\partial t} + \xi_\alpha \frac{\partial f_\alpha}{\partial X} = \Omega(f_\alpha) = \frac{1}{\tau} (f_\alpha^{eq} - f_\alpha), \quad 0 \leq \alpha \leq q-1 \quad (1)$$

- 1 Method of characteristics (classical stream-collide)
- 2 Finite volume method

$$\frac{df_\alpha}{ds} = \frac{\partial f_\alpha}{\partial t} \frac{\partial t}{\partial s} + \frac{\partial f_\alpha}{\partial X} \frac{\partial X}{\partial s} = -\frac{1}{\tau} [f_\alpha - f_\alpha^{eq}] \quad \text{with} \quad f_\alpha = f_\alpha(x(s), t(s)) \quad (2)$$

True for $\frac{\partial t}{\partial s} = 1$ and $\frac{\partial X}{\partial s} = \xi_\alpha$. Thus $t(s) = t(0) + s$ and $x(s) = x(0) + \xi_\alpha s$.
Integrating from $s = 0$ to $s = \Delta t$, we obtain

$$f_\alpha(x + \xi_\alpha \Delta t, t + \Delta t) - f_\alpha(x, t) = -\frac{1}{\tau} \int_0^{\Delta t} [f_\alpha(x + \xi_\alpha s, t + s) - f_\alpha^{eq}(x + \xi_\alpha s, t + s)] ds \quad (3)$$

with $\xi_\alpha = c \cdot e_\alpha$, c is the lattice speed $\Delta x / \Delta t$, usually set to 1.

Discretization in space and time of the DVBE

Solving the integral on the right hand side with the trapezoidal rule yields

$$f_{\alpha}(x + \xi_{\alpha} \Delta t, t + \Delta t) - f_{\alpha}(x, t) = -\frac{\Delta t}{2\tau} [f_{\alpha}(x + \xi_{\alpha} \Delta t, t + \Delta t) - f_{\alpha}^{eq}(x + \xi_{\alpha} \Delta t, t + \Delta t) + f_{\alpha}(x, t) - f_{\alpha}^{eq}(x, t)] + \mathcal{O}(\Delta t)^3 \quad (4)$$

Making equation (4) explicit and applying following change of variable

$$g_{\alpha} = f_{\alpha} + \frac{\Delta t}{2\tau} (f_{\alpha} - f_{\alpha}^{eq}), \quad (5)$$

we obtain

$$g_{\alpha}(x + \xi_{\alpha} \Delta t, t + \Delta t) = g_{\alpha}(x, t) - \frac{\Delta t}{\tau_g} (g_{\alpha}(x, t) - g_{\alpha}^{eq}(x, t)) \quad (6)$$

which is normally solved in a two-step *stream-collide* algorithm:

Collision:

$$\hat{g}_{\alpha}(x, t) = g_{\alpha}(x, t) - \frac{\Delta t}{\tau_g} (g_{\alpha}(x, t) - g_{\alpha}^{eq}(x, t))$$

Stream:

$$g_{\alpha}(x + \xi_{\alpha} \Delta t, t + \Delta t) = \hat{g}_{\alpha}(x, t)$$

Discretization in space and time of the DVBE

- 1 Method of characteristics (classical stream-collide)
- 2 Finite volume method

In order to derive the finite volume formulation of the BE we depart from the DVLBE(1) and integrate over the volume dV .

$$\int_V \frac{\partial f_\alpha}{\partial t} dV + \int_V \xi_\alpha \frac{\partial f_\alpha}{\partial x} dV = \int_V -\frac{1}{\tau} [f_\alpha - f_\alpha^{eq}] dV \quad (7)$$

Shrestha et al. [1] showed that it is possible to apply the same semi-implicit treatment of the collision operator to the finite volume formulation:

$$g_\alpha(x, t + \Delta t) = \underbrace{g_\alpha(x, t) - \frac{\Delta t}{\tau_g} (g_\alpha(x, t) - g_\alpha^{eq}(x, t))}_{\hat{g}_\alpha(x, t) = \text{post-collision}} - \boxed{\Delta t \frac{S_\beta}{V} \xi_\alpha} \cdot n_\beta [f_\alpha(x_\beta, t + \Delta t/2)] \quad (8)$$

For the evaluation of the surface fluxes we chose a 2^{nd} order scheme in **time** (Heun predictor-corrector, DUGKS) and 2^{nd} or higher order scheme in **space** (centred, QUICK)

Discretization in space and time of the DVBE

Evaluation of $F^{t+\frac{1}{2}} = \xi_\alpha \cdot n_\beta [f_\alpha(x_\beta, t + \Delta t/2)]$

- 1 Heun predictor-corrector scheme
- 2 DUGKS

$$F^{t+\frac{1}{2}} = \xi_\alpha \cdot n_\beta \left[\frac{f_\alpha^*(x_\beta, t + \Delta t)}{2} + \frac{f_\alpha(x_\beta, t)}{2} \right] \quad (9)$$

with $f_\alpha = g_\alpha + \frac{1}{2\bar{\tau}_g} (f_\alpha^{eq} - g_\alpha)$ and $f_\alpha = \hat{g}_\alpha + \frac{1}{2(\bar{\tau}_g - 1)} (\hat{g}_\alpha - f_\alpha^{eq})$ respectively

Three step (advect-collide) algorithm:

Step 1: Collision $\rightarrow \hat{g}_\alpha(x, t)$

Step 2: Prediction $\rightarrow g_\alpha^*(x, t + \Delta t)$, implies conversion $\hat{g} \rightarrow f$

Step 3: Correction $g_\alpha(x, t + \Delta t) = \hat{g}_\alpha(x, t) - F^{t+\frac{1}{2}}$, i. c. $g \rightarrow f$

Step 4: Step 1

Discretization in space and time of the DVBE

Evaluation of $F^{t+\frac{1}{2}} = \xi_\alpha \cdot n_\beta [f_\alpha(x_\beta, t + \Delta t/2)]$

- 1 Heun predictor-corrector scheme
- 2 DUGKS

$$f_\alpha(x_\beta, t + \Delta t/2) = g_\alpha^f(x_\beta, t + h) + \frac{\Delta t^f}{2\tau_g^f} (f_\alpha^{eq}(x_\beta, t + h) - g_\alpha^f(x_\beta, t + h))$$

with $h = \Delta t/2$ being the timestep of the stream-collide algorithm on a refined lattice.

$$g^f(x_\beta, t + h) = \hat{g}^f(x_\beta - \xi_\alpha h, t) \quad (10)$$

where the term on the RHS is approximated by

$$\hat{g}^f(x_\beta - \xi_\alpha h, t) = \hat{g}^f(x_\beta, t) - \xi_\alpha \cdot \sigma_\beta \quad (11)$$

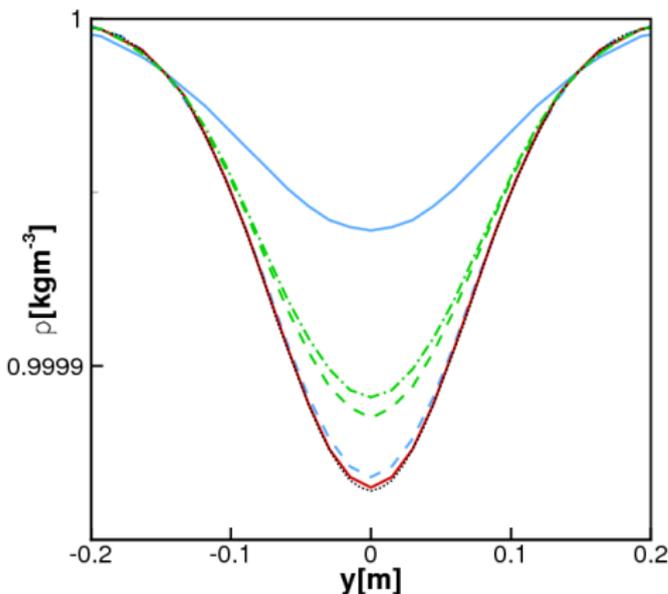
$\hat{g}^f(x_\beta, t)$ is obtained from the pre-collision coarse grained functions $g(x, t)$ (approximated with QUICK or centred scheme) and

$$\sigma_\beta = \frac{\hat{g}_\alpha^f(x + \Delta x, t) - \hat{g}_\alpha^f(x, t)}{\Delta x} \quad (12)$$

Single step algorithm!

Comparison of the two schemes

scheme	streaming	finite volume			
		DUGKS	<i>Heun</i> predictor-corrector	simple Euler	
CFL	1	0.5	0.5	1	0.125
flux approximation	-	centred, QUICK	centred	QUICK	QUICK



Solution of vortex after 5000 iterations:

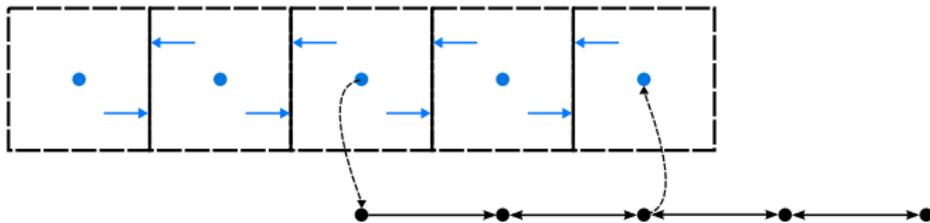
- ... analytic
- stream
- Heun QUICK $CFL = 1$
- - Heun centred $CFL = 0.5$
- - DUGKS centred $CFL = 0.5$
- · - DUGKS QUICK $CFL = 0.5$

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Algorithm

- In order to be congruent in the time step, we choose to couple the stream algorithm with the QUICK-Heun finite volume scheme that is stable at CFL=1 (in combination with a 7 point explicit filter).
- First results have shown that in order to couple the two algorithms, a multi-grid approach with a certain interface of at least $2\Delta x$ is required.



Sketch of the algorithm coupling on a uniform mesh

Algorithm

Streaming + *Heun* predictor corrector:

Step 1: Collision in every node of the domain

Step 2a: Memory shift in the domain where streaming is used

Step 2b: Prediction of $g_{\alpha}^*(x, t + \Delta t)$ with a simple Euler scheme

Step 3: Correction of $g_{\alpha}(x, t + \Delta t)$ with *Heun* scheme

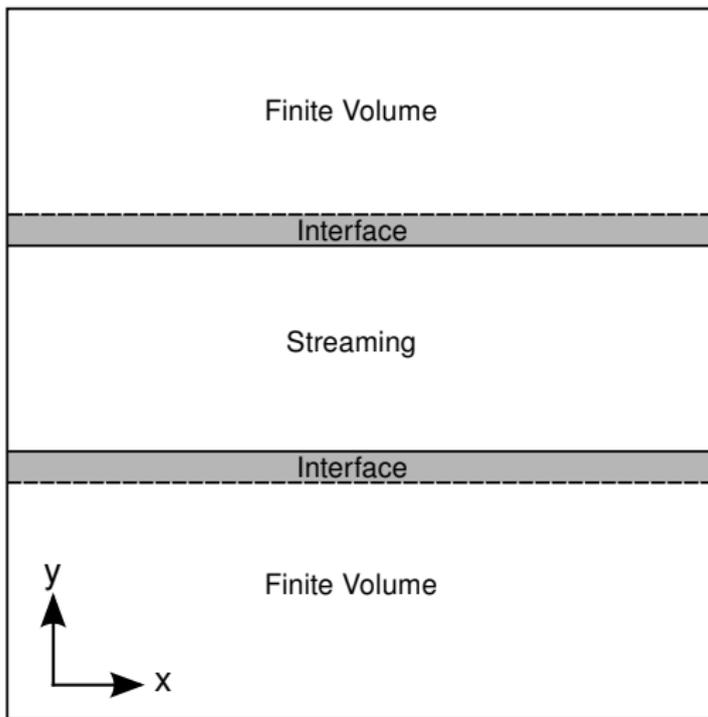
Step 4: Mutual transfer of information in the transition nodes

Step 5: Step 1

blue: two step algorithm, red: additional steps

Results of two test cases are presented in the following.

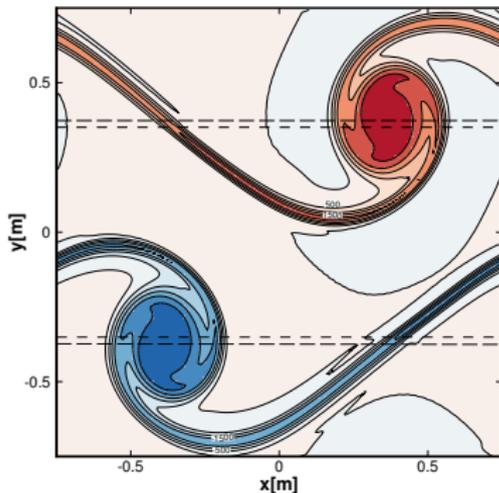
Numerical setup



mesh

- in-house LBM code
- periodic BC in X and Y
- $N = 101$, $N = 201$ and $N = 401$ respectively

Double Shear Layer



vorticity

$$u(x, y, t_0) = U_0 \tanh\left(\frac{y - L/4}{d_0}\right) \tanh\left(\frac{3/4L - y}{d_0}\right)$$

$$v(x, y, t_0) = U_0 a_0 \sin(2\pi(x/L + 0.25))$$

$$\rho(x, y, t_0) = \rho_0$$

$$\text{with } d_0 = 6\Delta x \text{ and } a_0 = 0.01\rho_0$$

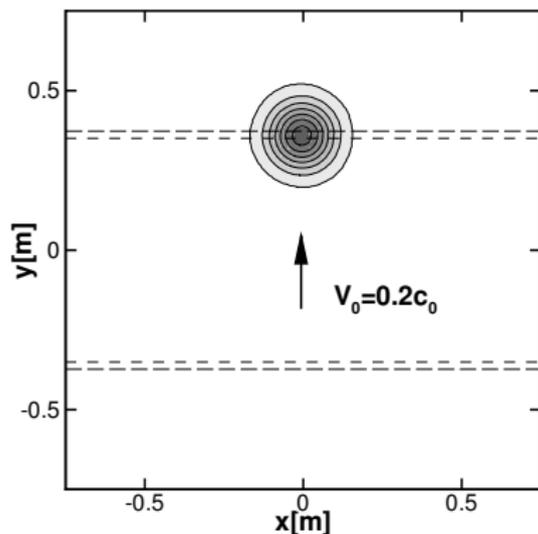
Vortex

$$u(x, y, t_0) = -\frac{\Gamma}{R^2} y \times \exp\left(\frac{1}{2} \left(1 - \frac{r^2}{2R^2}\right)\right)$$

$$v(x, y, t_0) = V_0 + \frac{\Gamma}{R^2} x \times \exp\left(\frac{1}{2} \left(1 - \frac{r^2}{2R^2}\right)\right)$$

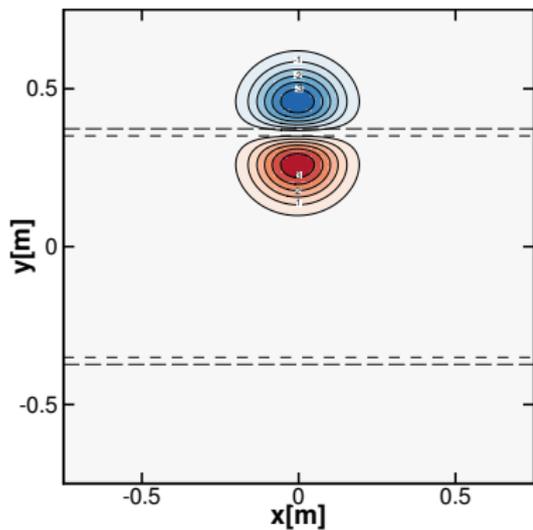
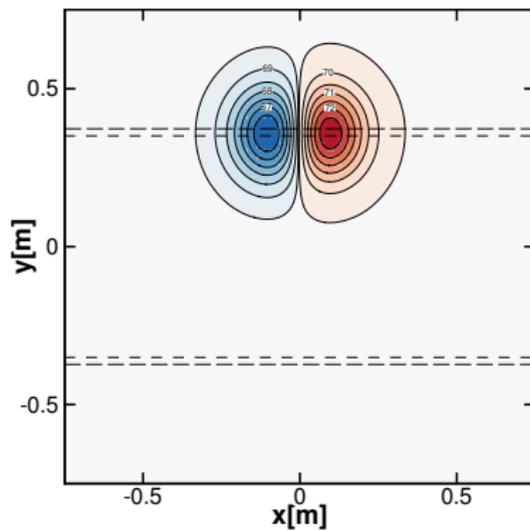
$$\rho(x, y, t_0) = \rho_0 - \frac{\Gamma^2}{2c_0^2} \rho_0 \times \exp\left(1 - \frac{r^2}{R^2}\right)$$

with $M = 0.2$, $\Gamma = 0.1 V_0$ and $R = 6\Delta x$
for $N=101$.



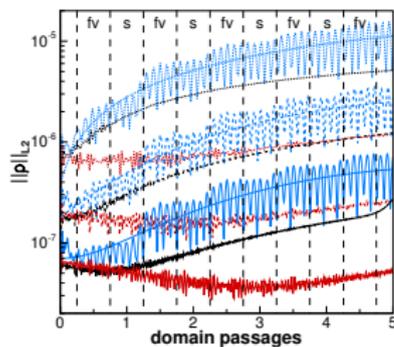
density

Vortex

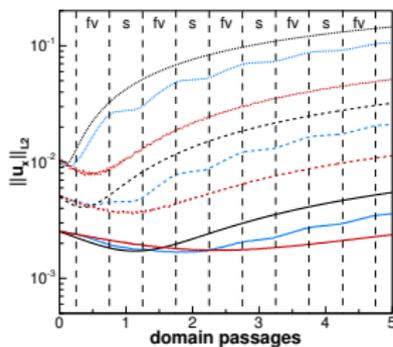
 u  v

velocity

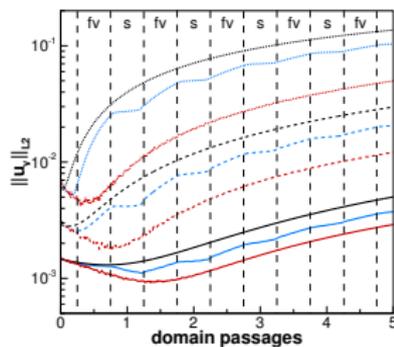
Vortex



$\|\rho\|_{L_2}$



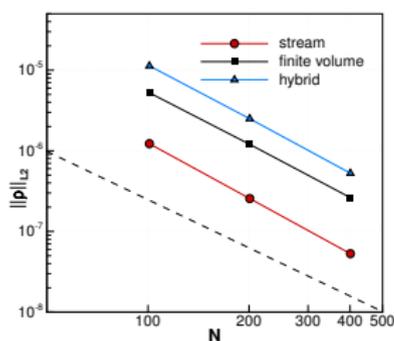
$\|u_x\|_{L_2}$



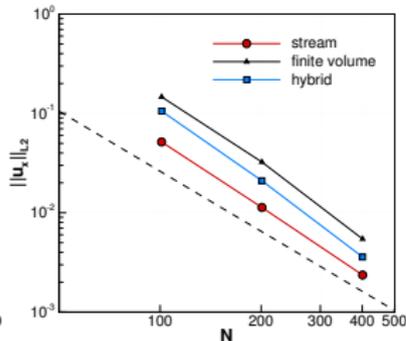
$\|u_y\|_{L_2}$

— stream, — finite volume, — hybrid
 ... N101, --- N201, — N401

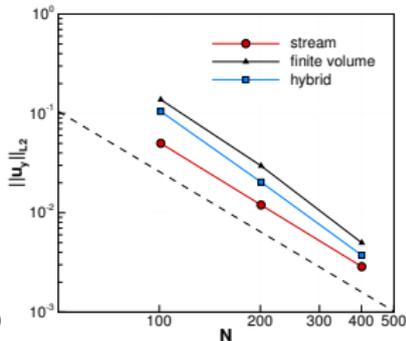
Vortex



$\|\rho\|_{L_2}$



$\|u_x\|_{L_2}$



$\|u_y\|_{L_2}$

L_2 of macroscopic variables for varying mesh size N after 5 domain passages:

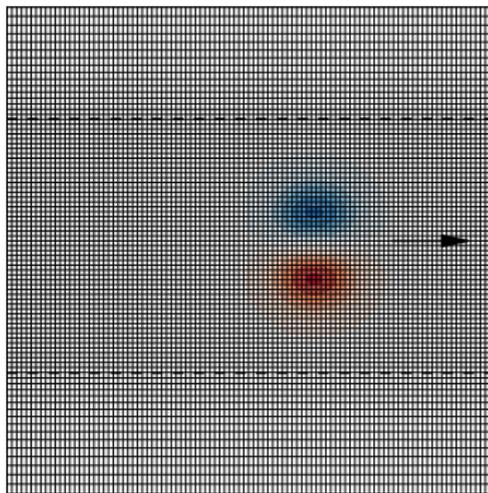
— stream, — finite volume, — hybrid

Content

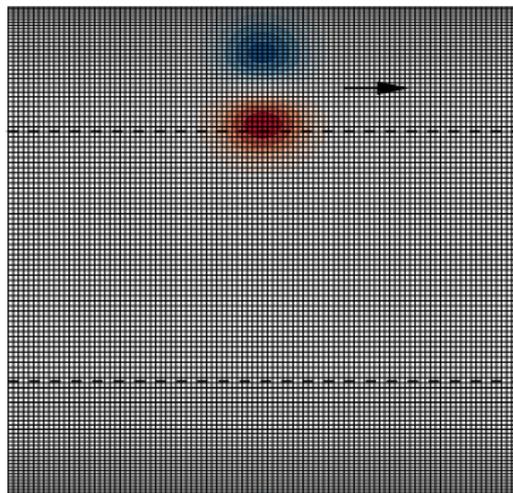
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Conclusion and application

- **First study** to show that it is possible to **couple** the streaming step in a LBM algorithm with a finite volume formulation
- This is of interest for regions with anisotropic flows



Coarse grading: $\Delta y \rightarrow 2\Delta y$



Fine grading: $\Delta y \rightarrow 0.5\Delta y$

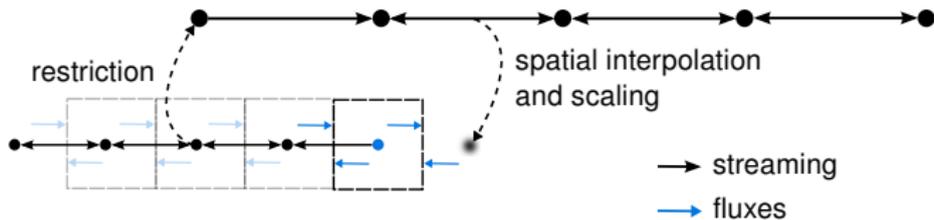
Mesh adaptation according to local flow features

Thank you for your attention!

References I

- [1] K. Shrestha G. Mompean, E. Calzavarini. *Finite-volume versus streaming-based lattice Boltzmann algorithm for fluid-dynamics simulations: A one-to-one accuracy and performance study*. Physical Review E 93, 2016.

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Sketch of the initially proposed algorithm for a non-uniform mesh