

Spatial branching processes: Superprocesses and snakes

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Outline

Spatial branching processes model the evolution of populations where individuals both

- reproduce themselves according to some branching distribution
- move in space according to a certain Markov process (e.g. Brownian motion)

Superprocesses (also called measure-valued branching processes) occur in the limit where:

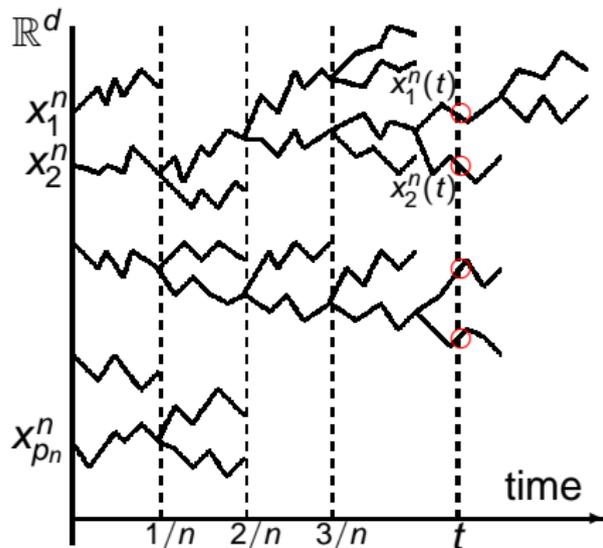
- the population is very large (but each individual has a very small “mass”)
- the mean time between two branching events is very small

Related model: **Fleming-Viot processes** used in population genetics (spatial position = genetic type of the individual)

Why study spatial branching processes, and in particular superprocesses ?

- These objects appear in the asymptotics of many other important probabilistic models:
 - ▶ **interacting particle systems**: voter model, contact process, etc. (Cox, Durrett, Perkins, ...)
 - ▶ models from **statistical physics**: lattice trees, oriented percolation, etc. (Slade, van der Hofstad, Hara, ...)
 - ▶ models from **mathematical biology**, where there is competition between several species (e.g. Lotka-Volterra models)
- Connections with the theory of stochastic partial differential equations.
- Connections with **partial differential equations** (probabilistic approach to an important class of nonlinear PDEs, cf Dynkin, Kuznetsov, LG, ...)
- Description of asymptotics in models of **combinatorics** (cf Lecture 3).

1. Branching particle systems and superprocesses



At time $t = 0$, p_n particles located at $x_1^n, x_2^n, \dots, x_{p_n}^n \in \mathbb{R}^d$.

Particles independently

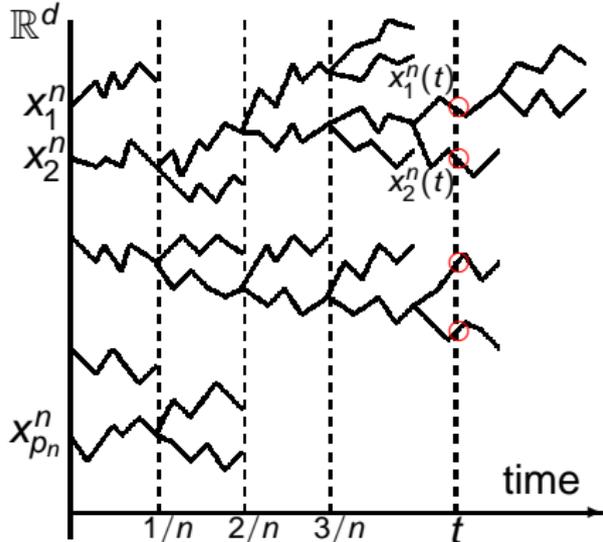
- move in space according to Brownian motion
- die at times $1/n, 2/n, 3/n, \dots$
- when a particle dies, it gives rise to children according to the offspring distribution γ

For every $t \geq 0$, $x_1^n(t), x_2^n(t), \dots$ positions of particles alive at time t ,

$$Z_t^n = \frac{1}{n} \sum_i \delta_{x_i^n(t)}$$

rescaled sum of Dirac masses at particles alive at time t .

Now let $n \rightarrow \infty \dots$



$$\text{Recall } Z_t^n = \frac{1}{n} \sum_i \delta_{x_i^n(t)}.$$

$$M_f(\mathbb{R}^d) = \{\text{finite measures on } \mathbb{R}^d\}.$$

Assumptions

- Convergence of initial values:

$$Z_0^n = \frac{1}{n} \sum_{i=1}^{p_n} \delta_{x_i^n} \xrightarrow[n \rightarrow \infty]{} \mu \in M_f(\mathbb{R}^d)$$

- The offspring distribution γ has mean 1 and finite variance ρ^2 .

Theorem (Watanabe)

Then,

$$(Z_t^n)_{t \geq 0} \xrightarrow[n \rightarrow \infty]{(d)} (Z_t)_{t \geq 0}$$

where $(Z_t)_{t \geq 0}$ is a Markov process with values in $M_f(\mathbb{R}^d)$, called *super-Brownian motion*.

$Z_t \in M_f(\mathbb{R}^d)$ is supported on “a cloud of particles alive at time t ”

Characterizing the law of super-Brownian motion

Notation: $C_b^+(\mathbb{R}^d) = \{\text{bounded continuous functions } g : \mathbb{R}^d \rightarrow \mathbb{R}_+\}$
 $\langle \mu, g \rangle = \int g d\mu$, for $\mu \in M_f(\mathbb{R}^d)$ and $g \in C_b^+(\mathbb{R}^d)$.

Then, for every $g \in C_b^+(\mathbb{R}^d)$,

$$E \left[\exp(-\langle Z_t, g \rangle) \mid Z_0 = \mu \right] = \exp -\langle \mu, u_t \rangle$$

where $(u_t(x), t \geq 0, x \in \mathbb{R}^d)$ is the unique nonnegative solution of

$$\begin{aligned} \frac{\partial u}{\partial t} &= \frac{1}{2} \Delta u - \frac{\rho^2}{2} u^2 \\ u_0 &= g \end{aligned}$$

The function $\psi(u) = \frac{\rho^2}{2} u^2$ is called the **branching mechanism** of Z .

Remark. The law of Z depends on the offspring distribution μ of the approximating system only through the parameter ρ^2 .

*Other characterizations via **martingale problems**, more appropriate for models with interactions.*

Path properties of super-Brownian motion (Dawson, Perkins, Shiga, ...)

$d = 1$: Then Z_t has a density with respect to Lebesgue measure

$$Z_t(dx) = Y_t(x) dx$$

and this density solves the SPDE

$$dY_t = \frac{1}{2} \Delta Y_t dt + c \sqrt{Y_t} dW_t$$

where W is space-time white noise.

$d \geq 2$: Then Z_t is almost surely supported on a set of zero Lebesgue measure, and **uniformly spread** on its support, in the sense of Hausdorff measure.

2. The Brownian snake approach

Idea. One can generate the **individual particle paths** (the “historical paths”) of a super-Brownian motion, as the values of a path-valued Markov process called the **Brownian snake**.

→ This construction is closely related to the fact that the underlying **genealogical structure** of a super-Brownian motion can be coded by Brownian excursions (in the same sense as the CRT is coded by a normalized Brownian excursion, cf Lecture 1).

The construction of the Brownian snake. Fix $x \in \mathbb{R}^d$ and set

$$\begin{aligned}\mathcal{W}_x &= \{\text{finite paths started from } x\} \\ &= \{w : [0, \zeta_w] \longrightarrow \mathbb{R}^d \text{ continuous}, w(0) = x\}.\end{aligned}$$

If $w \in \mathcal{W}_x$, ζ_w is called the **lifetime** of w .

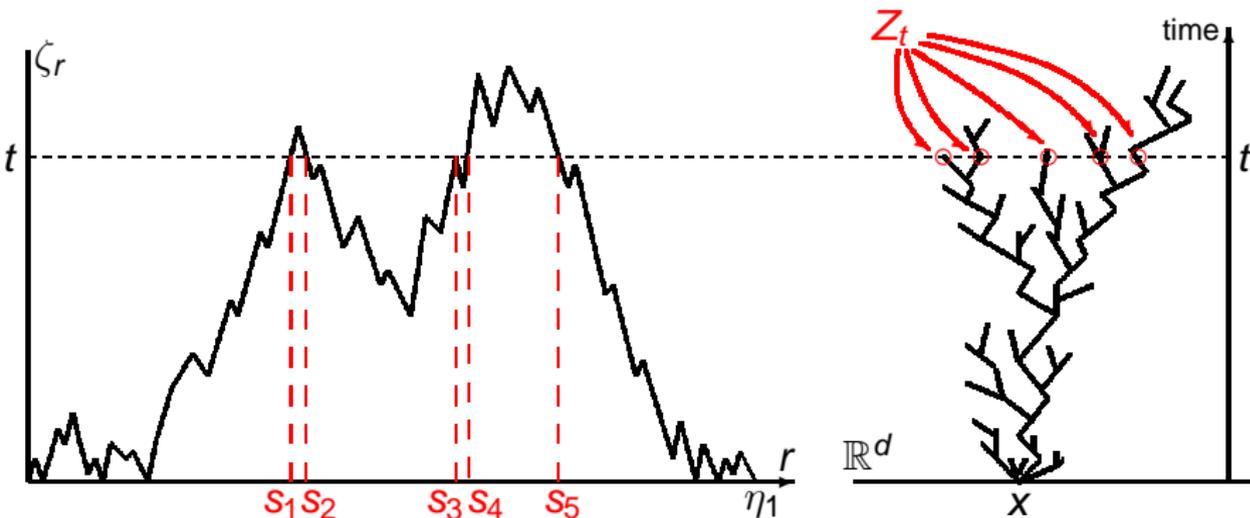
The terminal point or **tip** of w is $\hat{w} = w(\zeta_w)$.

Heuristic description of the Brownian snake $(W_s)_{s \geq 0}$

- For every $s \geq 0$, W_s is a random path in \mathbb{R}^d started at x , with a **random lifetime** ζ_s .
- The lifetime ζ_s evolves like linear Brownian motion reflected at 0 (**a lifetime cannot be negative !**)
- When ζ_s decreases, the path W_s is **shortened** from its tip.
- When ζ_s increases, the path W_s is **extended** by adding “little pieces” of d -dimensional Brownian motion at its tip.

Why consider such a process ?

In particular, because of its connection with super-Brownian motion.



For every $t \geq 0$, let $L^t = (L_s^t)_{s \geq 0}$ be the local time at level t of $(\zeta_s)_{s \geq 0}$ (the measure $L^t(ds)$ is supported on $\{s \geq 0 : \zeta_s = t\}$).

Theorem

Let $\eta_1 := \inf\{s \geq 0 : L_s^0 = 1\}$. The measure-valued process $(Z_t)_{t \geq 0}$

$$\langle Z_t, g \rangle = \int_0^{\eta_1} L^t(ds) g(W_s(t))$$

is a **super-Brownian motion** started from δ_x .

Applications

Many results about **super-Brownian motion** can be stated equivalently and proved more easily in terms of the **Brownian snake**.

This is true in particular for path properties:

- The values W_s of the Brownian snake are Hölder continuous with exponent $\frac{1}{2} - \varepsilon$. The **topological support** $\text{supp}(Z_t)$ of super-BM cannot move faster: for every $t \geq 0$, $0 < r < r_0(\omega)$,

$$\text{supp}(Z_{t+r}) \subset U_{r^{1/2-\varepsilon}}(\text{supp}(Z_t))$$

where $U_\delta(K)$ denotes the δ -enlargement of K .

- If $\widehat{W}_s = W_s(\zeta_s)$ denotes the tip of the path W_s , the map $s \rightarrow \widehat{W}_s$ is Hölder continuous with exponent $\frac{1}{4} - \varepsilon$. From the snake approach,

$$\{\widehat{W}_s : 0 \leq s \leq \eta_1\} = \overline{\bigcup_{t \geq 0} \text{supp}(Z_t)} =: \mathcal{R}$$

is the **range** of Z , that is the set of points touched by the cloud of particles. It follows that: **$\dim(\mathcal{R}) = 4 \wedge d$**

More precise results: **Perkins, Dawson, Iscoe, LG, etc.**

3. Connections with partial differential equations

Probabilistic potential theory: Classical connections between Brownian motion and the Laplace equation $\Delta u = 0$ or the heat equation $\frac{\partial u}{\partial t} = \Delta u$ (Doob, Kakutani, etc.)

In our setting: Similar remarkable connections between super-Brownian motion or the Brownian snake and semilinear equations of the form $\Delta u = u^\gamma$ or $\frac{\partial u}{\partial t} = \Delta u - u^\gamma$ (Dynkin, Kuznetsov, LG, etc.)

Why study these connections ? Because they

- Allow **explicit analytic calculations** of probabilistic quantities related to the Brownian snake and super-BM
- Give a **probabilistic representation** of solutions of PDE that has led to new analytic results

For simple statements of the connections with PDE, needs **excursion measures**.

The Itô excursion measure

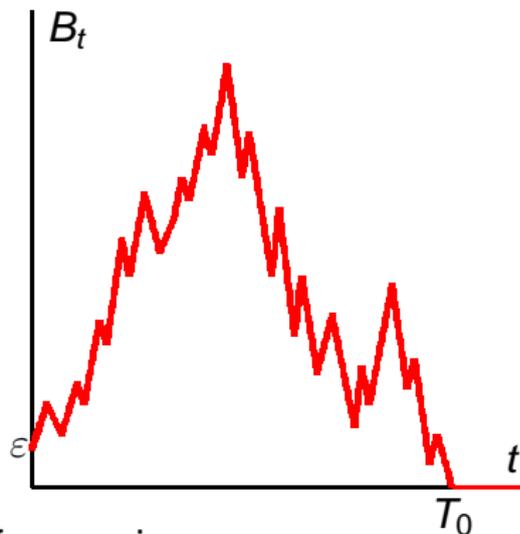
Consider a Brownian motion $(B_t)_{t \geq 0}$ with $B_0 = \varepsilon$.

Set $T_0 = \inf\{t \geq 0 : B_t = 0\}$.

Let P_ε be the law of $(B_{t \wedge T_0})_{t \geq 0}$

Then,

$$\varepsilon^{-1} P_\varepsilon \xrightarrow{\varepsilon \rightarrow 0} \Pi$$



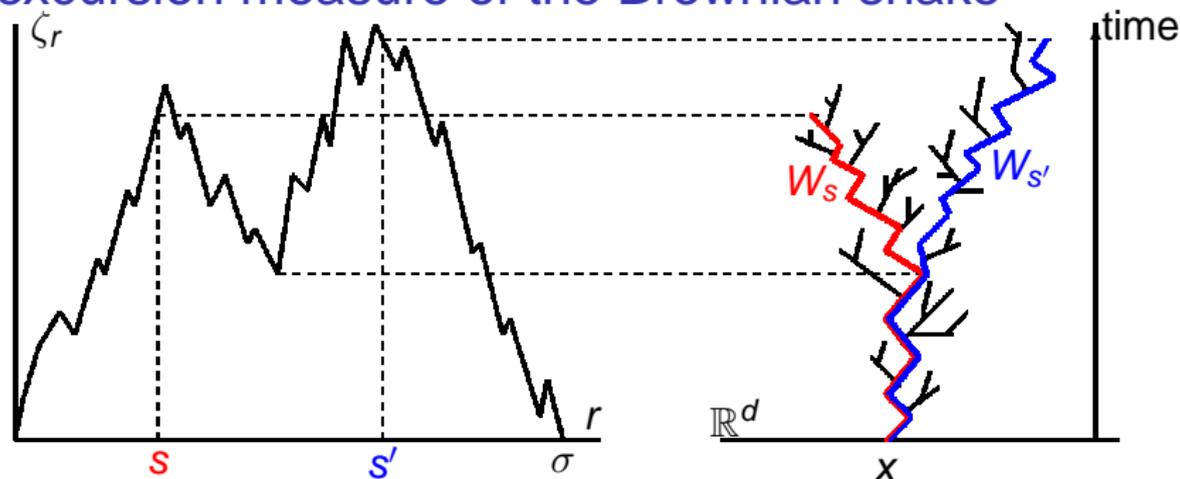
Π is a σ -finite measure on the set of excursions

$$E = \{e : [0, \infty) \rightarrow [0, \infty) \text{ continuous,} \\ \exists \sigma(e) > 0, e(s) > 0 \text{ iff } 0 < s < \sigma(e)\}$$

Π is called the **Itô excursion measure**.

(Note: $\Pi(\cdot \mid \sigma = 1)$ is the law of the **normalized** excursion, cf Lect.1)

The excursion measure of the Brownian snake



\mathbb{N}_x is the measure under which:

- $(\zeta_s)_{s \geq 0}$ is distributed according to $\Pi(de)$ (the Itô measure)
- Conditionally given $(\zeta_s)_{s \geq 0}$, $(W_s)_{s \geq 0}$ is distributed as **the snake driven by** $(\zeta_s)_{s \geq 0}$, with initial point x : W_s has lifetime ζ_s , and if $s < s'$, the conditional law of $W_{s'}$ given W_s is as described before.

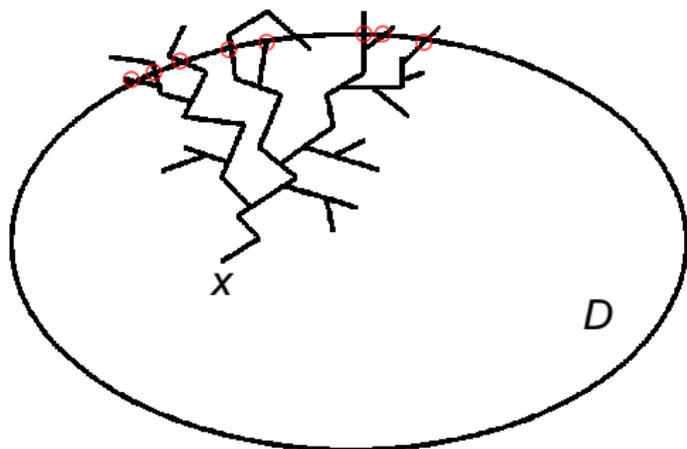
Under \mathbb{N}_x , the paths W_s , $s \in [0, \sigma]$ form a **“tree of Brownian paths”** with initial point x .

Warning. \mathbb{N}_x is an infinite measure (because so is Π).

Exit points from a domain

Classical theory of relations between Brownian motion and PDEs : A key role is played by the **first exit point** of Brownian motion from a domain D .

Here one constructs a measure supported on the **set of exit points** of the paths W_s from D (assuming that the initial point $x \in D$)



For every finite path $w \in \mathcal{W}_x$, set

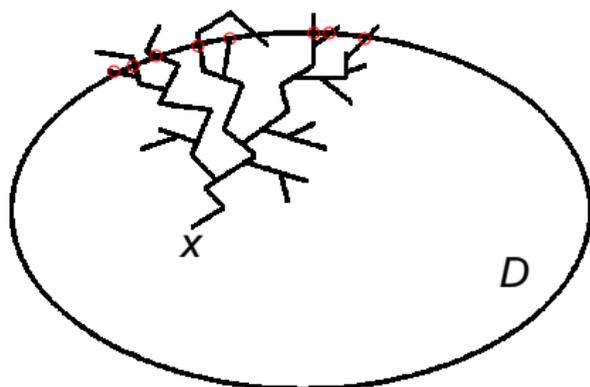
$$\tau(w) = \inf\{t \geq 0 : w(t) \notin D\}$$

and

$$\mathcal{E}^D = \{W_s(\tau(W_s)) : \tau(W_s) < \infty\}$$

(exit points of the paths W_s)

The exit measure of the Brownian snake



$$\mathcal{E}^D = \{\text{exit points of the paths } W_s\}$$

Proposition

The formula

$$\langle Z^D, g \rangle = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_0^\sigma ds \mathbf{1}_{\{\tau(W_s) < \zeta_s < \tau(W_s) + \varepsilon\}} g(W_s(\tau(W_s)))$$

defines \mathbb{N}_x a.e. a finite measure Z^D supported on \mathcal{E}^D .

Z^D is called the **exit measure** from D (Dynkin)

The key connection with PDE

Theorem (Reformulation of Dynkin 1991)

Let D be a regular domain (in the classical potential-theoretic sense), and $g \in C_b^+(\partial D)$. The formula

$$u(x) = \mathbb{N}_x(1 - \exp - \langle Z^D, g \rangle), \quad x \in \partial D \quad (1)$$

defines the unique (nonnegative) solution of the Dirichlet problem

$$\begin{aligned} \Delta u &= u^2 && \text{in } D \\ u|_{\partial D} &= g \end{aligned}$$

Remark. Similarity with the probabilistic formula $u(x) = \mathbb{E}_x[g(B_\tau)]$ for the classical Dirichlet problem.

Important point: Formula (1) is very robust with respect to passages to the limit, and yields probabilistic representations for “virtually any” positive solution of $\Delta u = u^2$ in a domain.

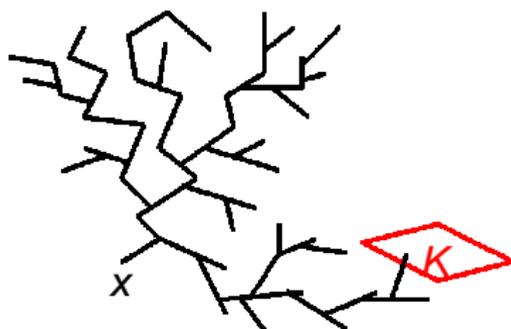
Maximal solutions

Corollary (Dynkin)

Let D be any domain. The formula

$$u(x) = \mathbb{N}_x(\mathcal{E}^D \neq \emptyset), \quad x \in D$$

gives the maximal nonnegative solution of $\Delta u = u^2$ in D .

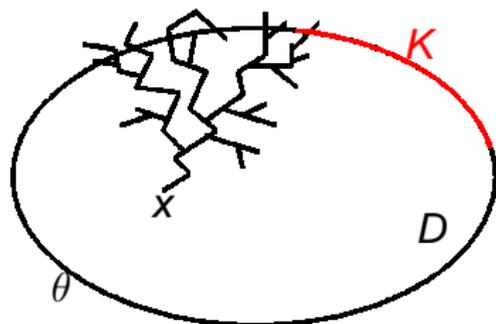


Application. $D = \mathbb{R}^d \setminus K$, K compact

The Brownian snake hits K with positive probability

- \Leftrightarrow There exists a non trivial solution of $\Delta u = u^2$ in $\mathbb{R}^d \setminus K$
- \Leftrightarrow K is **not** a removable singularity for $\Delta u = u^2$
- \Leftrightarrow $\text{cap}_{d-4}(K) > 0$ (Baras-Pierre)

The representation of solutions when $d = 2$



D smooth domain in \mathbb{R}^2

Fact. If $x \in D$, the exit measure Z^D has \mathbb{N}_x a.e. a continuous density with respect to Lebesgue measure on ∂D , denoted by $(z_D(y), y \in \partial D)$.

Recall $\mathcal{E}^D = \{\text{exit points of the paths } W_s\}$

Theorem (LG)

The formula

$$u_{K,\theta}(x) = \mathbb{N}_x \left(1 - \mathbf{1}_{\{\mathcal{E}^D \cap K = \emptyset\}} \exp - \langle \theta, z_D \rangle \right)$$

gives a **bijection** between $\{\text{positive solutions of } \Delta u = u^2 \text{ in } D\}$ and the set of all pairs (K, θ) , where:

- K is a **compact subset** of ∂D
- θ is a **Radon measure** on $\partial D \setminus K$

Extensions of the representation theorem

Consider more generally the equation

$$\Delta u = u^p$$

for any $p > 1$, in dimension $d \geq 2$

- **Subcritical case** $p < \frac{d+1}{d-1}$ (includes $p = 2, d = 2$)

The correspondence between solutions and traces (K, θ) remains valid as in the preceding theorem (cf [Marcus-Véron](#) (analytic methods), [Dynkin-Kuznetsov](#) and [LG-Mytnik](#))

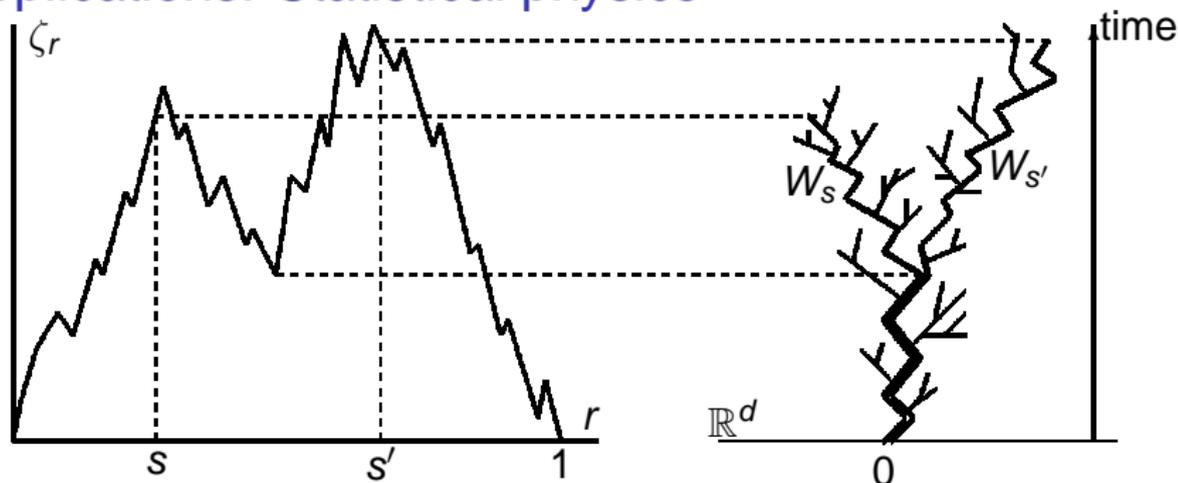
- **Supercritical case** $p \geq \frac{d+1}{d-1}$

Needs to introduce a notion of **fine trace** of a solution ([Dynkin](#))

[Dynkin](#) conjectured a **one-to-one correspondence** between solutions and admissible fine traces.

- ▶ Proved by [Mselati](#) (Memoirs AMS 2003) for $p = 2$ (using the Brownian snake)
- ▶ Proved by [Dynkin-Kuznetsov](#) for $1 < p < 2$
- ▶ Still open for $p > 2$ but recent analytic progress by [Marcus-Véron](#)

4. Applications: Statistical physics



Consider the Brownian snake (W_s)

- with initial point $x = 0$
- driven by a **normalized excursion** (condition on $\sigma = 1$)

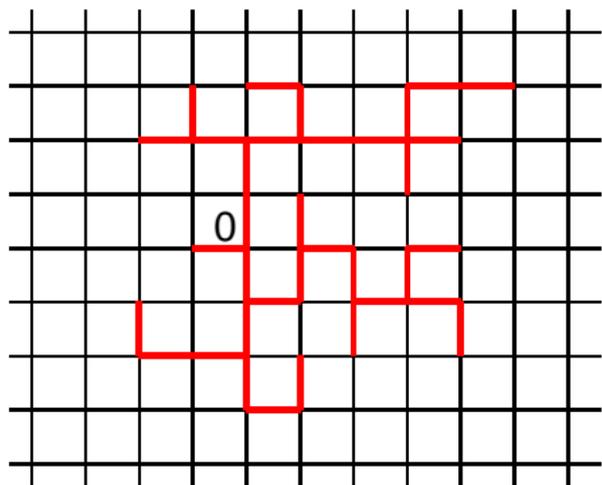
The random probability measure \mathcal{I} on \mathbb{R}^d defined by

$$\langle \mathcal{I}, g \rangle = \int_0^1 ds g(\widehat{W}_s) \quad (\text{recall } \widehat{W}_s = \text{terminal point of } W_s)$$

is called **ISE** (for integrated super-Brownian excursion, [Aldous](#)).

ISE has appeared in a number of limit theorems for models of **statistical physics** in high dimensions: Lattice trees, percolation clusters, etc.

A **lattice tree** is a finite subgraph of \mathbb{Z}^d with no loop.



A lattice tree in \mathbb{Z}^2
with 36 vertices

Question. What can we say about the shape (for instance the diameter) of a typical large lattice tree in \mathbb{Z}^d ?

→ Very hard question if d is small (self-avoiding constraint)

Let

$$\mathcal{T}_n = \{\text{lattice trees with } n \text{ vertices in } \mathbb{Z}^d \text{ containing } 0\}.$$

Let T_n be chosen uniformly over \mathcal{T}_n and let X_n be the random measure that assigns mass $\frac{1}{n}$ to each point of the form $c n^{-1/4} x$, x vertex of T_n . (X_n is uniformly spread over the rescaled tree $c n^{-1/4} T_n$)

Theorem (Derbez-Slade)

If d is large enough, we can choose $c = c_d > 0$ so that

$$X_n \xrightarrow[n \rightarrow \infty]{(d)} \mathcal{I}$$

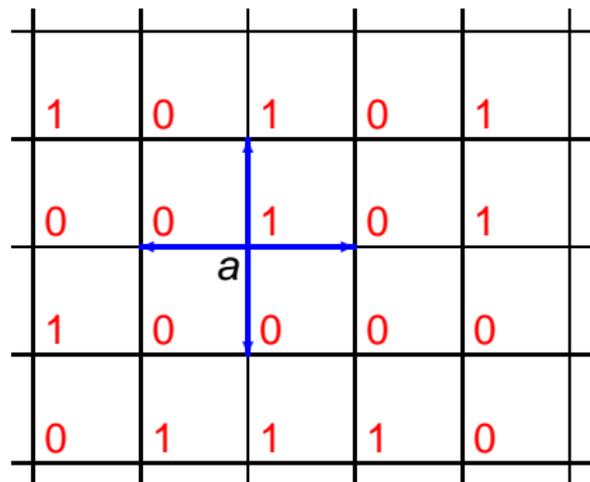
where \mathcal{I} is ISE.

Informally, a typical large lattice tree (suitably rescaled) looks like the support of ISE, or equivalently the range of a Brownian snake driven by a normalized Brownian excursion.

Conjecture. The preceding theorem holds for $d > 8$ (but not for $d \leq 8$).

5. Applications: Interacting particle systems

The voter model.



At each point of \mathbb{Z}^d sits an individual who can have opinion 0 or 1.

For each $a \in \mathbb{Z}^d$, after an exponential time with parameter 1, the individual sitting at a

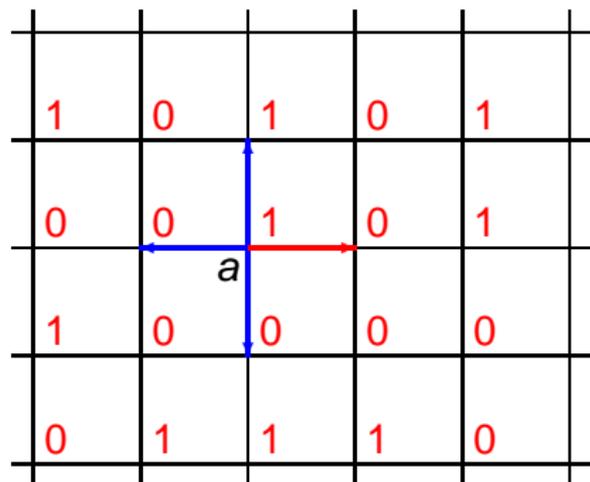
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- adopts his opinion

And so on.

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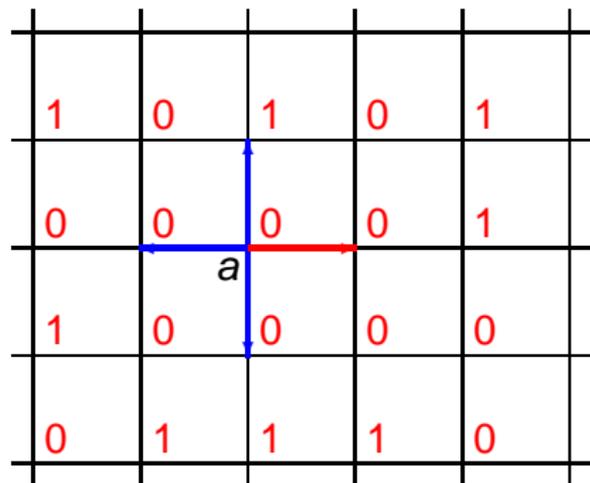
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Question. How do opinions propagate in space ?

Suppose $d \geq 2$. Write $\xi_t(\mathbf{a})$ for the opinion of \mathbf{a} at time t .

Suppose that

$$\xi_0(\mathbf{a}) = \begin{cases} 0 & \text{if } \mathbf{a} \neq 0 \\ 1 & \text{if } \mathbf{a} = 0 \end{cases}$$

(At time $t = 0$ only the origin has opinion 1)

Set $\mathcal{V}_t = \{\mathbf{a} \in \mathbb{Z}^d : \xi_t(\mathbf{a}) = 1\}$.

Bramson-Griffeath: estimates for $P(\mathcal{V}_t \neq \emptyset)$ (opinion 1 survives).

Theorem (Bramson-Cox-LG)

The law of $\frac{1}{\sqrt{t}} \mathcal{V}_t$ conditional on $\{\mathcal{V}_t \neq \emptyset\}$ converges as $t \rightarrow \infty$ to the law of the random set

$$\{W_s(1) : s \geq 0, \zeta_s \geq 1\}$$

under the conditional measure $\mathbb{N}_0(\cdot \mid \sup_{s \geq 0} \zeta_s > 1)$.

Asymptotically interactions disappear and opinions propagate like a spatial branching process: see also **Cox-Durrett-Perkins, ...**