# Random Trees, Lévy Processes and Spatial Branching Processes 

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## Introduction

The main goal of this work is to investigate the genealogical structure of continuousstate branching processes in connection with limit theorems for discrete GaltonWatson trees. Applications are also given to the construction and various properties of spatial branching processes including a general class of superprocesses.

Our starting point is the recent work of Le Gall and Le Jan [32] who proposed a coding of the genealogy of general continuous-state branching processes via a realvalued random process called the height process. Recall that continuous-state branching processes are the continuous analogues of discrete Galton-Watson branching processes, and that the law of any such process is characterized by a real function $\psi$ called the branching mechanism. Roughly speaking, the height process is a continuous analogue of the contour process of a discrete branching tree, which is easy to visualize (see Section 0.1, and note that the previous informal interpretation of the height process is made mathematically precise by the results of Chapter 2). In the important special case of the Feller branching diffusion $\left(\psi(u)=u^{2}\right)$, the height process is reflected linear Brownian motion: This unexpected connection between branching processes and Brownian motion, or random walk in a discrete setting has been known for long and exploited by a number of authors (see e.g. [3], [11], [18], [39], [42]). The key contribution of [32] was to observe that for a general subcritical continuous-state branching process, there is an explicit formula expressing the height process as a functional of a spectrally positive Lévy process whose Laplace exponent $\psi$ is precisely the branching mechanism. This suggests that many problems concerning the genealogy of continuous-state branching processes can be restated and solved in terms of spectrally positive Lévy processes, for which a lot of information is available (see e.g. Bertoin's recent monograph [5]). It is the principal aim of the present work to develop such applications.

In the first two sections below, we briefly describe the objects of interest in a discrete setting. In the next sections, we outline the main contributions of the present work.

### 0.1 Discrete trees

Let

$$
\mathcal{U}=\bigcup_{n=0}^{\infty} \mathbb{N}^{n}
$$

where $\mathbb{N}=\{1,2, \ldots\}$ and by convention $\mathbb{N}^{0}=\{\emptyset\}$. If $u=\left(u_{1}, \ldots, u_{n}\right) \in \mathbb{N}^{n}$, we set $|u|=n$, so that $|u|$ represents the "generation" of $u$. If $u=\left(u_{1}, \ldots u_{m}\right)$ and $v=$ $\left(v_{1}, \ldots, v_{n}\right)$ belong to $\mathcal{U}$, we write $u v=\left(u_{1}, \ldots u_{m}, v_{1}, \ldots, v_{n}\right)$ for the concatenation of $u$ and $v$. In particular $u \emptyset=\emptyset u=u$.

A (finite) rooted ordered tree $\mathcal{T}$ is a finite subset of $\mathcal{U}$ such that:
(i) $\emptyset \in \mathcal{T}$.
(ii) If $v \in \mathcal{T}$ and $v=u j$ for some $u \in \mathcal{U}$ and $j \in \mathbb{N}$, then $u \in \mathcal{T}$.
(iii) For every $u \in \mathcal{T}$, there exists a number $k_{u}(\mathcal{T}) \geq 0$ such that $u j \in \mathcal{T}$ if and only if $1 \leq j \leq k_{u}(\mathcal{T})$.

We denote by $\mathbf{T}$ the set of all rooted ordered trees. In what follows, we see each vertex of the tree $\mathcal{T}$ as an individual of a population whose $\mathcal{T}$ is the family tree. The cardinality $\#(\mathcal{T})$ of $\mathcal{T}$ is the total progeny.

If $\mathcal{T}$ is a tree and $u \in \mathcal{T}$, we define the shift of $\mathcal{T}$ at $u$ by $\theta_{u} \mathcal{T}=\{v \in U: u v \in \mathcal{T}\}$. Note that $\theta_{u} \tau \in \mathbf{T}$.

We now introduce the (discrete) height function associated with a tree $\mathcal{T}$. Let us denote by $u(0)=\emptyset, u(1), u(2), \ldots, u(\#(\mathcal{T})-1)$ the elements of $\mathcal{T}$ listed in lexicographical order. The height function $H(\mathcal{T})=\left(H_{n}(\mathcal{T}) ; 0 \leq n<\#(\mathcal{T})\right)$ is defined by

$$
H_{n}(\mathcal{T})=|u(n)|, \quad 0 \leq n<\#(\mathcal{T}) .
$$

The height function is thus the sequence of the generations of the individuals of $\mathcal{T}$, when these individuals are visited in the lexicographical order (see Fig. 1 for an example). It is easy to check that $H(\mathcal{T})$ characterizes the tree $\mathcal{T}$.


Figure 1
The contour function gives another way of characterizing the tree, which is easier to visualize on a picture (see Fig.1). Suppose that the tree is embedded in the half-plane in such a way that edges have length one. Informally, we imagine the motion of a particle that starts at time $t=0$ from the root of the tree and then explores the tree from the left to the right, moving continuously along the edges at unit speed, until it comes back to its starting point. Since it is clear that each edge will be crossed twice in this evolution, the total time needed to explore the tree is $\zeta(\mathcal{T}):=2(\#(\mathcal{T})-1)$. The value $C_{t}$ of the contour function at time $t$ is the distance (on the tree) between the position of the particle at time $t$ and the root. By convention $C_{t}=0$ if $t \geq \zeta(\mathcal{T})$. Fig. 1 explains the definition of the contour function better than a formal definition.

### 0.2 Galton-Watson trees

Let $\mu$ be a critical or subcritical offspring distribution. This means that $\mu$ is a probability measure on $\mathbb{Z}_{+}$such that

$$
\sum_{k=0}^{\infty} k \mu(k) \leq 1
$$

We exclude the trivial case where $\mu(1)=1$.
There is a unique probability distribution $\mathbf{Q}_{\mu}$ on $\mathbf{T}$ such that
(i) $\mathbf{Q}_{\mu}\left(k_{\emptyset}=j\right)=\mu(j), \quad j \in \mathbb{Z}_{+}$.
(ii) For every $j \geq 1$ with $\mu(j)>0$, the shifted trees $\theta_{1} \mathcal{T}, \ldots, \theta_{j} \mathcal{T}$ are independent under the conditional probability $\mathbf{Q}_{\mu}\left(\cdot \mid k_{\emptyset}=j\right)$ and their conditional distribution is $\mathbf{Q}_{\mu}$.

A random tree with distribution $\mathbf{Q}_{\mu}$ is called a Galton-Watson tree with offspring distribution $\mu$, or in short a $\mu$-Galton-Watson tree.

Let $\mathcal{T}_{1}, \mathcal{T}_{2}, \ldots$ be a sequence of independent $\mu$-Galton-Watson trees. We can associate with this sequence a height process obtained by concatenating the height functions of each of the trees $\mathcal{T}_{1}, \mathcal{T}_{2}, \ldots$. More precisely, for every $k \geq 1$, we set
$H_{n}=H_{n-\left(\#\left(\mathcal{T}_{1}\right)+\cdots+\#\left(\mathcal{T}_{k-1}\right)\right)}\left(\mathcal{T}_{k}\right)$ if $\#\left(\mathcal{T}_{1}\right)+\cdots+\#\left(\mathcal{T}_{k-1}\right) \leq n<\#\left(\mathcal{T}_{1}\right)+\cdots+\#\left(\mathcal{T}_{k}\right)$.
The process $\left(H_{n}, n \geq 0\right)$ codes the sequence of trees.
Similarly, we define a contour process $\left(C_{t}, t \geq 0\right)$ coding the sequence of trees by concatenating the contour functions $\left(C_{t}\left(\mathcal{T}_{1}\right), t \in\left[0, \zeta\left(\mathcal{T}_{1}\right)+2\right]\right),\left(C_{t}\left(\mathcal{T}_{2}\right), t \in\left[0, \zeta\left(\mathcal{T}_{2}\right)+\right.\right.$ 2]), etc. Note that $C_{t}\left(\mathcal{T}_{n}\right)=0$ for $t \in\left[\zeta\left(\mathcal{T}_{n}\right), \zeta\left(\mathcal{T}_{n}\right)+2\right]$, and that we are concatenating the functions $\left(C_{t}\left(\mathcal{T}_{n}\right), t \in\left[0, \zeta\left(\mathcal{T}_{n}\right)+2\right]\right)$ rather than the functions $\left(C_{t}\left(\mathcal{T}_{n}\right), t \in\right.$ $\left.\left[0, \zeta\left(\mathcal{T}_{n}\right)\right]\right)$. This is a technical trick that will be useful in Chapter 2 below. We may also observe that the process obtained by concatenating the functions $\left(C_{t}\left(\mathcal{T}_{n}\right), t \in\right.$ $\left.\left[0, \zeta\left(\mathcal{T}_{n}\right)\right]\right)$ would not determine the sequence of trees.

There is a simple relation between the height process and the contour process: See Section 2.4 in Chapter 2 for more details.

Although the height process is not a Markov process, except in very particular cases, it turns out to be a simple functional of a Markov chain, which is even a random walk. The next lemma is taken from [32], but was obtained independently by other authors: See [7] and [4].

Lemma Let $\mathcal{T}_{1}, \mathcal{T}_{2}, \ldots$ be a sequence of independent $\mu$-Galton-Watson trees, and let $\left(H_{n}, n \geq 0\right)$ be the associated height process. There exists a random walk $V$ on $\mathbb{Z}$ with initial value $V_{0}=0$ and jump distribution $\nu(k)=\mu(k+1)$, for $k=-1,0,1,2, \ldots$, such that for every $n \geq 0$,

$$
\begin{equation*}
H_{n}=\operatorname{Card}\left\{k \in\{0,1, \ldots, n-1\}: V_{k}=\inf _{k \leq j \leq n} V_{j}\right\} \tag{1}
\end{equation*}
$$

A detailed proof of this lemma would be cumbersome, and we only explain the idea. By definition, $H_{n}$ is the generation of the individual visited at time $n$, for a particle that visits the different vertices of the sequence of trees one tree after another and in lexicographical order for each tree. Write $R_{n}$ for the quantity equal to the
number of younger brothers (younger means greater in the lexicographical order) of the individual visited at time $n$ plus the number of younger brothers of his father, plus the number of younger brothers of his grandfather etc. Then the random walk that appears in the lemma may be defined by

$$
V_{n}=R_{n}-(j-1) \quad \text { if } \#\left(\mathcal{T}_{1}\right)+\cdots+\#\left(\mathcal{T}_{j-1}\right) \leq n<\#\left(\mathcal{T}_{1}\right)+\cdots+\#\left(\mathcal{T}_{j}\right)
$$

To verify that $V$ is a random walk with jump distribution $\nu$, note that because of the lexicographical order of visits, we have at time $n$ no information on the fact that the individual visited at that time has children or not. If he has say $k \geq 1$ children, which occurs with probability $\mu(k)$, then the individual visited at time $n+1$ will be the first of these children, and our definitions give $R_{n+1}=R_{n}+(k-1)$ and $V_{n+1}=V_{n}+(k-1)$. On the other hand if he has no child, which occurs with probability $\mu(0)$, then the individual visited at time $n+1$ is the first of the brothers counted in the definition of $R_{n}$ (or the ancestor of the next tree if $R_{n}=0$ ) and we easily see that $V_{n+1}=V_{n}-1$. We thus get exactly the transition mechanism of the random walk with jump distribution $\nu$.

Let us finally explain formula (1). From our definition of $R_{n}$ and $V_{n}$, it is easy to see that the condition $n<\inf \left\{j>k: V_{j}<V_{k}\right\}$ holds iff the individual visited at time $n$ is a descendant of the individual visited at time $k$ (more precisely, $\inf \left\{j>k: V_{j}<V_{k}\right\}$ is the time of the first visit after $k$ of an individual that is not a descendant of individual $k$ ). Put in a different way, the condition $V_{k}=\inf _{k \leq j \leq n} V_{j}$ holds iff the individual visited at time $k$ is an ascendant of the individual visited at time $n$. It is now clear that the right-hand side of (1) just counts the number of ascendants of the individual visited at time $n$, that is the generation of this individual.

### 0.3 The continuous height process

To define the height process in a continuous setting, we use an analogue of the discrete formula (1). The role of the random walk $V$ in this formula is played by a Lévy process $X=\left(X_{t}, t \geq 0\right)$ without negative jumps. We assume that $X$ does not drift to $+\infty$ (this corresponds to the subcriticality of $\mu$ in the discrete setting), and that the paths of $X$ are of infinite variation a.s.: The latter assumption implies in particular that the process $X$ started at the origin will immediately hit both $(0, \infty)$ and $(-\infty, 0)$. The law of $X$ can be characterized by its Laplace functional $\psi$, which is the nonnegative function on $\mathbb{R}_{+}$defined by

$$
E\left[\exp \left(-\lambda X_{t}\right)\right]=\exp (t \psi(\lambda))
$$

By the Lévy-Khintchine formula and our special assumptions on $X$, the function $\psi$ has to be of the form

$$
\psi(\lambda)=\alpha \lambda+\beta \lambda^{2}+\int \pi(d r)\left(e^{-\lambda r}-1+\lambda r\right)
$$

where $\alpha, \beta \geq 0$ and $\pi$ is a $\sigma$-finite measure on $(0, \infty)$ such that $\int \pi(d r)\left(r \wedge r^{2}\right)<\infty$. We write

$$
S_{t}=\sup _{s \leq t} X_{s}, \quad I_{t}=\inf _{s \leq t} X_{s}
$$

By analogy with the discrete case, we would like to define $H_{t}$ as the "measure" of the set

$$
\begin{equation*}
\left\{s \leq t: X_{s}=\inf _{s \leq r \leq t} X_{r}\right\} \tag{2}
\end{equation*}
$$

However, under our assumptions on $X$, the Lebesgue measure of this set is always zero, and so we need to use some sort of local time. The key idea is to introduce for every fixed $t>0$ the time-reversed process

$$
\widehat{X}_{s}^{(t)}=X_{t}-X_{(t-s)-}, \quad 0 \leq s \leq t
$$

and its associated supremum

$$
\widehat{S}_{s}^{(t)}=\sup _{r \leq s} \widehat{X}_{r}^{(t)}
$$

We observe that via time-reversal $s \rightarrow t-s$, the set (2) corresponds to $\{s \leq t$ : $\left.\widehat{S}_{s}^{(t)}=\widehat{X}_{s}^{(t)}\right\}$. This leads to the rigorous definition of $H: H_{t}$ is defined as the local time at level 0 , at time $t$ of the process $\widehat{S}^{(t)}-\widehat{X}^{(t)}$. This definition makes sense because $\widehat{S}^{(t)}-\widehat{X}^{(t)}$ has the same law over $[0, t]$ as the so-called reflected process $S-X$ for which 0 is a regular point under our assumptions. Note that the normalization of local time has to be specified in some way: See Section 1.1. The process $\left(H_{t}, t \geq 0\right)$ is called the $\psi$-height process, or simply the height process.

Why is the $\psi$-height process $H$ an interesting object of study? In the same way as the discrete height process codes the genealogy of a sequence of independent GaltonWatson trees, we claim that the continuous height process represents the genealogical structure of continuous-state branching processes, which are the continuous analogues of Galton-Watson processes. This informal claim is at the heart of the developments of the present work. Perhaps the best justification for it can be found in the limit theorems of Chapter 2 that relate the discrete and continuous height processes (see Section 0.4 below). Another justification is the Ray-Knight theorem for the height process that will be discussed below.

The goal of Chapter 1 is to present a self-contained construction and to derive several new properties of the $\psi$-height process. Although there is some overlap with [32], our approach is different and involves new approximations. It is important to realize that $H_{t}$ is defined as the local time at time $t$ of a process which itself depends on $t$. For this reason, it is not clear whether the paths of $H$ have any regularity properties. Also $H$ is not Markov, except in the very special case where $X$ has no jumps. To circumvent these difficulties, we rely on the important tool of the exploration process: For every $t \geq 0$, we define a random measure $\rho_{t}$ on $\mathbb{R}_{+}$by setting

$$
\begin{equation*}
\left\langle\rho_{t}, f\right\rangle=\int_{[0, t]} d_{s} I_{t}^{s} f\left(H_{s}\right) \tag{3}
\end{equation*}
$$

where

$$
I_{t}^{s}=\inf _{s \leq r \leq t} X_{r}
$$

and the notation $d_{s} I_{t}^{s}$ refers to integration with respect to the nondecreasing function $s \rightarrow I_{t}^{s}$. The exploration process $\left(\rho_{t}, t \geq 0\right)$ is a Markov process with values in the space $M_{f}\left(\mathbb{R}_{+}\right)$of finite measures on $\mathbb{R}_{+}$. It was introduced and studied in [32], where its definition was motivated by a model of a LIFO queue (see [35] for some applications to queueing theory).

The exploration process has several interesting properties. In particular it is càdlàg (right-continuous with left limits) and it has an explicit invariant measure in terms of the subordinator with Laplace exponent $\psi(\lambda) / \lambda$ (see Proposition 1.2.5). Despite its apparently complicated definition, the exploration process is the crucial tool that makes it possible to answer most questions concerning the height process. A first illustration of this is the choice of a "good" lower-semicontinuous modification of $H_{t}$, which is obtained by considering for every $t \geq 0$ the supremum of the support of the
measure $\rho_{t}$ (beforehand, to make sense of the definition of $\rho_{t}$, one needs to use a first version of $H$ that can be defined by suitable approximations of local times).

An important feature of both the height process and the exploration process is the fact that both $H_{t}$ and $\rho_{t}$ depend only on the values of $X$, or of $X-I$, on the excursion interval of $X-I$ away from 0 that straddles $t$. For this reason, it is possible to define and to study both the height process and the exploration process under the excursion measure of $X-I$ away from 0 . This excursion measure, which is denoted by $N$, plays a major role throughout this work, and many results are more conveniently stated under $N$. Informally, the height process under $N$ codes exactly one continuous tree, in the same way as each excursion away from 0 of the discrete height process corresponds to one Galton-Watson tree in the sequence (cf Section 0.2).

As a typical application of the exploration process, we introduce and study the local times of the height process, which had not been considered in earlier work. These local times play an important role in the sequel, in particular in the applications to spatial branching processes. The local time of $H$ at level $a \geq 0$ and at time $t$ is denoted by $L_{t}^{a}$ and these local times can be defined through the approximation

$$
\lim _{\varepsilon \rightarrow 0} E\left[\sup _{s \leq t}\left|\varepsilon^{-1} \int_{0}^{s} 1_{\left\{a<H_{r}<a+\varepsilon\right\}} d r-L_{s}^{a}\right|\right]=0
$$

(Proposition 1.3.3). The proof of this approximation depends in a crucial way on properties of the exploration process derived in Section 1.3: Since $H$ is in general not Markovian nor a semimartingale, one cannot use the standard methods of construction of local time.

The Ray-Knight theorem for the height process states that if $T_{r}=\inf \{t \geq 0$ : $\left.X_{t}=-r\right\}$, for $r>0$, the process $\left(L_{T_{r}}^{a}, a \geq 0\right)$ is a continuous-state branching process with branching mechanism $\psi$ (in short a $\psi$-CSBP) started at $r$. Recall that the $\psi$ CSBP is the Markov process $\left(Y_{a}, a \geq 0\right)$ with values in $\mathbb{R}_{+}$whose transition kernels are characterized by their Laplace transform: For $\lambda>0$ and $b>a$,

$$
E\left[\exp -\lambda Y_{b} \mid Y_{a}\right]=\exp \left(-Y_{a} u_{b-a}(\lambda)\right)
$$

where $u_{t}(\lambda), t \geq 0$ is the unique nonnegative solution of the differential equation

$$
\frac{\partial u_{t}(\lambda)}{\partial t}=-\psi\left(u_{t}(\lambda)\right), \quad u_{0}(\lambda)=\lambda
$$

By analogy with the discrete setting, we can think of $L_{T_{r}}^{a}$ as "counting" the number of individuals at generation $a$ in a Poisson collection of continuous trees (those trees coded by the excursions of $X-I$ away from 0 before time $T_{r}$ ). The Ray-Knight theorem corresponds to the intuitive fact that the population at generation $a$ is a branching process.

The previous Ray-Knight theorem had already been derived in [32] although in a less precise form (local times of the height process had not been constructed). An important consequence of the Ray-Knight theorem, also derived in [32], is a criterion for the path continuity of $H: H$ has continuous sample paths iff

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d \lambda}{\psi(\lambda)}<\infty \tag{4}
\end{equation*}
$$

This condition is in fact necessary and sufficient for the a.s. extinction of the $\psi$-CSBP. If it does not hold, the paths of $H$ have a very wild behavior: The values of $H$ over any nontrivial interval $[s, t]$ contain a half-line $[a, \infty)$. On the other hand, (4) holds
if $\beta>0$, and in the stable case $\psi(\lambda)=c \lambda^{\gamma}, 1<\gamma \leq 2$ (the values $\gamma \in(0,1]$ are excluded by our assumptions).

In view of applications in Chapter 4, we derive precise information about the Hölder continuity of $H$. We show that if

$$
\gamma=\sup \left\{r \geq 0: \lim _{\lambda \rightarrow \infty} \lambda^{-r} \psi(\lambda)=+\infty\right\}
$$

then the height process $H$ is a.s. Hölder continuous with exponent $r$ for any $r \in$ $\left(0,1-\gamma^{-1}\right)$, and a.s. not Hölder continuous with exponent $r$ if $r>1-\gamma^{-1}$.

### 0.4 From discrete to continuous trees

Chapter 2 discusses limit theorems for rescaled Galton-Watson trees. These results demonstrate that the $\psi$-height process is the correct continuous analogue of the discrete height process coding Galton-Watson trees.

It is well known [27] that continuous-state branching processes are the only possible scaling limits of discrete-time Galton-Watson branching processes. One may then ask for finer limit theorems involving the genealogy. Precisely, starting from a sequence of rescaled Galton-Watson processes that converge in distribution towards a continuousstate branching process, can one say that the corresponding discrete Galton-Watson trees also converge, in some sense, towards a continuous genealogical structure? The results of Chapter 2 show that the answer is yes.

To be specific, consider a sequence ( $\mu_{p}$ ) of (sub)critical offspring distributions. For every $p \geq 1$, let $Y^{p}$ be a (discrete-time) Galton-Watson process with offspring distribution $\mu_{p}$ started at $Y_{0}^{p}=p$. Suppose that the processes $Y^{p}$ converge after rescaling towards a $\psi$-CSBP, where $\psi$ satisfies the conditions introduced in Section 0.3. Precisely, we assume that there is a sequence $\gamma_{p} \uparrow \infty$ such that

$$
\begin{equation*}
\left(p^{-1} Y_{\left[\gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(Y_{t}, t \geq 0\right) \tag{5}
\end{equation*}
$$

where $Y$ is a $\psi$-CSBP, and the symbol (d) indicates convergence in distribution in the Skorokhod space. Let $H^{p}$ be the discrete height process associated with $\mu_{p}$ in the sense of Section 0.2. Then Theorem 2.2.1 shows that

$$
\begin{equation*}
\left(\gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \xrightarrow[p \rightarrow \infty]{(\mathrm{fd})}\left(H_{t}, t \geq 0\right) \tag{6}
\end{equation*}
$$

where $H$ is the $\psi$-height process and (fd) indicates convergence of finite-dimensional marginals. A key ingredient of the proof is the observation due to Grimvall [21] that the convergence (5) implies the convergence in distribution (after suitable rescaling) of the random walks $V^{p}$ with jump distribution $\nu_{p}(k)=\mu_{p}(k+1), k=-1,0,1, \ldots$, towards the Lévy process with Laplace exponent $\psi$. The idea is then to pass to the limit in the formula for $H^{p}$ in terms of $V^{p}$, recalling that the $\psi$-height process is given by an analogous formula in terms of the Lévy process $X$. In the special case $\beta=0$ and under more restrictive assumptions, the convergence (6) had already appeared in [32].

In view of applications, the limiting result (6) is not satisfactory because the convergence of finite-dimensional marginals is too weak. In order to reinforce (6) to a functional convergence, it is necessary to assume some regularity of the paths of $H$. We assume that condition (4) ensuring the path continuity of $H$ holds (recall that if this condition does not hold, the paths of $H$ have a very wild behavior). Then, we can
prove (Theorem 2.3.1) that the convergence (6) holds in the sense of weak convergence on the Skorokhod space, provided that the following condition is satisfied: For every $\delta>0$,

$$
\begin{equation*}
\liminf _{p \rightarrow \infty} P\left[Y_{\left[\delta \gamma_{p}\right]}^{p}=0\right]>0 \tag{7}
\end{equation*}
$$

Roughly speaking this means that the rescaled Galton-Watson process $\left(p^{-1} Y_{\left[\gamma_{p} t\right]}^{p}\right)_{t \geq 0}$ may die out at a time of order 1, as its weak limit $Y$ does (recall that we are assuming (4)). The technical condition (7) is both necessary and sufficient for the reinforcement of (6) to a functional convergence. Simple examples show that this condition cannot be omitted in general.

However, in the important special case where $\mu_{p}=\mu$ for every $p$, we are able to show (Theorem 2.3.2) that the technical condition (7) is always satisfied . In that case, $\psi$ must be of the form $\psi(u)=c u^{\gamma}$ with $1<\gamma \leq 2$, so that obviously (4) also holds. Thus when $\mu_{p}=\mu$ for every $p$, no extra condition is needed to get a functional convergence.

In Section 2.4, we show that the functional convergence derived for rescaled discrete height processes can be stated as well in terms of the contour processes (cf Section 0.1). Let $C^{p}=\left(C_{t}^{p}, t \geq 0\right)$ be the contour process for a sequence of independent $\mu_{p}$-Galton-Watson trees. Under the assumptions that warrant the functional convergence in (6), Theorem 2.4.1 shows that we have also

$$
\left(p^{-1} C_{p \gamma_{p} t}^{p} t \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(H_{t / 2}, t \geq 0\right) .
$$

Thus scaling limits are the same for the discrete height process and for the contour process.

In the remaining part of Chapter 2, we give applications of (6) assuming that the functional convergence holds. In particular, rather than considering a sequence of $\mu_{p}$-Galton-Watson trees, we discuss the height process associated with a single tree conditioned to be large. Precisely, let $\tilde{H}^{p}$ be the height process for one $\mu_{p}$-GaltonWatson tree conditioned to non-extinction at generation $\left[\gamma_{p} T\right]$, for some fixed $T>0$. Then, Proposition 2.5.2 gives

$$
\left(\gamma_{p}^{-1} \tilde{H}_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \xrightarrow[p \rightarrow \infty]{(\mathrm{d})}\left(\tilde{H}_{t}, t \geq 0\right)
$$

where the limiting process is an excursion of the $\psi$-height process conditioned to hit level $T$. This is of course reminiscent of a result of Aldous [3] who proved that in the case of a critical offspring distribution $\mu$ with finite variance, the contour process of a $\mu$-Galton-Watson tree conditioned to have exactly $p$ vertices converges after a suitable rescaling towards a normalized Brownian excursion (see also [19] and [36] for related results including the convergence of the height process in Aldous' setting). Note that in Aldous' result, the conditioning becomes degenerate in the limit, since the "probability" that a Brownian excursion has length exactly one is zero. This makes it more difficult to derive this result from our approach, although it seems very related to our limit theorems. See however Duquesne [10] for an extension of Aldous' theorem to the stable case using the tools of the present work (a related result in the stable case was obtained by Kersting [26]).

The end of Chapter 2 is devoted to reduced trees. We consider again a single Galton-Watson tree conditioned to non-extinction at generation $\left[\gamma_{p} T\right]$. For every $k<\left[\gamma_{p} T\right]$, we denote by $Z_{k}^{(p),\left[\gamma_{p} T\right]}$ the number of vertices at generation $k$ that have
descendants at generation $\left[\gamma_{p} T\right]$. Under the assumptions and as a consequence of Proposition 2.5.2, we can prove that

$$
\left(Z_{\left[\gamma_{p} t\right]}^{(p),\left[\gamma_{p} T\right]}, 0 \leq t<T\right) \underset{p \rightarrow \infty}{(\mathrm{fd})}\left(Z_{t}^{T}, 0 \leq t<T\right)
$$

where the limit $Z^{T}$ has a simple definition in terms of $\tilde{H}: Z_{t}^{T}$ is the number of excursions of $\tilde{H}$ above level $t$ that hit level $T$. Thanks to the properties of the height process and the exploration process that have been derived in Chapter 1, it is possible to calculate the distribution of the time-inhomogeneous branching process $\left(Z_{t}^{T}, t \geq 0\right)$. This distribution is derived in Theorem 2.7.1. Of course in the stable case, corresponding to $\mu_{p}=\mu$ for every $p$, the distribution of $Z^{T}$ had been computed previously. See in particular Zubkov [50] and Fleischmann and Siegmund-Schultze [17].

### 0.5 Duality properties of the exploration process

In the applications developed in Chapters 3 and 4, a key role is played by the duality properties of the exploration process $\rho$. We first observe that formula (3) defining the exploration process can be rewritten in the following equivalent way

$$
\rho_{t}(d r)=\beta 1_{\left[0, H_{t}\right]}(r) d r+\sum_{s \leq t, X_{s-}<I_{t}^{s}}\left(I_{t}^{s}-X_{s-}\right) \delta_{H_{s}}(d r)
$$

where $\delta_{H_{s}}$ is the Dirac measure at $H_{s}$, and we recall that $I_{t}^{s}=\inf _{s \leq r \leq t} X_{r}$. We then define another measure $\eta_{t}$ by setting

$$
\eta_{t}(d r)=\beta 1_{\left[0, H_{t}\right]}(r) d r+\sum_{s \leq t, X_{s-}<I_{t}^{s}}\left(X_{s}-I_{t}^{s}\right) \delta_{H_{s}}(d r) .
$$

To motivate this definition, we may come back to the discrete setting of GaltonWatson trees. In that setting, the value at time $n$ of the discrete height process $H_{n}$ is the generation of the $n$-th visited vertex by a "particle" that visits vertices in lexicographical order one tree after another, and the analogue of $\rho_{t}$ gives for every $k \leq H_{n}$ the number of younger (i.e. coming later in the lexicographical order) brothers of the ancestor at generation $k$ of the $n$-th visited vertex. Then the analogue of $\eta_{t}$ gives for every $k \leq H_{n}$ the number of older brothers of the ancestor at generation $k$ of the $n$-th visited vertex.

It does not seem easy to study directly the Markovian properties or the regularity of paths of the process $\left(\eta_{t}, t \geq 0\right)$. The right point of view is to consider the pair $\left(\rho_{t}, \eta_{t}\right)$, which is easily seen to be a Markov process in $M_{f}\left(\mathbb{R}_{+}\right)^{2}$. The process $\left(\rho_{t}, \eta_{t}\right)$ has an invariant measure $\mathbb{M}$ determined in Proposition 3.1.3. The key result (Theorem 3.1.4) then states that the Markov processes $(\rho, \eta)$ and $(\eta, \rho)$ are in duality under $\mathbb{M}$. A consequence of this is the fact that $\left(\eta_{t}, t \geq 0\right)$ also has a càdlàg modification. More importantly, we obtain a crucial time-reversal property: Under the excursion measure $N$ of $X-I$, the processes $\left(\rho_{s}, \eta_{s} ; 0 \leq s \leq \sigma\right)$ and ( $\left.\eta_{(\sigma-s)-}, \rho_{(\sigma-s)-} ; 0 \leq s \leq \sigma\right)$ have the same distribution (here $\sigma$ stands for the duration of the excursion under $N$ ). This time-reversal property plays a major role in many subsequent calculations. It implies in particular that the law of $H$ under $N$ is invariant under time-reversal. This property is natural in the discrete setting, if we think of the contour process of a Galton-Watson tree, but not obvious in the continuous case.

### 0.6 Marginals of trees coded by the height process

Let us explain more precisely how an excursion of the $\psi$-height process codes a continuous branching structure. We consider first a deterministic continuous function $e: \mathbb{R}_{+} \longrightarrow \mathbb{R}_{+}$such that $e(t)>0$ iff $0<t<\sigma$, for some $\sigma=\sigma(e)>0$. For any $s, s^{\prime} \geq 0$, set

$$
m_{e}\left(s, s^{\prime}\right)=\inf _{s \wedge s^{\prime} \leq t \leq s \vee s^{\prime}} e(t)
$$

Then $e$ codes a continuous genealogical structure via the following simple prescriptions:
(i) To each $s \in[0, \sigma]$ corresponds a vertex at generation $e(s)$.
(ii) Vertex $s$ is an ancestor of vertex $s^{\prime}$ if $e(s)=m_{e}\left(s, s^{\prime}\right)$. In general, $m_{e}\left(s, s^{\prime}\right)$ is the generation of the last common ancestor to $s$ and $s^{\prime}$.
(iii) We put $d\left(s, s^{\prime}\right)=e(s)+e\left(s^{\prime}\right)-2 m_{e}\left(s, s^{\prime}\right)$ and identify $s$ and $s^{\prime}\left(s \sim s^{\prime}\right)$ if $d\left(s, s^{\prime}\right)=0$.

Formally, the tree coded by $e$ can be defined as the quotient set $[0, \sigma] / \sim$, equipped with the distance $d$ and the genealogical relation specified in (ii).

With these definitions, the line of ancestors of a vertex $s$ is isometric to the segment $[0, e(s)]$. If we pick two vertices $s$ and $s^{\prime}$, their lines of ancestors share a common part isometric to $\left[0, m_{e}\left(s, s^{\prime}\right)\right]$, and then become distinct. In general, if we consider $p$ instants $t_{1}, \ldots, t_{p}$ with $0 \leq t_{1} \leq \cdots \leq t_{p} \leq \sigma$, we can associate with these $p$ instants a genealogical tree $\theta\left(e, t_{1}, \ldots, t_{p}\right)$, which consists of a discrete rooted ordered tree with $p$ leaves, denoted by $\mathcal{T}\left(e, t_{1}, \ldots, t_{p}\right)$ and marks $h_{v}\left(e, t_{1}, \ldots, t_{p}\right) \geq 0$ for $v \in \mathcal{T}\left(e, t_{1}, \ldots, t_{p}\right)$, that correspond to the lifetimes of vertices in $\mathcal{T}\left(e, t_{1}, \ldots, t_{p}\right)$. See subsection 3.2.1 for a precise definition.

In the second part of Chapter 3, we use the duality results proved in the first part to calculate the distribution of the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{p}\right)$ under certain excursion laws of $H$ and random choices of the instants $\tau_{1}, \ldots, \tau_{p}$. We assume that the continuity condition (4) holds. We first consider Poissonnian marks with intensity $\lambda$, and the height process $H$ under the excursion measure $N$ of $X-I$. Let $\tau_{1}, \ldots, \tau_{M}$ be the marks that fall into the duration interval $[0, \sigma]$ of the excursion. Theorem 3.2.1 shows that under the probability measure $N(\cdot \mid M \geq 1)$, the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$ is distributed as the family tree of a continuous-time Galton-Watson process starting with one individual at time 0 and where

- lifetimes have an exponential distribution with parameter $\psi^{\prime}\left(\psi^{-1}(\lambda)\right)$;
- the offspring distribution is the law of the variable $\xi$ with generating function

$$
E\left[r^{\xi}\right]=r+\frac{\psi\left((1-r) \psi^{-1}(\lambda)\right)}{\psi^{-1}(\lambda) \psi^{\prime}\left(\psi^{-1}(\lambda)\right)}
$$

In the quadratic case, we get a critical binary branching $E\left[r^{\xi}\right]=\frac{1}{2}\left(1+r^{2}\right)$. The result in that case had been obtained by Hobson [22].

We finally specialize to the stable case $\psi(\lambda)=\lambda^{\gamma}, \gamma \in(1,2]$. By scaling arguments, we can then make sense of the law $N_{(1)}=N(\cdot \mid \sigma=1)$ of the normalized excursion of $H$. Using the case of Poissonnian marks, we compute explicitly the law of the tree $\theta\left(H, t_{1}, \ldots, t_{p}\right)$ under $N_{(1)}$, when $\left(t_{1}, \ldots, t_{p}\right)$ are chosen independently and uniformly over $[0,1]^{p}$. In the quadratic case $\psi(u)=u^{2}, H$ is under $N_{(1)}$ a normalized Brownian
excursion, and the corresponding tree is called the continuum random tree (see Aldous [1],[2], [3]). By analogy, in our more general case $\psi(u)=u^{\gamma}$, we may call the tree coded by $H$ under $N_{(1)}$ the stable continuum random tree. Our calculations give what Aldous calls the finite-dimensional marginals of the tree. In the case $\gamma=2$, these marginals were computed by Aldous (see also Le Gall [31] for a different approach closer to the present work). In that case, the discrete skeleton $\mathcal{T}\left(H, t_{1}, \ldots, t_{p}\right)$ is uniformly distributed over all binary rooted ordered trees with $k$ leaves. When $\gamma<2$, things become different as we can get nonbinary trees (the reason why we get only binary trees in the Brownian case is the fact that local minima of Brownian motion are distinct). Theorem 3.3 .3 shows in particular that if $\mathcal{T}$ is a tree with $p$ leaves such that $k_{u}(\mathcal{T}) \neq 1$ for every $u \in \mathcal{T}$ (this condition must be satisfied by our trees $\left.\mathcal{T}\left(e, t_{1}, \ldots, t_{p}\right)\right)$ then the probability that $\mathcal{T}\left(H, t_{1}, \ldots, t_{p}\right)=\mathcal{T}$ is

$$
\frac{p!}{(\gamma-1)(2 \gamma-1) \cdots((p-1) \gamma-1)} \prod_{v \in \mathcal{N}_{\mathcal{T}}} \frac{\left|(\gamma-1)(\gamma-2) \cdots\left(\gamma-k_{v}+1\right)\right|}{k_{v}!}
$$

where $\mathcal{N}_{\mathcal{T}}=\left\{v \in \mathcal{T}: k_{v}>0\right\}$ is the set of nodes of $\mathcal{T}$. It would be interesting to know whether this distribution on discrete trees has occurred in other settings.

### 0.7 The Lévy snake

Chapters $1-3$ explore the continuous genealogical structure coded by the $\psi$-height process $H$. In Chapter 4, we examine the probabilistic objects obtained by combining this branching structure with a spatial motion given by a càdlàg Markov process $\xi$ with state space $E$. Informally, "individuals" do not only reproduce themselves, but they also move in space independently according to the law of $\xi$. The $(\xi, \psi)$ superprocess is then a Markov process taking values in the space of finite measures on $E$, whose value at time $t$ is a random measure putting mass on the set of positions of "individuals" alive at time $t$. Note that the previous description is very informal since in the continuous branching setting there are no individual particles but rather a continuum of infinitesimal particles. Recent accounts of the theory of superprocesses can be found in Dynkin [13], Etheridge [15] and Perkins [40].

Our coding of the genealogy by the height process leads to introducing a Markov process whose values will give the historical paths followed by the "individuals" in the population. This a generalization of the Brownian snake introduced in [28] and studied in particular in [31]. To give a precise definition, fix a starting point $x \in E$, consider the $\psi$-height process $\left(H_{s}, s \geq 0\right)$ and recall the notation $m_{H}\left(s, s^{\prime}\right)=\inf _{\left[s, s^{\prime}\right]} H_{r}$ for $s \leq s^{\prime}$. We assume that the continuity condition (4) holds. Then conditionally on $\left(H_{s}, s \geq 0\right)$ we consider a time-inhomogeneous Markov process ( $W_{s}, s \geq 0$ ) whose distribution is described as follows:

- For every $s \geq 0, W_{s}=\left(W_{s}(t), 0 \leq t<H_{s}\right)$ is a path of $\xi$ started at $x$ and with finite lifetime $H_{s}$.
- If we consider two instants $s$ and $s^{\prime}$, the corresponding paths $W_{s}$ and $W_{s^{\prime}}$ are the same up to time $m_{H}\left(s, s^{\prime}\right)$ and then behave independently.

The latter property is consistent with the fact that in our coding of the genealogy, vertices attached to $s$ and $s^{\prime}$ have the same ancestors up to generation $m_{H}\left(s, s^{\prime}\right)$. See Section 4.1 for a more precise definition.

The pair $\left(\rho_{s}, W_{s}\right)$ is then a Markov process with values in the product space $M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}$, where $\mathcal{W}$ stands for the set of all finite càdlàg paths in $E$. This process
is called the Lévy snake (with initial point $x$ ). It was introduced and studied in [33], where a form of its connection with superprocesses was established. Chapter 4 gives much more detailed information about its properties. In particular, we prove the strong Markov property of the Lévy snake (Theorem 4.1.2), which plays a crucial role in several applications.

We also use the local times of the height process to give a nicer form of the connection with superprocesses. Write $\widehat{W}_{s}$ for the left limit of $W_{s}$ at its lifetime $H_{s}$ (which exists a.s. for each fixed $s$ ), and recall the notation $T_{r}=\inf \left\{t \geq 0: X_{t}=-r\right\}$. For every $t \geq 0$, we can define a random measure $Z_{t}$ on $E$ by setting

$$
\left\langle Z_{t}, f\right\rangle=\int_{0}^{T_{r}} d_{s} L_{s}^{t} f\left(\widehat{W}_{s}\right)
$$

Then $\left(Z_{t}, t \geq 0\right)$ is a $(\xi, \psi)$-superprocess with initial value $r \delta_{x}$. This statement is in fact a special case of Theorem 4.2 .1 which constructs a $(\xi, \psi)$-superprocess with an arbitrary initial value. For this more general statement, it is necessary to use excursion measures of the Lévy snake: Under the excursion measure $\mathbb{N}_{x}$, the process ( $\rho_{s}, s \geq 0$ ) is distributed according to its excursion measure $N$, and ( $W_{s}, s \geq 0$ ) is constructed by the procedure explained above, taking $x$ for initial point.

As a second application, we use local time techniques to construct exit measures from an open set and to establish the integral equation satisfied by the Laplace functional of exit measures (Theorem 4.3.3). Recall that exit measures of superprocesses play a fundamental role in the connections with partial differential equations studied recently by Dynkin and Kuznetsov (a detailed account of these connections can be found in the book [13]).

We then study the continuity of the path-valued process $W_{s}$ with respect to the uniform topology on paths. This question is closely related to the compact support property for superprocesses. In the case when $\xi$ is Brownian motion in $\mathbb{R}^{d}$, Theorem 4.5.2 shows that the condition

$$
\int_{1}^{\infty}\left(\int_{0}^{t} \psi(u) d u\right)^{-1 / 2} d t<\infty
$$

is necessary and sufficient for $W_{t}$ to be continuous with respect to the uniform topology on paths. The proof relies on connections of the exit measure with partial differential equations and earlier work of Sheu [45], who was interested in the compact support property for superprocesses. In the case of a general spatial motion, assuming only that $\xi$ has Hölder continuous paths, we use the continuity properties of $H$ derived in Chapter 1 to give simple sufficient conditions ensuring that the same conclusion holds.

Although we do not develop such applications in the present work, we expect that the Lévy snake will be a powerful tool to study connections with partial differential equations, in the spirit of [30], as well as path properties of superprocesses (see [34] for a typical application of the Brownian snake to super-Brownian motion).

In the last two sections of Chapter 4, we compute certain explicit distributions related to the Lévy snake and the $(\xi, \psi)$-superprocess, under the excursion measures $\mathbb{N}_{x}$. We assume that the path-valued process $W_{s}$ is continuous with respect to the uniform topology on paths, and then the value $W_{s}\left(H_{s}\right)$ can be defined as a left limit at the lifetime, simultaneously for all $s \geq 0$. If $D$ is an open set in $E$ such that $x \in D$, we consider the first exit time

$$
T_{D}=\inf \left\{s \geq 0: \tau\left(W_{s}\right)<\infty\right\}
$$

where $\tau\left(W_{s}\right)=\inf \left\{t \in\left[0, H_{s}\right]: W_{s}(t) \notin D\right\}$. Write $u(y)=\mathbb{N}_{y}\left(T_{D}<\infty\right)<\infty$ for every $y \in D$. Then the distribution of $W_{T_{D}}$ under $\mathbb{N}_{x}\left(\cdot \cap\left\{T_{D}<\infty\right\}\right)$ is characterized by the function $u$ and the distribution $\Pi_{x}$ of $\xi$ started at $x$ via the formula: For every $a \geq 0$,

$$
\begin{aligned}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} 1_{\left\{a<H_{T_{D}}\right\}} F\left(W_{T_{D}}(t), 0 \leq t \leq a\right)\right) \\
& \quad=\Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) \exp \left(-\int_{0}^{a} \tilde{\psi}\left(u\left(\xi_{r}\right)\right) d r\right)\right]
\end{aligned}
$$

where $\tau$ stands for the first exit time of $\xi$ from $D$, and $\tilde{\psi}(r)=\psi(r) / r$. Theorem 4.6.2 gives more generally the law of the pair $\left(W_{T_{D}}, \rho_{T_{D}}\right)$ under $\mathbb{N}_{x}\left(\cdot \cap\left\{T_{D}<\infty\right\}\right)$. In the special case when $\xi$ is Brownian motion in $\mathbb{R}^{d}$, the function $u$ can be identified as the maximal nonnegative solution of $\frac{1}{2} \Delta u=\psi(u)$ in $D$, and the law of $W_{T_{D}}$ is that of a Brownian motion with drift $\nabla u / u$ up to its exit time from $D$. This considerably extends a result of [29] proved in the quadratic branching case by a very different method.

The last section of Chapter 4 investigates reduced spatial trees, again under the assumption that the path-valued process $W_{s}$ is continuous with respect to the uniform topology on paths. We consider a spatial open set $D$ with $x \in D$, and the Lévy snake under its excursion measure $\mathbb{N}_{x}$ (in the superprocess setting this means that we are looking at all historical paths corresponding to one given ancestor at time 0 ). We condition on the event that $\left\{T_{D}<\infty\right\}$, that is one at least of the paths $W_{s}$ exits $D$, and we want to describe the spatial structure of all the paths that exit $D$, up to their respective exit times. This is an analogue (and in fact a generalization) of the reduced tree problem studied in Chapter 2. In the spatial situation, all paths $W_{s}$ that exit $D$ will be the same up to a certain time $m_{D}$ at which there is a branching point with finitely many branches, each corresponding to an excursion of the height process $H$ above level $m_{D}$ during which the Lévy snake exits $D$. In each such excursion the paths $W_{s}$ that exit $D$ will be the same up to a level strictly greater than $m_{D}$, at which there is another branching point, and so on.

To get a full description of the reduced spatial tree, one only needs to compute the joint distribution of the path $W_{0}^{D}=W_{T_{D}}\left(. \wedge m_{D}\right)$, that is the common part to all paths that do exit $D$, and the number $N_{D}$ of branches at the first branching point. Indeed, conditionally on the pair $\left(W_{0}^{D}, N_{D}\right)$, the "subtrees of paths" that originate from the first branching point will be independent and distributed according to the full reduced tree with initial point $\widehat{W}_{0}^{D}=W_{0}^{D}\left(m_{D}\right)$ (see Theorem 4.7.2 for more precise statements). Theorem 4.7.2 gives explicit formulas for the joint distribution of $\left(W_{0}^{D}, N_{D}\right)$, again in terms of the function $u(y)=\mathbb{N}_{y}\left(T_{D}<\infty\right)<\infty$. Precisely, the law of the "first branch" $W_{0}^{D}$ is given by

$$
\begin{aligned}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} F\left(W_{0}^{D}\right)\right) \\
& =\int_{0}^{\infty} d b \Pi_{x}\left[1_{\{b<\tau\}} u\left(\xi_{b}\right) \theta\left(u\left(\xi_{b}\right)\right) \exp \left(-\int_{0}^{b} \psi^{\prime}\left(u\left(\xi_{r}\right)\right) d r\right) F\left(\xi_{r}, 0 \leq r \leq b\right)\right]
\end{aligned}
$$

where $\theta(r)=\psi^{\prime}(r)-\tilde{\psi}(r)$. Furthermore the conditional distribution of $N_{D}$ given $W_{0}^{D}$ depends only on the branching point $\widehat{W}_{0}^{D}$ and is given by

$$
\mathbb{N}_{x}\left[r^{N_{D}} \mid T_{D}<\infty, W_{0}^{D}\right]=r \frac{\psi^{\prime}(U)-\gamma_{\psi}(U,(1-r) U)}{\psi^{\prime}(U)-\gamma_{\psi}(U, 0)}, \quad 0 \leq r \leq 1
$$

where $U=u\left(\widehat{W}_{0}^{D}\right)$ and $\gamma_{\psi}(a, b)=\frac{\psi(a)-\psi(b)}{a-b}$. In the stable case $\psi(u)=u^{\gamma}$, the
variable $N_{D}$ is independent of $W_{0}^{D}$ and its generating function is $(\gamma-1)^{-1}\left((1-r)^{\gamma}-\right.$ $1+\gamma r)$.

In the special case $\psi(u)=u^{2}$, the previous result was used by Mselati [38] to investigate absolute continuity properties of the law of the exit measure of superBrownian motion, in view of applications to the semilinear partial differential equation $\Delta u=u^{2}$. Again in the quadratic branching case, our description of the reduced spatial tree is reminiscent of the recent work of Salisbury and Verzani [43],[44] who study exit measures of super-Brownian motion conditioned to hit a number of specified points on the boundary. This conditioning leads to a spatial tree described by a branching particle backbone process with immigration of mass along the paths of the particles.

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## Chapter 1

## The height process

### 1.1 Preliminaries on Lévy processes

### 1.1.1 Basic assumptions

In this section, we introduce the class of Lévy processes that will be relevant to our study and we record some of their basic properties. For almost all facts about Lévy processes that we need, we refer to the recent book of Bertoin [5] (especially Chapter VII).

We consider a Lévy process $X$ on the real line. It will be convenient to assume that $X$ is the canonical process on the Skorokhod space $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ of càdlàg (rightcontinuous with left limits) real-valued paths. Unless otherwise noted, the underlying probability measure $P$ is the law of the process started at 0 . We denote by ( $\mathcal{G}_{t}, t \in$ $[0, \infty]$ ) the canonical filtration, completed as usual by the class of all $P$-negligible sets of $\mathcal{G}_{\infty}$.

We assume that the following three properties hold a.s.:
(H1) $X$ has no negative jumps.
(H2) $X$ does not drift to $+\infty$.
(H3) The paths of $X$ are of infinite variation.
Thanks to (H1), the "Laplace transform" $E\left[\exp -\lambda X_{t}\right]$ is well defined for every $\lambda \geq 0$ and $t \geq 0$, and can be written as

$$
E\left[\exp -\lambda X_{t}\right]=\exp (t \psi(\lambda)),
$$

with a function $\psi$ of the form

$$
\psi(\lambda)=\alpha_{0} \lambda+\beta \lambda^{2}+\int_{(0, \infty)} \pi(d r)\left(e^{-\lambda r}-1+1_{\{r<1\}} \lambda r\right)
$$

where $\alpha_{0} \in \mathbb{R}, \beta \geq 0$ and the Lévy measure $\pi$ is a Radon measure on $(0, \infty)$ such that $\int_{(0, \infty)}\left(1 \wedge r^{2}\right) \pi(d r)<\infty$.

Assumption (H2) then holds iff $X$ has first moments and $E\left[X_{1}\right] \leq 0$. The first moment assumption is equivalent to saying that $\pi$ satisfies the stronger integrability condition

$$
\int_{(0, \infty)}\left(r \wedge r^{2}\right) \pi(d r)<\infty
$$

Then $\psi$ can be written in the form

$$
\begin{equation*}
\psi(\lambda)=\alpha \lambda+\beta \lambda^{2}+\int_{(0, \infty)} \pi(d r)\left(e^{-\lambda r}-1+\lambda r\right) \tag{1.1}
\end{equation*}
$$

Note that $\psi$ is then convex and that we have $E\left[X_{t}\right]=-t \psi^{\prime}(0)=-t \alpha$. The condition $E\left[X_{1}\right] \leq 0$ thus holds iff $\alpha \geq 0$. The process $X$ is recurrent or drifts to $-\infty$ according as $\alpha=0$ or $\alpha>0$.

Finally, according to [5] (Corollary VII.5), assumption (H3) holds iff at least one of the following two conditions is satisfied: $\beta>0$, or

$$
\int_{(0,1)} r \pi(d r)=\infty
$$

Summarizing, we assume that $X$ is a Lévy process with no negative jumps, whose Laplace exponent $\psi$ has the form (1.1), where $\alpha \geq 0, \beta \geq 0$ and $\pi$ is a $\sigma$-finite measure on $(0, \infty)$ such that $\int\left(r \wedge r^{2}\right) \pi(d r)<\infty$, and we exclude the case where both $\beta=0$ and $\int_{(0,1)} r \pi(d r)<\infty$.
Remark. Only assumption (H1) is crucial to the connections with branching processes that are presented in this work. Assumption (H2) means that we restrict our attention to the critical or subcritical case. We impose assumption (H3) in order to concentrate on the most interesting cases: A simpler parallel theory can be developed in the finite variation case, see Section 3 of [32].

We will use the notation $T_{y}=\inf \left\{t \geq 0: X_{t}=-y\right\}$ for $y \in \mathbb{R}$. By convention $\inf \emptyset=+\infty$.

Under our assumptions, the point 0 is regular for $(0, \infty)$ and for $(-\infty, 0)$, meaning that $\inf \left\{t>0: X_{t}>0\right\}=0$ and $\inf \left\{t>0: X_{t}<0\right\}=0$ a.s. (see [5], Theorem VII. 1 and Corollary VII.5). We sometimes use this property in connection with the so-called duality property: For every $t>0$, define a process $\widehat{X}^{(t)}=\left(\widehat{X}_{s}^{(t)}, 0 \leq s \leq t\right)$ by setting

$$
\widehat{X}_{s}^{(t)}=X_{t}-X_{(t-s)-}, \quad \text { if } 0 \leq s<t
$$

and $\widehat{X}_{t}^{(t)}=X_{t}$. Then $\left(\widehat{X}_{s}^{(t)}, 0 \leq s \leq t\right)$ has the same law as $\left(X_{s}, 0 \leq s \leq t\right)$.
Let $S$ and $I$ be respectively the supremum and the infimum process of $X$, defined by

$$
S_{t}=\sup _{s \leq t} X_{s}, \quad I_{t}=\inf _{s \leq t} X_{s}
$$

If we combine the duality property with the regularity of 0 for both $(0, \infty)$ and $(-\infty, 0)$, we easily get that the set

$$
\left\{s>0: X_{s-}=I_{s} \text { or } X_{s-}=S_{s-}\right\}
$$

almost surely does not intersect $\left\{s \geq 0: \Delta X_{s} \neq 0\right\}$. This property will be used implicitly in what follows.

### 1.1.2 Local times at the maximum and the minimum

Both processes $X-S$ and $X-I$ are strong Markov processes, and the results recalled at the end of the previous subsection imply that the point 0 is regular for itself with respect to each of these two Markov processes. We can thus define the corresponding

Markovian local times and excursion measures, which both play a fundamental role in this work.

Consider first $X-S$. We denote by $L=\left(L_{t}, t \geq 0\right)$ a local time at 0 for $X-S$. Observe that $L$ is only defined up to a positive multiplicative constant, that will be specified later. Let $N^{*}$ be the associated excursion measure, which is a $\sigma$-finite measure on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. It will be important for our purposes to keep track of the final jump under $N^{*}$. This can be achieved by the following construction. Let $\left(a_{j}, b_{j}\right), j \in$ $J$ be the excursion intervals of $X-S$ away from 0 . In the transient case ( $\alpha>0$ ), there is exactly one value $j \in J$ such that $b_{j}=+\infty$. For every $j \in J$ let $\omega^{j} \in \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ be defined by

$$
\omega^{j}(s)=X_{\left(a_{j}+s\right) \wedge b_{j}}-X_{a_{j}}, \quad s \geq 0 .
$$

Then the point measure

$$
\sum_{j \in J} \delta_{\left(L_{a_{j}}, \omega^{j}\right)}
$$

is distributed as $1_{\{l \leq \eta\}} \mathcal{N}(d l d \omega)$, where $\mathcal{N}$ denotes a Poisson point measure with intensity $d l N^{*}(d \omega)$, and $\eta=\inf \{l: \mathcal{N}([0, l] \times\{\sigma=+\infty\}) \geq 1\}$, if

$$
\sigma(\omega)=\inf \{t>0: \omega(r)=\omega(t) \text { for every } r \geq t\}
$$

stands for the duration of the excursion $\omega$. This statement characterizes the excursion measure $N^{*}$, up to the multiplicative constant already mentioned. Note that $X_{0}=0$ and $X_{t}=X_{\sigma} \geq 0$ for $t \geq \sigma, N^{*}$ a.e.

Consider then $X-I$. It is easy to verify that the continuous increasing process $-I$ is a local time at 0 for the Markov process $X-I$. We will denote by $N$ the associated excursion measure, which can be characterized in a way similar to $N^{*}$ (with the difference that we have always $-I_{\infty}=+\infty$ a.s., in contrast to the property $L_{\infty}<\infty$ a.s. in the transient case). We already noticed that excursions of $X-I$ cannot start with a jump. Hence, $X_{0}=0, N$ a.e. It is also clear from our assumptions on $X$ that $\sigma<\infty, X_{t}>0$ for every $t \in(0, \sigma)$ and $X_{\sigma-}=0, N$ a.e.

We will now specify the normalization of $N^{*}$, or equivalently of $L$. Let $m$ denote Lebesgue measure on $\mathbb{R}$.

Lemma 1.1.1 We can fix the normalization of $L$, or equivalently of $N^{*}$, so that, for every Borel subset $B$ of $(-\infty, 0)$,

$$
\begin{equation*}
N^{*}\left(\int_{0}^{\sigma} d s 1_{B}\left(X_{s}\right)\right)=m(B) . \tag{1.2}
\end{equation*}
$$

Proof. For every $x \in \mathbb{R}$, write $P_{x}$ for the law of the Lévy process started at $x$. Also set $\tau=\inf \left\{s \geq 0: X_{s} \geq 0\right\}$ and recall that $\left(X_{t}, t>0\right)$ is Markovian under $N^{*}$ with the transition kernels of the underlying Lévy process stopped when hitting $[0, \infty)$. Thanks to this observation, it is enough to prove that, for every $\varepsilon>0$, there exists a constant $c(\varepsilon)$ such that for every Borel subset $B$ of $(-\infty,-\varepsilon)$,

$$
E_{-\varepsilon}\left[\int_{0}^{\tau} d s 1_{B}\left(X_{s}\right)\right]=c(\varepsilon) m(B) .
$$

Consider first the transient case. By applying the strong Markov property at hitting times of negative values, it is easy to verify that the measure on $(-\infty,-\varepsilon)$ defined by

$$
B \longrightarrow E_{-\varepsilon}\left[\int_{0}^{\infty} d s 1_{B}\left(X_{s}\right)\right]
$$

must be a multiple of Lebesgue measure. However, writing $T_{0}^{\varepsilon}=0, T_{1}^{\varepsilon}, \ldots, T_{n}^{\varepsilon}$, etc. for the successive visits of $-\varepsilon$ via $[0, \infty)$, we have

$$
E_{-\varepsilon}\left[\int_{0}^{\infty} d s 1_{B}\left(X_{s}\right)\right]=\sum_{i=0}^{\infty} E_{-\varepsilon}\left[1_{\left\{T_{i}^{\varepsilon}<\infty\right\}} \int_{T_{i}^{\varepsilon}}^{T_{i+1}^{\varepsilon}} d s 1_{B}\left(X_{s}\right)\right]=\frac{E_{-\varepsilon}\left[\int_{0}^{\tau} d s 1_{B}\left(X_{s}\right)\right]}{P_{-\varepsilon}[\tau=\infty]}
$$

The desired result follows.
In the recurrent case, the ergodic theorem gives

$$
\frac{1}{n} \int_{0}^{T_{n}^{\varepsilon}} d s 1_{B}\left(X_{s}\right) \underset{n \rightarrow \infty}{\text { a.s. }} E_{-\varepsilon}\left[\int_{0}^{\tau} d s 1_{B}\left(X_{s}\right)\right]
$$

whereas the Chacon-Ornstein ergodic theorem implies

$$
\frac{\int_{0}^{T_{n}^{\varepsilon}} d s 1_{B}\left(X_{s}\right)}{\int_{0}^{T_{n}^{\varepsilon}} d s 1_{(-2 \varepsilon,-\varepsilon)}\left(X_{s}\right)} \xrightarrow[n \rightarrow \infty]{\text { a.s. }} \frac{m(B)}{\varepsilon}
$$

The conclusion easily follows.
In what follows we always assume that the normalization of $L$ or of $N^{*}$ is fixed as in Lemma 1.1.1.

Let $L^{-1}(t)=\inf \left\{s, L_{s}>t\right\}$. By convention, $X_{L^{-1}(t)}=+\infty$ if $t \geq L_{\infty}$. The process $\left(X_{L^{-1}(t)}, t \geq 0\right)$ is a subordinator (the so-called ladder height process) killed at an independent exponential time in the transient case.

Lemma 1.1.2 For every $\lambda>0$,

$$
E\left[\exp -\lambda X_{L^{-1}(t)}\right]=\exp (-t \tilde{\psi}(\lambda))
$$

where

$$
\tilde{\psi}(\lambda)=\frac{\psi(\lambda)}{\lambda}=\alpha+\beta \lambda+\int_{0}^{\infty}\left(1-e^{-\lambda r}\right) \pi([r, \infty)) d r
$$

Proof. By a classical result of fluctuation theory (see e.g. [6] Corollary p.724), we have

$$
E\left[\exp -\lambda X_{L^{-1}(t)}\right]=\exp (-c t \tilde{\psi}(\lambda))
$$

where $c$ is a positive constant. We have to verify that $c=1$ under our normalization.
Suppose first that $\pi \neq 0$. Then notice that the Lévy measure $c \pi([r, \infty)) d r$ of $X_{L^{-1}(t)}$ is the "law" of $X_{\sigma}$ under $N^{*}\left(\cdot \cap\left\{X_{\sigma}>0\right\}\right)$. However, for any nonnegative measurable function $f$ on $[0, \infty)^{2}$, we get by a predictable projection

$$
\begin{aligned}
N^{*}\left(f\left(\Delta X_{\sigma}, X_{\sigma}\right) 1_{\left\{X_{\sigma}>0\right\}}\right) & =N^{*}\left(\sum_{0<s \leq \sigma} f\left(\Delta X_{s}, X_{s}\right) 1_{\left\{X_{s}>0\right\}}\right) \\
& =N^{*}\left(\int_{0}^{\sigma} d s \int \pi(d x) f\left(x, X_{s-}+x\right) 1_{\left\{X_{s-}+x>0\right\}}\right) \\
& =\int_{-\infty}^{0} d y \int \pi(d x) f(x, y+x) 1_{\{y+x>0\}}
\end{aligned}
$$

using Lemma 1.1.1 in the last equality. It follows that

$$
\begin{equation*}
N^{*}\left(f\left(\Delta X_{\sigma}, X_{\sigma}\right) 1_{\left\{X_{\sigma}>0\right\}}\right)=\int \pi(d x) \int_{0}^{x} d z f(x, z) \tag{1.3}
\end{equation*}
$$

and we get $c=1$ by comparing with the Lévy measure of $X_{L^{-1}(t)}$.
In the case $\pi=0, X$ is a scaled linear Brownian motion with drift, and the same conclusion follows from direct computations.

Note that we have in particular $P\left[L^{-1}(t)<\infty\right]=e^{-\alpha t}$, which shows that $L_{\infty}$ has an exponential distribution with parameter $\alpha$ in the transient case.

When $\beta>0$, we can get a simple expression for $L_{t}$. From well-known results on subordinators, we have a.s. for every $u \geq 0$,

$$
m\left(\left\{X_{L^{-1}(t)} ; t \leq u, L^{-1}(t)<\infty\right\}\right)=\beta\left(u \wedge L_{\infty}\right)
$$

Since the sets $\left\{X_{L^{-1}(t)} ; t \leq u, L^{-1}(t)<\infty\right\}$ and $\left\{S_{r} ; r \leq L^{-1}(u)\right\}$ coincide except possibly for a countable set, we have also

$$
\begin{equation*}
m\left(\left\{S_{r} ; r \leq t\right\}\right)=\beta L_{t} \tag{1.4}
\end{equation*}
$$

for every $t \geq 0$ a.s.
The next lemma provides a useful approximation of the local time $L_{t}$.
Lemma 1.1.3 For every $x>0$,

$$
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \int_{0}^{L^{-1}(x)} 1_{\left\{S_{s}-X_{s}<\varepsilon\right\}} d s=x \wedge L_{\infty}
$$

in the $L^{2}$-norm. Consequently, for every $t \geq 0$,

$$
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \int_{0}^{t} 1_{\left\{S_{s}-X_{s}<\varepsilon\right\}} d s=L_{t}
$$

in probability.
Proof. It is enough to prove the first assertion. Let $\mathcal{N}$ be as previously a Poisson point measure on $\mathbb{R}_{+} \times \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ with intensity $d l N^{*}(d \omega)$, and $\eta=\inf \{l: \mathcal{N}([0, l] \times\{\sigma=$ $+\infty\}) \geq 1\}$. For every $x>0$ set

$$
J_{\varepsilon}(x)=\frac{1}{\varepsilon} \int \mathcal{N}(d l d \omega) 1_{\{l \leq x\}} \int_{0}^{\sigma(\omega)} 1_{(-\varepsilon, 0]}(\omega(s)) d s
$$

Then,

$$
E\left[J_{\varepsilon}(x)\right]=\frac{x}{\varepsilon} N^{*}\left(\int_{0}^{\sigma} 1_{(-\varepsilon, 0]}\left(X_{s}\right) d s\right)=x
$$

by (1.2). Furthermore,

$$
E\left[J_{\varepsilon}(x)^{2}\right]=\left(E\left[J_{\varepsilon}(x)\right]\right)^{2}+x \varepsilon^{-2} N^{*}\left(\left(\int_{0}^{\sigma} 1_{(-\varepsilon, 0]}\left(X_{s}\right) d s\right)^{2}\right)
$$

and

$$
\begin{aligned}
N^{*}\left(\left(\int_{0}^{\sigma} 1_{(-\varepsilon, 0]}\left(X_{s}\right) d s\right)^{2}\right) & =2 N^{*}\left(\int_{0 \leq s \leq t \leq \sigma} 1_{(-\varepsilon, 0]}\left(X_{s}\right) 1_{(-\varepsilon, 0]}\left(X_{t}\right) d s d t\right) \\
& =2 N^{*}\left(\int_{0}^{\sigma} d s 1_{(-\varepsilon, 0]}\left(X_{s}\right) E_{X_{s}}\left[\int_{0}^{\tau} d t 1_{(-\varepsilon, 0]}\left(X_{t}\right)\right]\right) \\
& \leq 2 \varepsilon \sup _{0 \geq y>-\varepsilon} E_{y}\left[\int_{0}^{\tau} d t 1_{(-\varepsilon, 0]}\left(X_{t}\right)\right],
\end{aligned}
$$

using the same notation $\tau=\inf \left\{t \geq 0: X_{t} \geq 0\right\}$ as previously. We then claim that

$$
\begin{equation*}
\sup _{0 \geq y>-\varepsilon} E_{y}\left[\int_{0}^{\tau} d t 1_{(-\varepsilon, 0]}\left(X_{t}\right)\right]=o(\varepsilon) \tag{1.5}
\end{equation*}
$$

as $\varepsilon \rightarrow 0$. Indeed, by applying the strong Markov property at $T_{y}$, we have for $y>0$,

$$
N^{*}\left(T_{y}<\infty\right) E_{-y}\left[\int_{0}^{\tau} d t 1_{(-\varepsilon, 0]}\left(X_{t}\right)\right] \leq N^{*}\left(\int_{0}^{\sigma} d t 1_{(-\varepsilon, 0]}\left(X_{t}\right)\right)=\varepsilon
$$

and the claim follows since $N^{*}\left[T_{y}<\infty\right] \uparrow+\infty$ as $y \downarrow 0$. From (1.5) and the preceding calculations, we get

$$
\lim _{\varepsilon \rightarrow 0} E\left[\left(J_{\varepsilon}(x)-x\right)^{2}\right]=0 .
$$

By Doob's inequality (or a monotonicity argument), we have also

$$
\lim _{\varepsilon \rightarrow 0} E\left[\sup _{0 \leq z \leq x}\left(J_{\varepsilon}(z)-z\right)^{2}\right]=0
$$

The lemma now follows, since the pair

$$
\left(\frac{1}{\varepsilon} \int_{0}^{L^{-1}(x)} 1_{\left\{S_{s}-X_{s}<\varepsilon\right\}} d s, L_{\infty}\right)
$$

has the same distribution as $\left(J_{\varepsilon}(x \wedge \eta), \eta\right)$.
As a consequence of Lemma 1.1.3, we may choose a sequence $\left(\varepsilon_{k}, k=1,2, \ldots\right)$ of positive real numbers decreasing to 0 , such that

$$
\begin{equation*}
L_{t}=\lim _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{0}^{t} 1_{\left\{S_{s}-X_{s}<\varepsilon_{k}\right\}} d s, \quad P \text { a.s. } \tag{1.6}
\end{equation*}
$$

Using monotonicity arguments and a diagonal subsequence, we may and will assume that the previous convergence holds simultaneously for every $t \geq 0$ outside a single set of zero probability. In particular, if we set for $\omega \in \mathbb{D}([0, t], \mathbb{R})$,

$$
\Phi_{t}(\omega)=\liminf _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{0}^{t} 1_{\left\{\sup _{[0, s]} \omega(r)-\omega(s)<\varepsilon_{k}\right\}} d s
$$

we have $L_{t}=\Phi_{t}\left(X_{s}, 0 \leq s \leq t\right)$, for every $t \geq 0, P$ a.s.
Recall the notation $\widehat{X}^{(t)}$ for the process $X$ time-reversed at time $t$.
Proposition 1.1.4 For any nonnegative measurable functional $F$ on the Skorokhod space $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$,

$$
N\left(\int_{0}^{\sigma} d t F\left(\widehat{X}_{s \wedge t}^{(t)}, s \geq 0\right)\right)=E\left[\int_{0}^{L_{\infty}} d x F\left(X_{s \wedge L^{-1}(x)}, s \geq 0\right)\right]
$$

Proof. We may assume that $F$ is bounded and continuous. Fix $t>0$ and if $\omega \in$ $\mathbb{D}([0, t], \mathbb{R})$, set $T_{\max }(\omega)=\inf \left\{s \in[0, t]: \sup _{[0, s]} \omega(r)=\sup _{[0, t]} \omega(r)\right\}$ and let $\theta \omega \in$ $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ be defined by $\theta \omega(t)=\omega\left(t \wedge T_{\max }(\omega)\right)$. Let $z>0$. Excursion theory for $X-I$ shows that, for every $\varepsilon>0$,

$$
N\left(\int_{0}^{\sigma} d t 1_{\left\{\Phi_{t}\left(\hat{X}^{(t)} \leq z\right\}\right.} F\left(\widehat{X}_{s \wedge t}^{(t)}, s \geq 0\right)\right)=\frac{1}{\varepsilon} E\left[\int_{0}^{T_{\varepsilon}} d t 1_{\left\{\Phi_{t}\left(\hat{X}^{(t)} \leq z\right\}\right.} F \circ \theta\left(\widehat{X}^{(t)}\right)\right] .
$$

In deriving this equality, we also apply to the time-reversed process $\widehat{X}^{(t)}$ the fact that the local time $L_{s}$ does not increase after the (first) time of the maximum over $[0, t]$. Then,

$$
\begin{aligned}
\frac{1}{\varepsilon} E & {\left[\int_{0}^{T_{\varepsilon}} d t 1_{\left\{\Phi_{t}\left(\hat{X}^{(t)}\right) \leq z\right\}} F \circ \theta\left(\widehat{X}^{(t)}\right)\right] } \\
& =\frac{1}{\varepsilon} E\left[\int_{0}^{\infty} d t 1_{\left\{I_{t}>-\varepsilon\right\}} 1_{\left\{\Phi_{t}\left(\hat{X}^{(t)}\right) \leq z\right\}} F \circ \theta\left(\widehat{X}^{(t)}\right)\right] \\
& =\frac{1}{\varepsilon} \int_{0}^{\infty} d t E\left[1_{\left\{S_{t}-X_{t}<\varepsilon\right\}} 1_{\left\{L_{t} \leq z\right\}} F \circ \theta\left(X_{s}, s \leq t\right)\right] \\
& =E\left[\frac{1}{\varepsilon} \int_{0}^{L^{-1}(z)} d t 1_{\left\{S_{t}-X_{t}<\varepsilon\right\}} F \circ \theta\left(X_{s}, s \leq t\right)\right]
\end{aligned}
$$

We then take $\varepsilon=\varepsilon_{k}$ and pass to the limit $k \rightarrow \infty$, using the $L^{2}$ bounds provided by Lemma 1.1.3. Note that the measures

$$
\frac{1}{\varepsilon_{k}} 1_{\left[0, L^{-1}(z)\right]}(t) 1_{\left\{S_{t}-X_{t}<\varepsilon_{k}\right\}} d t
$$

converge weakly to the finite measure $1_{\left[0, L^{-1}(z)\right]}(t) d L_{t}$. Furthermore, $\theta\left(X_{s}, s \leq t\right)=$ $\left(X_{s \wedge t}, s \geq 0\right), d L_{t}$ a.e., a.s., and it is easy to verify that the mapping $t \longrightarrow F \circ \theta\left(X_{s}, s \leq\right.$ $t$ ) is continuous on a set of full $d L_{t}$-measure. We conclude that

$$
\begin{aligned}
N\left(\int_{0}^{\sigma} d t 1_{\left\{\Phi_{t}\left(\hat{X}^{(t)}\right) \leq z\right\}} F\left(\widehat{X}_{s \wedge t}^{(t)}, s \geq 0\right)\right) & =E\left[\int_{0}^{L^{-1}(z)} F\left(X_{s \wedge t}, s \geq 0\right) d L_{t}\right] \\
& =E\left[\int_{0}^{z \wedge L_{\infty}} F\left(X_{s \wedge L^{-1}(x)}, s \geq 0\right) d x\right]
\end{aligned}
$$

and the desired result follows by letting $z \rightarrow \infty$.

### 1.2 The height process and the exploration process

We write $\widehat{S}_{s}^{(t)}=\sup _{[0, s]} \widehat{X}_{r}^{(t)}(0 \leq s \leq t)$ for the supremum process of $\widehat{X}^{(t)}$.
Definition 1.2.1 The height process is the real-valued process $\left(H_{t}, t \geq 0\right)$ defined as follows. First $H_{0}=0$ and for every $t>0, H_{t}$ is the local time at level 0 at time $t$ of the process $\widehat{X}^{(t)}-\widehat{S}^{(t)}$.

The normalization of local time is of course that prescribed by Lemma 1.1.1.
Note that the existence of a modification of the process $\left(H_{t}, t \geq 0\right)$ with good continuity properties is not clear from the previous definition. When $\beta>0$ however, we can use (1.4) to see that

$$
\begin{equation*}
H_{t}=\frac{1}{\beta} m\left(\left\{I_{t}^{s} ; s \leq t\right\}\right) \tag{1.7}
\end{equation*}
$$

where for $0 \leq s \leq t$,

$$
I_{t}^{s}=\inf _{s \leq r \leq t} X_{r}
$$

Clearly the right-hand side of (1.7) gives a continuous modification of $\left(H_{t}, t \geq 0\right)$. When $\beta=0$, this argument does not apply and we will see later that there may exist no continuous (or even càdlàg) modification of ( $H_{t}, t \geq 0$ ).

At the present stage, we will use the measurable modification of $\left(H_{t}, t \geq 0\right)$ with values in $[0, \infty]$ obtained by taking

$$
\begin{equation*}
H_{t}^{o}=\Phi_{t}\left(\widehat{X}_{s}^{(t)}, 0 \leq s \leq t\right)=\liminf _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{0}^{t} 1_{\left\{X_{s}<I_{t}^{s}+\varepsilon_{k}\right\}} d s \tag{1.8}
\end{equation*}
$$

where the $\varepsilon_{k}$ 's are as in (1.6). The liminf in (1.8) is a limit (and is finite) a.s. for every fixed $t \geq 0$. The following lemma shows that more is true.

Lemma 1.2.1 Almost surely for every $t \geq 0$, we have

$$
H_{s}^{o}=\lim _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{0}^{s} 1_{\left\{X_{r}<I_{s}^{r}+\varepsilon_{k}\right\}} d r<\infty,
$$

for every $s<t$ such that $X_{s-} \leq I_{t}^{s}$, and for $s=t$ if $\Delta X_{t}>0$.
Proof. Let $s$ and $t$ be as in the statement. Then there must exist a rational $u \in(s, \infty)$ such that $X_{s-} \leq I_{u}^{s}$. We can then apply to the time-reversed process $\widehat{X}^{(u)}$ the approximation result (1.6) at times $u$ and $u-s$ respectively. The desired result follows.

We denote by $M_{f}\left(\mathbb{R}_{+}\right)$the space of all finite measures on $\mathbb{R}_{+}$, which is equipped with the topology of weak convergence.

Definition 1.2.2 The exploration process is the process $\left(\rho_{t}, t \geq 0\right)$ with values in $M_{f}\left(\mathbb{R}_{+}\right)$defined as follows. For every nonnegative measurable function $f$,

$$
\begin{equation*}
\left\langle\rho_{t}, f\right\rangle=\int_{[0, t]} d_{s} I_{t}^{s} f\left(H_{s}^{o}\right) \tag{1.9}
\end{equation*}
$$

where the notation $d_{s} I_{t}^{s}$ refers to integration with respect to the nondecreasing function $s \rightarrow I_{t}^{s}$.

Since we did not exclude the value $+\infty$ for $H_{t}^{o}$ (as defined by (1.8)), it may not be obvious that the measure $\rho_{t}$ is supported on $[0, \infty)$. However, this readily follows from the previous lemma since the measure $d_{s} I_{t}^{s}$ is supported on the set $\left\{s<t: X_{s-} \leq I_{t}^{s}\right\}$ (to which we need to add the point $t$ if $\Delta X_{t}>0$ ).

Notice that if $u$ and $v$ belong to the set $\left\{s \leq t: X_{s-} \leq I_{t}^{s}\right\}$, and if $u \leq v$, then for every $r \in[0, u)$ the condition $X_{r}<I_{u}^{r}+\varepsilon_{k}$ implies $X_{r}<I_{v}^{r}+\varepsilon_{k}$, and by construction it follows that $H_{u}^{o} \leq H_{v}^{o}$. Using the previous remark on the support of the measure $d_{s} I_{t}^{s}$, we see that the measure $\rho_{t}$ is supported on $\left[0, H_{t}^{o}\right]$, for every $t \geq 0$, a.s.

The total mass of $\rho_{t}$ is

$$
\left\langle\rho_{t}, 1\right\rangle=I_{t}^{t}-I_{t}^{0}=X_{t}-I_{t} .
$$

In particular $\rho_{t}=0$ iff $X_{t}=I_{t}$.
It will be useful to rewrite the definition of $\rho_{t}$ in terms of the time-reversed process $\widehat{X}^{(t)}$. Denote by $\widehat{L}^{(t)}=\left(\widehat{L}_{s}^{(t)}, 0 \leq s \leq t\right)$ the local time at 0 of $\widehat{X}^{(t)}-\widehat{S}^{(t)}$ (in particular $\left.H_{t}^{o}=\widehat{L}_{t}^{(t)}\right)$. Note that for $t \geq 0$ fixed, we have $H_{s}^{o}=\widehat{L}_{t}^{(t)}-\widehat{L}_{s}^{(t)}$ for every $s \in[0, t]$ such that $X_{s-} \leq I_{t}^{s}$, a.s. (compare (1.6) and (1.8)). Hence, for every $t \geq 0$,

$$
\begin{equation*}
\left\langle\rho_{t}, f\right\rangle=\int_{[0, t]} d \widehat{S}_{s}^{(t)} f\left(\widehat{L}_{t}^{(t)}-\widehat{L}_{s}^{(t)}\right), \quad \text { a.s.. } \tag{1.10}
\end{equation*}
$$

If $\mu$ is a nonzero measure in $M_{f}\left(\mathbb{R}_{+}\right)$, we write supp $\mu$ for the topological support of $\mu$ and set $H(\mu)=\sup (\operatorname{supp} \mu)$. By convention $H(0)=0$. By a preceding remark, $H\left(\rho_{t}\right) \leq H_{t}^{o}$ for every $t \geq 0$, a.s.

Lemma 1.2.2 For every $t \geq 0, P\left[H\left(\rho_{t}\right)=H_{t}^{o}\right]=1$. Furthermore, almost surely for every $t>0$, we have
(i) $\rho_{t}(\{0\})=0$;
(ii) $\operatorname{supp} \rho_{t}=\left[0, H\left(\rho_{t}\right)\right]$ if $\rho_{t} \neq 0$;
(iii) $H\left(\rho_{s}\right)=H_{s}^{o}$ for every $s \in[0, t)$ such that $X_{s-} \leq I_{t}^{s}$ and for $s=t$ if $\Delta X_{t}>0$.

Proof. It is well known, and easy to prove from the strong Markov property, that the two random measures $d S_{s}$ and $d L_{s}$ have the same support a.s. Then (1.10) implies that supp $\rho_{t}=\left[0, H_{t}^{o}\right]$ a.s. for every fixed $t>0$. In particular, $P\left[H_{t}^{o}=H\left(\rho_{t}\right)\right]=1$. Similarly (1.10) implies that $P\left[\rho_{t}(\{0\})>0\right]=0$ for every fixed $t>0$. However, if we have $\rho_{t}(\{0\})>0$ for some $t \geq 0$, our definitions and the right-continuity of paths show that the same property must hold for some rational $r>t$. Property (i) follows.

Let us now prove (ii), which is a little more delicate. We already noticed that (ii) holds for every fixed $t$, a.s., hence for every rational outside a set of zero probability. Let $t>0$ with $X_{t}>I_{t}$, and set

$$
\gamma_{t}=\sup \left\{s<t: I_{t}^{s}<X_{t}\right\}
$$

We consider two different cases.
(a) Suppose first that $X_{\gamma_{t}-}<X_{t}$, which holds in particular if $\Delta X_{t}>0$. Then note that

$$
\left\langle\rho_{t}, f\right\rangle=\int_{\left[0, \gamma_{t}\right)} d_{s} I_{t}^{s} f\left(H_{s}^{o}\right)+\left(X_{t}-X_{\gamma_{t}-}\right) f\left(H_{t}^{o}\right)
$$

Thus we can find a rational $r>t$ sufficienty close to $t$, so that $\rho_{r}$ and $\rho_{t}$ have the same restriction to $\left[0, H_{t}^{o}\right)$. The fact that property (ii) holds for $r$ implies that it holds for $t$, and we see also that $H_{t}^{o}=H\left(\rho_{t}\right)$ in that case.
(b) Suppose that $X_{\gamma_{t}-}=X_{t}$. Then we set for every $\varepsilon>0$,

$$
\left\langle\rho_{t}^{\varepsilon}, f\right\rangle=\int_{[0, t]} d_{s} I_{t}^{s} 1_{\left\{I_{t}^{s}<X_{t}-\varepsilon\right\}} f\left(H_{s}^{o}\right)
$$

From the remarks following the definition of $\rho_{t}$, it is clear that there exists some $a \geq 0$ such that $\rho_{t}^{\varepsilon}$ is bounded below by the restriction of $\rho_{t}$ to $[0, a)$, and bounded above by the restriction of $\rho_{t}$ to $[0, a]$. Also note that $\rho_{t}=\lim \uparrow \rho_{t}^{\varepsilon}$ as $\varepsilon \downarrow 0$. Now, for every $\varepsilon>0$, we can pick a rational $r>t$ so that $I_{r}^{t}>X_{t}-\varepsilon$, and we have by construction

$$
\rho_{t}^{\varepsilon}=\rho_{r}^{\varepsilon+X_{r}-X_{t}}
$$

From the rational case, the support of $\rho_{r}^{\varepsilon+X_{r}-X_{t}}$ must be an interval [ $\left.0, a\right]$, and thus the same is true for $\rho_{t}^{\varepsilon}$. By letting $\varepsilon \downarrow 0$, we get (ii) for $t$.

We already obtained (iii) for $s=t$ when $\Delta X_{s}>0$ (see (a) above). If $s \in(0, t)$ is such that $X_{s-} \leq I_{t}^{s}$, we will have also $X_{s-} \leq I_{r}^{s}$ for any rational $r \in(s, t)$. Then $H_{s}^{o} \leq H_{r}^{o}=H\left(\rho_{r}\right)$, and on the other hand, it is clear that the measures $\rho_{s}$ and $\rho_{r}$ have the same restriction to $\left[0, H_{s}^{o}\right)$. Thus (ii) implies that the support of $\rho_{s}$ contains [ $0, H_{s}^{o}$ ), and so $H\left(\rho_{s}\right) \geq H_{s}^{o}$. This gives the desired result since the reverse inequality always holds.

Proposition 1.2.3 The process $\left(\rho_{t}, t \geq 0\right)$ is a càdlàg strong Markov process in $M_{f}\left(\mathbb{R}_{+}\right)$. Furthermore, $\left(\rho_{t}, t \geq 0\right)$ is càdlàg with respect to the variation distance on finite measures.

Proof. We first explain how to define the process $\rho$ started at an arbitrary initial value $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$. To this end, we introduce some notation. Let $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$and $a \geq 0$. If $a \leq\langle\mu, 1\rangle$, we let $k_{a} \mu$ be the unique finite measure on $\mathbb{R}_{+}$such that, for every $r \geq 0$,

$$
k_{a} \mu([0, r])=\mu([0, r]) \wedge(\langle\mu, 1\rangle-a)
$$

In particular, $\left\langle k_{a} \mu, 1\right\rangle=\langle\mu, 1\rangle-a$. If $a \geq\langle\mu, 1\rangle$, we take $k_{a} \mu=0$.
If $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$has compact support and $\nu \in M_{f}\left(\mathbb{R}_{+}\right)$, we define the concatenation $[\mu, \nu] \in M_{f}\left(\mathbb{R}_{+}\right)$by the formula

$$
\int[\mu, \nu](d r) f(r)=\int \mu(d r) f(r)+\int \nu(d r) f(H(\mu)+r)
$$

With this notation, the law of the process $\rho$ started at $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$is the distribution of the process $\rho^{\mu}$ defined by

$$
\begin{equation*}
\rho_{t}^{\mu}=\left[k_{-I_{t}} \mu, \rho_{t}\right], \quad t>0 . \tag{1.11}
\end{equation*}
$$

Note that this definition makes sense because $k_{-I_{t}} \mu$ has compact support, for every $t>0$ a.s.

We then verify that the process $\rho$ has the stated properties. For simplicity, we consider only the case when the initial value is 0 , that is when $\rho$ is defined as in Definition 1.2 .2 . The right-continuity of paths is straightforward from the definition since the measures $1_{\left[0, t^{\prime}\right]}(s) d_{s} I_{t^{\prime}}^{s}$ converge to $1_{[0, t]}(s) d_{s} I_{t}^{s}$ in the variation norm as $t^{\prime} \downarrow t$. Similarly, we get the existence of left-limits from the fact that the measures $1_{\left[0, t^{\prime}\right]}(s) d_{s} I_{t^{\prime}}^{s}$ converge to $1_{[0, t)}(s) d_{s} I_{t}^{s}$ in the variation norm as $t^{\prime} \uparrow t, t^{\prime}<t$. We see in particular that $\rho$ and $X$ have the same discontinuity times and that

$$
\begin{equation*}
\rho_{t}=\rho_{t-}+\Delta X_{t} \delta_{H_{t}^{o}} . \tag{1.12}
\end{equation*}
$$

We now turn to the strong Markov property. Let $T$ be a stopping time of the canonical filtration. We will express $\rho_{T+t}$ in terms of $\rho_{T}$ and the shifted process $X_{t}^{(T)}=X_{T+t}-X_{T}$. We claim that a.s. for every $t>0$

$$
\begin{equation*}
\rho_{T+t}=\left[k_{-I_{t}^{(T)}} \rho_{T}, \rho_{t}^{(T)}\right] \tag{1.13}
\end{equation*}
$$

where $\rho_{t}^{(T)}$ and $I_{t}^{(T)}$ obviously denote the analogues of $\rho_{t}$ and $I_{t}$ when $X$ is replaced by $X^{(T)}$. When we have proved (1.13), the strong Markov property of the process $\rho$ follows by standard arguments, using also our definition of the process started at a general initial value.

For the proof of (1.13), write

$$
\left\langle\rho_{T+t}, f\right\rangle=\int_{[0, T]} d_{s} I_{T+t}^{s} f\left(H_{s}^{o}\right)+\int_{(T, T+t]} d_{s} I_{T+t}^{s} f\left(H_{s}^{o}\right)
$$

We consider separately each term in the right-hand side. Introduce $u=\sup \{r \in$ $\left.(0, T]: X_{r-}<I_{T+t}^{T}\right\}$, with the usual convention $\sup \emptyset=0$. We have $I_{T+t}^{s}=I_{T}^{s}$ for $s \in[0, u)$ and $I_{T+t}^{s}=I_{T+t}^{T}$ for $s \in[u, T]$. Since $X_{T}-I_{T+t}^{T}=-I_{t}^{(T)}$, it then follows from our definitions that
$\int_{[0, T]} d_{s} I_{T+t}^{s} f\left(H_{s}^{o}\right)=\int_{[0, u)} d_{s} I_{T}^{s} f\left(H_{s}^{o}\right)+1_{\{u>0\}}\left(I_{T+t}^{T}-X_{u-}\right) f\left(H_{u}^{o}\right)=\left\langle k_{-I_{t}^{(T)}} \rho_{T}, f\right\rangle$.

Also notice that the measures $\rho_{u}$ and $k_{-I_{t}^{(T)}} \rho_{T}$ coincide, except possibly at the point $H_{u}^{o}$. In any case, $H\left(\rho_{u}\right)=H\left(k_{-I_{t}^{(T)}} \rho_{T}\right)$, and we have also $H_{u}^{o}=H\left(\rho_{u}\right)$ by Lemma 1.2.2 (iii).

Now observe that for $d_{s} I_{T+t}^{s}$ almost every $s \in(T, T+t]$, we have $H_{s}^{o}=H_{u}^{o}+H_{s-T}^{o,(T)}$, with an obvious notation. To see this, pick a rational $r>T+t$ such that $I_{r}^{T+t}>X_{s-}$ and argue on the time-reversed process $\widehat{X}^{(r)}$ as in the proof of Lemma 1.2.1. Hence,

$$
\begin{equation*}
\int_{(T, T+t]} d_{s} I_{T+t}^{s} f\left(H_{s}^{o}\right)=\int_{(T, T+t]} d_{s} I_{T+t}^{s} f\left(H_{u}^{o}+H_{s-T}^{o,(T)}\right)=\int \rho_{t}^{(T)}(d x) f\left(H_{u}^{o}+x\right) \tag{1.15}
\end{equation*}
$$

Formula (1.13) follows by combining (1.14) and (1.15).
We now come back to the problem of finding a modification of the height process with good continuity properties. By the first assertion of Lemma 1.2.2, $\left(H\left(\rho_{t}\right), t \geq 0\right)$ is a modification of $\left(H_{t}^{o}, t \geq 0\right)$. From now on, we will systematically use this modification and write $H_{t}=H\left(\rho_{t}\right)$. From Lemma 1.2.2 (iii), we see that formula (1.9) defining $\left(\rho_{t}, t \geq 0\right)$ remains true if $H_{s}^{o}$ is replaced by $H_{s}$. The same applies to formula (1.12) giving the jumps of $\rho$. Furthermore, the continuity properties of the process $\rho_{t}$ (and especially the form of its jumps) imply that the mapping $t \longrightarrow H\left(\rho_{t}\right)=H_{t}$ is lower semicontinuous a.s.

Let us make an important remark at this point. Write

$$
g_{t}=\sup \left\{s \leq t: X_{s}=I_{s}\right\}
$$

for the beginning of the excursion of $X-I$ that straddles $t$. Then a simple timereversal argument shows that a.s. for every $t$ such that $X_{t}-I_{t}>0$, we have

$$
\lim _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{0}^{g_{t}} 1_{\left\{X_{s}<I_{t}^{s}+\varepsilon_{k}\right\}} d s=0
$$

and thus we can replace (1.8) by

$$
H_{t}^{o}=\liminf _{k \rightarrow \infty} \frac{1}{\varepsilon_{k}} \int_{g_{t}}^{t} 1_{\left\{X_{s}<I_{t}^{s}+\varepsilon_{k}\right\}} d s
$$

Recalling (1.9), we see that, a.s. for every $t \geq 0$ such that $X_{t}-I_{t}>0$, we can write $\rho_{t}$ and $H_{t}$ as measurable functions of the excursion of $X-I$ that straddles $t$ (and of course $\rho_{t}=0$ and $H_{t}=0$ if $X_{t}=I_{t}$ ). We can thus define both the height process and the exploration process under the excursion measure $N$. Furthermore, if $\left(\alpha_{j}, \beta_{j}\right), j \in J$, denote the excursion intervals of $X-I$, and if $\omega_{j}, j \in J$, denote the corresponding excursions, we have $H_{t}=H_{t-\alpha_{j}}\left(\omega_{j}\right)$ and $\rho_{t}=\rho_{t-\alpha_{j}}\left(\omega_{j}\right)$ for every $t \in\left(\alpha_{j}, \beta_{j}\right)$ and $j \in J$, a.s.

Since 0 is a regular point for $X-I$, we also see that the measure 0 is a regular point for the exploration process $\rho$. It is immediate from the previous remark that the excursion measure of $\rho$ away from 0 is the "law" of $\rho$ under $N$. Similarly, the process $-I$, which is the local time at 0 for $X-I$, is also the local time at 0 for $\rho$.

We now state and prove a useful technical lemma about the process $H$.
Lemma 1.2.4 (Intermediate values property) Almost surely for every $t<t^{\prime}$, the process $H$ takes all values between $H_{t}$ and $H_{t^{\prime}}$ on the time interval $\left[t, t^{\prime}\right]$.

Proof. First consider the case when $H_{t}>H_{t^{\prime}}$. By using the lower semi-continuity of $H$, we may assume that $t$ is rational. From (1.13), we have $\rho_{t+s}=\left[k_{-I_{s}^{(t)}} \rho_{t}, \rho_{s}^{(t)}\right]$ for every $s>0$, a.s. Hence, if

$$
\gamma_{r}=\inf \left\{s>0: I_{s}^{(t)}=-r\right\}
$$

we have $\rho_{t+\gamma_{r}}=k_{r} \rho_{t}$, and so $H_{t+\gamma_{r}}=H\left(k_{r} \rho_{t}\right)$ for every $r \geq 0$, a.s. However, Lemma 1.2.2 (ii) implies that the mapping $r \rightarrow H\left(k_{r} \rho_{t}\right)$ is continuous. Now note that for $r=0, H_{t+\gamma_{r}}=H_{t}$, whereas for $r=X_{t}-I_{t^{\prime}}^{t}=-I_{t^{\prime}-t}^{(t)}$ we have $t+\gamma_{r} \leq t^{\prime}$ and our definitions easily imply $\rho_{t+\gamma_{r}} \leq \rho_{t^{\prime}}$ and $H_{t+\gamma_{r}} \leq H_{t^{\prime}}$.

Consider then the case when $H_{t}<H_{t^{\prime}}$. By lower semi-continuity again, we may assume that $t^{\prime}$ is rational. In terms of the process time-reversed at time $t^{\prime}$, we have $H_{t^{\prime}}=\widehat{L}_{t^{\prime}}^{\left(t^{\prime}\right)}$. Set

$$
\sigma_{r}=\inf \left\{s \geq 0: \widehat{S}_{s}^{\left(t^{\prime}\right)} \geq r\right\}
$$

which is well defined for $0 \leq r \leq X_{t^{\prime}}-I_{t^{\prime}}$. Since the subordinator $S_{L^{-1}(t)}$ is strictly increasing, we see that the mapping $r \rightarrow \widehat{L}_{\sigma_{r}}^{\left(t^{\prime}\right)}$ is continuous for $r \in\left[0, X_{t^{\prime}}-I_{t^{\prime}}\right]$, a.s. Now note that

$$
H_{t^{\prime}-\sigma_{r}}=\widehat{L}_{t^{\prime}}^{\left(t^{\prime}\right)}-\widehat{L}_{\sigma_{r}}^{\left(t^{\prime}\right)}
$$

for every $r \in\left[0, X_{t^{\prime}}-I_{t^{\prime}}\right]$, a.s. For $r=X_{t^{\prime}}-I_{t^{\prime}}^{t}=\widehat{S}_{t^{\prime}-t}^{\left(t^{\prime}\right)}$, we have $t^{\prime}-\sigma_{r} \geq t$ and $H_{t^{\prime}-\sigma_{r}} \leq H_{t}$ by construction. The desired result follows.

The next proposition is a corollary of Proposition 1.1.4. We denote by $U$ a subordinator defined on an auxiliary probability space $\left(\Omega^{0}, P^{0}\right)$, with Laplace exponent

$$
E^{0}\left[\exp -\lambda U_{t}\right]=\exp \left(-t\left(\beta \lambda+\int_{0}^{\infty}\left(1-e^{-\lambda r}\right) \pi([r, \infty)) d r\right)\right)=\exp (-t(\tilde{\psi}(\lambda)-\alpha))
$$

For every $a \geq 0$, we let $J_{a}$ be the random element of $M_{f}\left(\mathbb{R}_{+}\right)$defined by $J_{a}(d r)=$ $1_{[0, a]}(r) d U_{r}$.

Proposition 1.2.5 For every nonnegative measurable function $\Phi$ on $M_{f}\left(\mathbb{R}_{+}\right)$,

$$
N\left(\int_{0}^{\sigma} d t \Phi\left(\rho_{t}\right)\right)=\int_{0}^{\infty} d a e^{-\alpha a} E^{0}\left[\Phi\left(J_{a}\right)\right]
$$

Let $b \geq 0$. Then $\rho_{s}(\{b\})=0$ for every $s \geq 0, N$ a.e. or $P$ a.s.
Proof. We have $\rho_{t}=\Sigma\left(\widehat{X}_{s \wedge t}^{(t)}, s \geq 0\right)$, with a functional $\Sigma$ that is made explicit in (1.10). We then apply Proposition 1.1.4 to obtain

$$
N\left(\int_{0}^{\sigma} d t \Phi\left(\rho_{t}\right)\right)=E\left[\int_{0}^{L \infty} d a \Phi \circ \Sigma\left(X_{s \wedge L^{-1}(a)}, s \geq 0\right)\right] .
$$

However, for $a<L_{\infty}$,

$$
\left\langle\Sigma\left(X_{s \wedge L^{-1}(a)}, s \geq 0\right), f\right\rangle=\int_{0}^{L^{-1}(a)} d S_{s} f\left(a-L_{s}\right)
$$

The first assertion is now a consequence of Lemma 1.1.2, which shows that $P[a<$ $\left.L_{\infty}\right]=e^{-\alpha a}$ and that, conditionally on $\left\{L^{-1}(a)<\infty\right\}, S_{L^{-1}(r)}=X_{L^{-1}(r)}, 0 \leq r \leq a$ has the same distribution as $U$ (we also use the fact that $\left(U_{a}-U_{a-r}, 0 \leq r \leq a\right)$ has the same distribution as $\left(U_{r}, 0 \leq r \leq a\right)$ ).

Consider now the second assertion. Note that the case $b=0$ is a consequence of Lemma 1.2.2 (i). So we may assume that $b>0$ and it is enough to prove the result under the excursion measure $N$. However, since $b$ is a.s. not a jump time of $U$, the right side of the formula of the proposition vanishes for $\Phi(\mu)=\mu(\{b\})$. The desired result follows, using also the fact that $\rho$ is càdlàg in the variation norm.

We denote by $\mathbf{M}$ the measure on $M_{f}\left(\mathbb{R}_{+}\right)$defined by:

$$
\langle\mathbf{M}, \Phi\rangle=\int_{0}^{\infty} d a e^{-\alpha a} E^{0}\left[\Phi\left(J_{a}\right)\right]
$$

Proposition 1.2.5 implies that the measure $\mathbf{M}$ is invariant for $\rho$.
The last proposition of this section describes the potential kernel of the exploration process killed when it hits 0 . We fix $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$and let $\rho^{\mu}$ be as in (1.11) the exploration process started at $\mu$. We use the notation introduced in the proof of Proposition 1.2.3.

Proposition 1.2.6 Let $\tau_{0}=\inf \left\{s \geq 0, \rho_{s}^{\mu}=0\right\}$. Then,

$$
E\left[\int_{0}^{\tau_{0}} d s \Phi\left(\rho_{s}^{\mu}\right)\right]=\int_{0}^{<\mu, 1>} d r \int \mathbf{M}(d \theta) \Phi\left(\left[k_{r} \mu, \theta\right]\right)
$$

Proof. First note that $\tau_{0}=T_{<\mu, 1>}$ by an immediate application of the definition of $\rho^{\mu}$. Then, denote by $\left(\alpha_{j}, \beta_{j}\right), j \in J$ the excursion intervals of $X-I$ away from 0 before time $T_{<\mu, 1>}$, and by $\omega_{j}, j \in J$ the corresponding excursions. As we observed before Proposition 1.2.5, we have $\rho_{t}=\rho_{t-\alpha_{j}}\left(\omega_{j}\right)$ for every $t \in\left(\alpha_{j}, \beta_{j}\right), j \in J$, a.s. Since $\left\{s \geq 0: X_{s}=I_{s}\right\}$ has zero Lebesgue measure a.s., it follows that

$$
E\left[\int_{0}^{\tau_{0}} d s \Phi\left(\rho_{s}^{\mu}\right)\right]=E\left[\sum_{j \in J} \int_{0}^{\beta_{j}-\alpha_{j}} d r \Phi\left(\left[k_{-I_{\alpha_{j}}} \mu, \rho_{r}\left(\omega_{j}\right)\right]\right)\right]
$$

By excursion theory, the point measure

$$
\sum_{j \in J} \delta_{I_{\alpha_{j}}, \omega_{j}}(d u d e)
$$

is a Poisson point measure with intensity $1_{[-<\mu, 1>, 0]}(u) d u N(d \omega)$. Hence,

$$
E\left[\int_{0}^{\tau_{0}} d s \Phi\left(\rho_{s}^{\mu}\right)\right]=\int_{0}^{<\mu, 1>} d u N\left(\int_{0}^{\sigma} d r \Phi\left(\left[k_{u} \mu, \rho_{r}\right]\right)\right)
$$

and the desired result follows from Proposition 1.2.5.

### 1.3 Local times of the height process

### 1.3.1 The construction of local times

Our goal is to construct a local time process for $H$ at each level $a \geq 0$. Since $H$ is in general not Markovian (and not a semimartingale) we cannot apply a general theory, but still we will use certain ideas which are familiar in the Brownian motion setting. In the case $a=0$, we can already observe that $H_{t}=0$ iff $\rho_{t}=0$ or equivalently $X_{t}-I_{t}=0$. Therefore the process $-I$ is the natural candidate for the local time of $H$ at level 0 .

Let us fix $a \geq 0$. Since $t \rightarrow \rho_{t}$ is càdlàg in the variation norm, it follows that the mapping $t \rightarrow \rho_{t}((a, \infty))$ is càdlàg. Furthermore, it follows from (1.12) that the discontinuity times of this mapping are exactly those times $t$ such that $\Delta X_{t}>0$ and $H_{t}>a$, and the corresponding jump is $\Delta X_{t}$.

Set

$$
\tau_{t}^{a}=\inf \left\{s \geq 0: \int_{0}^{s} 1_{\left\{H_{r}>a\right\}} d r>t\right\}=\inf \left\{s \geq 0: \int_{0}^{s} 1_{\left\{\rho_{r}((a, \infty))>0\right\}} d r>t\right\}
$$

From Proposition 1.2.5, we get that $\int_{0}^{\infty} 1_{\left\{H_{r}>a\right\}} d r=\infty$ a.s., so that the random times $\tau_{t}^{a}$ are a.s. finite.

When $a>0$, we also set

$$
\tilde{\tau}_{t}^{a}=\inf \left\{s \geq 0: \int_{0}^{s} 1_{\left\{H_{r} \leq a\right\}} d r>t\right\}
$$

and we let $\mathcal{H}_{a}$ be the $\sigma$-field generated by the càdlàg process $\left(X_{\tilde{\tau}_{t}^{a}}, \rho_{\tilde{\tau}_{t}^{a}} ; t \geq 0\right)$ and the class of $P$-negligible sets of $\mathcal{G}_{\infty}$. We also define $\mathcal{H}_{0}$ as the $\sigma$-field generated by the class of $P$-negligible sets of $\mathcal{G}_{\infty}$.

Proposition 1.3.1 For every $t \geq 0$, let $\rho_{t}^{a}$ be the random measure on $\mathbb{R}_{+}$defined by

$$
\left\langle\rho_{t}^{a}, f\right\rangle=\int_{(a, \infty)} \rho_{\tau_{t}^{a}}(d r) f(r-a)
$$

The process $\left(\rho_{t}^{a}, t \geq 0\right)$ has the same distribution as $\left(\rho_{t}, t \geq 0\right)$ and is independent of $\mathcal{H}_{a}$.

Proof. First step. We first verify that the process $\left(\left\langle\rho_{t}^{a}, 1\right\rangle, t \geq 0\right)$ has the same distribution as $\left(\left\langle\rho_{t}, 1\right\rangle, t \geq 0\right)$.

Let $\varepsilon>0$. We introduce two sequences of stopping times $S_{\varepsilon}^{k}, T_{\varepsilon}^{k}, k \geq 1$, defined inductively as follows:

$$
\begin{aligned}
& S_{\varepsilon}^{1}=\inf \left\{s \geq 0: \rho_{s}((a, \infty)) \geq \varepsilon\right\} \\
& T_{\varepsilon}^{k}=\inf \left\{s \geq S_{\varepsilon}^{k}: \rho_{s}((a, \infty))=0\right\} \\
& S_{\varepsilon}^{k+1}=\inf \left\{s \geq T_{\varepsilon}^{k}: \rho_{s}((a, \infty)) \geq \varepsilon\right\}
\end{aligned}
$$

It is easy to see that these stopping times are a.s. finite, and $S_{\varepsilon}^{k} \uparrow \infty, T_{\varepsilon}^{k} \uparrow \infty$ as $k \uparrow \infty$.

From (1.13) applied with $T=S_{\varepsilon}^{k}$, we obtain that, for every $k \geq 1$,

$$
\begin{equation*}
T_{\varepsilon}^{k}=\inf \left\{s>S_{\varepsilon}^{k}: X_{s}=X_{S_{\varepsilon}^{k}}-\rho_{S_{\varepsilon}^{k}}((a, \infty))\right\} \tag{1.16}
\end{equation*}
$$

Formula (1.13) also implies that, for every $s \in\left[0, T_{\varepsilon}^{k}-S_{\varepsilon}^{k}\right]$,

$$
\begin{align*}
\rho_{S_{\varepsilon}^{k}+s}((a, \infty)) & =\left(\rho_{S_{\varepsilon}^{k}}((a, \infty))+I_{s}^{\left(S_{\varepsilon}^{k}\right)}\right)+\left(X_{s}^{\left(S_{\varepsilon}^{k}\right)}-I_{s}^{\left(S_{\varepsilon}^{k}\right)}\right)  \tag{1.17}\\
& =X_{S_{\varepsilon}^{k}+s}-\left(X_{S_{\varepsilon}^{k}}-\rho_{S_{\varepsilon}^{k}}((a, \infty))\right)
\end{align*}
$$

We set

$$
Y_{s}^{k, \varepsilon}=\rho_{\left(S_{\varepsilon}^{k}+s\right) \wedge T_{\varepsilon}^{k}}((a, \infty)) .
$$

As a straightforward consequence of (1.16) and (1.17), conditionally on $\mathcal{G}_{S_{\varepsilon}^{k}}$, the process $Y^{k, \varepsilon}$ is distributed as the underlying Lévy process started at $\rho_{S_{\varepsilon}^{k}}((a, \infty))$ and stopped at its first hitting time of 0 .

We then claim that, for every $t \geq 0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \sup _{\left\{k \geq 1, S_{\varepsilon}^{k} \leq t\right\}} \rho_{S_{\varepsilon}^{k}}((a, \infty))=0, \quad \text { a.s. } \tag{1.18}
\end{equation*}
$$

Indeed, by previous observations about the continuity properties of the mapping $s \longrightarrow$ $\rho_{s}((a, \infty))$, we have

$$
\sup _{\left\{k \geq 1, S_{\varepsilon}^{k} \leq t\right\}} \rho_{S_{\varepsilon}^{k}}((a, \infty)) \leq \varepsilon+\sup \left\{\Delta X_{s} ; s \leq t, H_{s}>a, \rho_{s}((a, \infty)) \leq \varepsilon\right\}
$$

However, the sets $\left\{s \leq t ; \Delta X_{s}>0, H_{s}>a, \rho_{s}((a, \infty)) \leq \varepsilon\right\}$ decrease to $\emptyset$ as $\varepsilon \downarrow 0$, and so

$$
\lim _{\varepsilon \rightarrow 0}\left(\sup \left\{\Delta X_{s} ; s \leq t, H_{s}>a, \rho_{s}((a, \infty)) \leq \varepsilon\right\}\right)=0
$$

a.s., which yields the desired claim.

Set

$$
Z_{s}^{k, \varepsilon}=Y_{s}^{k, \varepsilon}-\inf _{0 \leq r \leq s} Y_{r}^{k, \varepsilon} \leq Y_{s}^{k, \varepsilon}
$$

Then, conditionally on $\mathcal{G}_{S_{\varepsilon}^{k}}, Z^{k, \varepsilon}$ is distributed as an independent copy of the reflected process $X-I$, stopped when its local time at 0 hits $\rho_{S_{\varepsilon}^{k}}((a, \infty))$.

Denote by $U^{\varepsilon}=\left(U_{s}^{\varepsilon}, s \geq 0\right)$ the process obtained by pasting together the paths $\left(Z_{s}^{k, \varepsilon}, 0 \leq s \leq T_{\varepsilon}^{k}-S_{\varepsilon}^{k}\right)$. By the previous remarks, $U^{\varepsilon}$ is distributed as the reflected Lévy process $X-I$.

We then set

$$
\tau_{s}^{a, \varepsilon}=\inf \left\{t \geq 0: \int_{0}^{t} \sum_{k=1}^{\infty} 1_{\left[S_{\varepsilon}^{k}, T_{\varepsilon}^{k}\right]}(r) d r>s\right\}
$$

Observe that the time-changed process $\left(\rho_{\tau_{s}^{a, \varepsilon}}((a, \infty)), s \geq 0\right)$ is obtained by patching together the paths $\left(Y_{s}^{k, \varepsilon}, 0 \leq s \leq T_{\varepsilon}^{k}-S_{\varepsilon}^{k}\right)$. Moreover, we have for every $k \geq 1$,

$$
\sup _{0 \leq s \leq T_{\varepsilon}^{k}-S_{\varepsilon}^{k}}\left(Y_{s}^{k, \varepsilon}-Z_{s}^{k, \varepsilon}\right)=\rho_{S_{\varepsilon}^{k}}((a, \infty))=Y_{0}^{k, \varepsilon}
$$

From (1.18), we conclude that for every $t \geq 0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left(\sup _{s \leq t}\left|U_{s}^{\varepsilon}-\rho_{\tau_{s}^{a, \varepsilon}}((a, \infty))\right|\right)=0 \tag{1.19}
\end{equation*}
$$

Notice that $\tau_{s}^{a, \varepsilon} \downarrow \tau_{s}^{a}$ as $\varepsilon \downarrow 0$ and recall that for every $\varepsilon>0, U^{\varepsilon}$ is distributed as the reflected Lévy process $X-I$. We then get from (1.19) that the process $\left\langle\rho_{s}^{a}, 1\right\rangle=$ $\rho_{\tau_{s}^{a}}((a, \infty))$ is distributed as the reflected proces $X-I$, which completes the first step.

Second step. We will now verify that $\rho^{a}$ can be obtained as a functional of the total mass process $\left\langle\rho^{a}, 1\right\rangle$ in the same way as $\rho$ is obtained from $\langle\rho, 1\rangle$. It will be enough to argue on one excursion of $\left\langle\rho^{a}, 1\right\rangle$ away from 0 . Thus, let $(u, v)$ be the interval corresponding to one such excursion. We can associate with ( $u, v$ ) a unique connected component $(p, q)$ of the open set $\left\{s, H_{s}>a\right\}$, such that $\tau_{u+r}^{a}=p+r$ for every $r \in[0, v-u)$, and $q=\tau_{v-}^{a}$. By the intermediate values property, we must have $H_{p}=H_{q}=a$.

We also claim that $X_{r}>X_{p}$ for every $r \in(p, q)$. If this were not the case, we could find $r \in(p, q)$ such that $X_{r}=\inf \left\{X_{s}, p \leq s \leq r\right\}$, which forces $H_{r} \leq H_{p}=a$ and gives a contradiction.

The previous observations and the definition of the process $\rho$ imply that, for every $r \in(p, q)$, the restriction of $\rho_{r}$ to $[0, a]$ is exactly $\rho_{p}=\rho_{q}$. Define

$$
\omega(r)=X_{(p+r) \wedge q}-X_{p}=\left\langle\rho_{(p+r) \wedge q}, 1\right\rangle-\left\langle\rho_{p}, 1\right\rangle=\left\langle\rho_{(u+r) \wedge v}^{a}, 1\right\rangle
$$

so that $\omega$ is the excursion of $\left\langle\rho^{a}, 1\right\rangle$ corresponding to $(u, v)$. The construction of the process $\rho$ implies that, for $0<r<q-p=v-u$,

$$
\rho_{p+r}=\left[\rho_{p}, \rho_{r}(\omega)\right],
$$

and so, for the same values of $r$,

$$
\rho_{u+r}^{a}=\rho_{r}(\omega)
$$

This completes the second step of the proof.
Third step. It remains to prove that $\rho^{a}$ is independent of the $\sigma$-field $\mathcal{H}_{a}$. For $\varepsilon>0$, denote by $\mathcal{H}_{a}^{\varepsilon}$ the $\sigma$-field generated by the processes

$$
\left(X_{\left(T_{\varepsilon}^{k}+s\right) \wedge S_{\varepsilon}^{k+1}}, s \geq 0\right)
$$

for $k=0,1, \ldots$ (by convention $T_{\varepsilon}^{0}=0$ ), and the negligible sets of $\mathcal{G}_{\infty}$. From our construction (in particular the fact that $X_{s}>X_{T_{\varepsilon}^{k}}$ for $s \in\left[S_{\varepsilon}^{k}, T_{\varepsilon}^{k}\right)$ ), it is easy to verify that the processes $\left(\rho_{\left(T_{\varepsilon}^{k}+s\right) \wedge S_{\varepsilon}^{k+1}}, s \geq 0\right)$ are measurable with respect to $\mathcal{H}_{a}^{\varepsilon}$, and since $H_{t}>a$ for $t \in\left(S_{\varepsilon}^{k}, T_{\varepsilon}^{k}\right)$, it follows that $\mathcal{H}_{a} \subset \mathcal{H}_{a}^{\varepsilon}$.

By the arguments of the first step, the processes $Z^{k, \varepsilon}$ are independent conditionally on $\mathcal{H}_{a}^{\varepsilon}$, and the conditional law of $Z^{k, \varepsilon}$ is the law of an independent copy of the reflected process $X-I$, stopped when its local time at 0 hits $\rho_{S_{\varepsilon}^{k}}((a, \infty))$. It follows that the process $U^{\varepsilon}$ of the first step is independent of $\mathcal{H}_{a}^{\varepsilon}$, hence also of $\mathcal{H}_{a}$. By passing to the limit $\varepsilon \rightarrow 0$, we obtain that the total mass process $\left\langle\rho^{a}, 1\right\rangle$ is independent of $\mathcal{H}_{a}$. As we know that $\rho^{a}$ can be reconstructed as a measurable functional of its total mass process, this completes the proof.

We let $l^{a}=\left(l^{a}(s), s \geq 0\right)$ be the local time at 0 of $\left\langle\rho^{a}, 1\right\rangle$, or equivalently of $\rho^{a}$.
Definition 1.3.1 The local time at level $a$ and at time $s$ of the height process $H$ is defined by

$$
L_{s}^{a}=l^{a}\left(\int_{0}^{s} 1_{\left\{H_{r}>a\right\}} d r\right)
$$

This definition will be justified below: See in particular Proposition 1.3.3.

### 1.3.2 Some properties of local times

The next lemma can be seen as dual to Lemma 1.1.3.
Lemma 1.3.2 For every $t \geq 0$,

$$
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_{0}^{t} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s=-I_{t}
$$

in the $L^{1}$-norm.
Proof. We use arguments similar to the proof of Lemma 1.1.3. Recall that $T_{x}=$ $\inf \left\{t \geq 0: X_{t}=-x\right\}$. We first establish that for every $x>0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_{0}^{T_{x}} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s=x \tag{1.20}
\end{equation*}
$$

in probability. Note that Proposition 1.2 .5 gives for any nonnegative measurable function $g$

$$
N\left(\int_{0}^{\sigma} d s g\left(H_{s}\right)\right)=\int_{0}^{\infty} d a e^{-\alpha a} g(a)
$$

Let $\omega^{j}, j \in J$ denote the excursions of $X-I$ away from 0 and let $\left(\alpha_{j}, \beta_{j}\right)$ be the corresponding time intervals. We already noticed that $H_{s}=H_{s-\alpha_{j}}\left(\omega^{j}\right)$ for $s \in$ $\left(\alpha_{j}, \beta_{j}\right)$. Hence, using also the previous displayed formula, we have

$$
\begin{equation*}
E\left[\frac{1}{\varepsilon} \int_{0}^{T_{x}} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s\right]=\frac{x}{\varepsilon} N\left(\int_{0}^{\sigma} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s\right)=\frac{x}{\varepsilon}\left(\frac{1-e^{-\alpha \varepsilon}}{\alpha}\right) \leq x \tag{1.21}
\end{equation*}
$$

and in particular,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} E\left[\frac{1}{\varepsilon} \int_{0}^{T_{x}} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s\right]=x \tag{1.22}
\end{equation*}
$$

We then want to get a second moment estimate. To this end, it is necessary to introduce a suitable truncation. Fix $K>0$. A slight modification of the proof of (1.22) gives

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} E\left[\frac{1}{\varepsilon} \int_{0}^{T_{x}} 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s}-I_{s} \leq K\right\}} d s\right]=x \tag{1.23}
\end{equation*}
$$

If $H^{(s)}$ denotes the height process for the shifted process $X_{t}^{(s)}=X_{s+t}-X_{s}$, the bound $H_{t-s}^{(s)} \leq H_{t}$ (for $0 \leq s \leq t$ ) is obvious from our construction. We can use this simple observation to bound

$$
\begin{align*}
N & \left(\left(\int_{0}^{\sigma} 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s} \leq K\right\}} d s\right)^{2}\right) \\
& \leq 2 N\left(\int_{0}^{\sigma} d s 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s} \leq K\right\}} \int_{s}^{\sigma} d t 1_{\left\{H_{t} \leq \varepsilon\right\}}\right) \\
& \leq 2 N\left(\int_{0}^{\sigma} d s 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s} \leq K\right\}} \int_{s}^{\sigma} d t 1_{\left\{H_{t-s}^{(s)} \leq \varepsilon\right\}}\right) \\
& =2 N\left(\int_{0}^{\sigma} d s 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s} \leq K\right\}} E_{X_{s}}\left[\int_{0}^{T_{0}} d t 1_{\left\{H_{t} \leq \varepsilon\right\}}\right]\right) \\
& \leq 2 \varepsilon N\left(\int_{0}^{\sigma} d s 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s} \leq K\right\}} X_{s}\right)  \tag{1.21}\\
& =2 \varepsilon \int_{0}^{\varepsilon} d y E\left[X_{L^{-1}(y)} 1_{\left\{L^{-1}(y)<\infty, X_{L^{-1}(y)} \leq K\right\}}\right] \\
& \leq 2 \varepsilon^{2} E\left[X_{L^{-1}(\varepsilon)} \wedge K\right] \\
& =o\left(\varepsilon^{2}\right)
\end{align*}
$$

(Proposition 1.1.4)
by dominated convergence. As in the proof of Lemma 1.1.3, we can conclude from (1.22) and the previous estimate that

$$
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_{0}^{T_{x}} 1_{\left\{H_{s} \leq \varepsilon\right\}} 1_{\left\{X_{s}-I_{s} \leq K\right\}} d s=x
$$

in the $L^{2}$-norm. Since this holds for every $K>0,(1.20)$ follows.
From (1.20) and a monotonicity argument we deduce that the convergence of Lemma 1.3.2 holds in probability. To get $L^{1}$-convergence, we need a few other estimates. First observe that

$$
\begin{equation*}
E\left[\int_{0}^{t} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s\right]=\int_{0}^{t} d s P\left[H_{s} \leq \varepsilon\right]=\int_{0}^{t} d s P\left[L_{s} \leq \varepsilon\right]=E\left[L^{-1}(\varepsilon) \wedge t\right] \leq C \varepsilon \tag{1.24}
\end{equation*}
$$

with a constant $C$ depending only on $t$ (in the last bound we applied to $L^{-1}$ an estimate valid for any subordinator). Then,

$$
\begin{aligned}
E\left[\left(\int_{0}^{t} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s\right)^{2}\right] & =2 E\left[\iint_{\{0<r<s<t\}} d r d s 1_{\left\{H_{r} \leq \varepsilon\right\}} 1_{\left\{H_{s} \leq \varepsilon\right\}}\right] \\
& \leq 2 E\left[\iint_{\{0<r<s<t\}} d r d s 1_{\left\{H_{r} \leq \varepsilon\right\}} 1_{\left\{H_{s-r}^{(r)} \leq \varepsilon\right\}}\right] \\
& =2 E\left[\int_{0}^{t} d r 1_{\left\{H_{r} \leq \varepsilon\right\}} E\left[\int_{0}^{t-r} d s 1_{\left\{H_{s} \leq \varepsilon\right\}}\right]\right] \\
& \leq 2\left(E\left[\int_{0}^{t} d r 1_{\left\{H_{r} \leq \varepsilon\right\}}\right]\right)^{2}
\end{aligned}
$$

As a consequence of the last estimate and (1.24), the variables $\varepsilon^{-1} \int_{0}^{t} 1_{\left\{H_{s} \leq \varepsilon\right\}} d s, \varepsilon>0$ are bounded in $L^{2}$. This completes the proof of Lemma 1.3.2.

We can now give a useful approximation result for local times of the height process.
Proposition 1.3.3 For every $t \geq 0$,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{a \geq 0} E\left[\sup _{s \leq t}\left|\varepsilon^{-1} \int_{0}^{s} 1_{\left\{a<H_{r} \leq a+\varepsilon\right\}} d r-L_{s}^{a}\right|\right]=0 .
$$

Similarly, for every $t \geq 0$,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{a \geq \varepsilon} E\left[\sup _{s \leq t}\left|\varepsilon^{-1} \int_{0}^{s} 1_{\left\{a-\varepsilon<H_{r} \leq a\right\}} d r-L_{s}^{a}\right|\right]=0
$$

There exists a jointly measurable modification of the collection ( $L_{s}^{a}, a \geq 0, s \geq 0$ ), which is continuous and nondecreasing in the variable s, and such that, a.s. for any nonnegative measurable function $g$ on $\mathbb{R}_{+}$and any $s \geq 0$,

$$
\begin{equation*}
\int_{0}^{s} g\left(H_{r}\right) d r=\int_{\mathbb{R}_{+}} g(a) L_{s}^{a} d a . \tag{1.25}
\end{equation*}
$$

Proof. First consider the case $a=0$. Then, $\rho^{0}=\rho$ and $L_{t}^{0}=l^{0}(t)=-I_{t}$. Lemma 1.3.2 and a simple monotonicity argument, using the continuity of $L_{s}^{0}$, give

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} E\left[\sup _{s \leq t}\left|\varepsilon^{-1} \int_{0}^{s} 1_{\left\{0<H_{r} \leq \varepsilon\right\}} d r-L_{s}^{0}\right|\right]=0 \tag{1.26}
\end{equation*}
$$

For $a>0$, set $A_{t}^{a}=\int_{0}^{t} 1_{\left\{H_{s}>a\right\}} d s$. Note that $\left\{a<H_{s} \leq a+\varepsilon\right\}=\left\{\rho_{s}((a, \infty))>\right.$ $0\} \cap\left\{\rho_{s}((a+\varepsilon, \infty))=0\right\}$, and so

$$
\begin{aligned}
\int_{0}^{s} 1_{\left\{a<H_{r} \leq a+\varepsilon\right\}} d r & =\int_{0}^{t} 1_{\left\{\rho_{s}((a, \infty))>0\right\} \cap\left\{\rho_{s}((a+\varepsilon, \infty))=0\right\}} d s \\
& =\int_{0}^{A_{t}^{a}} 1_{\left\{\rho_{r}^{a}((\varepsilon, \infty))=0\right\}} d r \\
& =\int_{0}^{A_{t}^{a}} 1_{\left\{0<H_{r}^{a} \leq \varepsilon\right\}} d r,
\end{aligned}
$$

where $H_{t}^{a}=H\left(\rho_{t}^{a}\right)$. The first convergence of the proposition then follows from (1.26), the trivial bound $A_{t}^{a} \leq t$ and the fact that $\rho^{a}$ has the same distribution as $\rho$.

The second convergence is easily derived from the first one by elementary arguments. Let us only sketch the method. For any fixed $\delta>0$, we can choose $\varepsilon_{0}>0$ sufficiently small so that for every $a \geq 0, \varepsilon \in\left(0, \varepsilon_{0}\right]$,

$$
\begin{equation*}
E\left[\left|\varepsilon^{-1} \int_{0}^{t} 1_{\left\{a<H_{r} \leq a+\varepsilon\right\}} d r-L_{s}^{a}\right|\right] \leq \delta \tag{1.27}
\end{equation*}
$$

Then, if $0<\varepsilon<\varepsilon_{0} \wedge a$,

$$
E\left[\left|\varepsilon^{-1} \int_{0}^{t} 1_{\left\{a-\varepsilon<H_{r} \leq a\right\}} d r-\varepsilon_{0}^{-1} \int_{0}^{t} 1_{\left\{a-\varepsilon<H_{r} \leq a-\varepsilon+\varepsilon_{0}\right\}} d r\right|\right]<2 \delta
$$

However, if $\varepsilon$ is very small in comparison with $\varepsilon_{0}$, one also gets the bound

$$
E\left[\left|\varepsilon_{0}^{-1} \int_{0}^{t} 1_{\left\{a-\varepsilon<H_{r} \leq a-\varepsilon+\varepsilon_{0}\right\}} d r-\left(\varepsilon_{0}-\varepsilon\right)^{-1} \int_{0}^{t} 1_{\left\{a<H_{r} \leq a+\varepsilon_{0}-\varepsilon\right\}} d r\right|\right] \leq \delta
$$

We get the desired result by combining the last two bounds and (1.27).
The existence of a jointly measurable modification of the process ( $L_{s}^{a}, a \geq 0, s \geq 0$ ) that satisfies the density of occupation time formula (1.25) follows from the first assertion of the proposition by standard arguments.

From now on, we will only deal with the jointly measurable modification of the local times ( $L_{s}^{a}, a \geq 0, s \geq 0$ ) given by Proposition 1.3.3. We observe that it is easy to extend the definition of these local times under the excursion measure $N$. First notice that, as an obvious consequence of the first assertion of Proposition 1.3.3, we have also for every $a \geq 0, t \geq 0$

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} E\left[\sup _{0 \leq r \leq s \leq t}\left|\varepsilon^{-1} \int_{r}^{s} 1_{\left\{a<H_{u} \leq a+\varepsilon\right\}} d u-\left(L_{s}^{a}-L_{r}^{a}\right)\right|\right]=0 \tag{1.28}
\end{equation*}
$$

Then, let $V$ be a measurable subset of $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ such that $N[V]<\infty$. For instance we may take $V=\left\{\sup _{s>0} X_{s}>\delta\right\}$ for $\delta>0$. By considering the first excursion of $X-I$ that belongs to $V$, and then using (1.28), we immediately obtain the existence under $N$ of a continuous increasing process, still denoted by ( $L_{s}^{a}, s \geq 0$ ), such that

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} N\left(1_{V} \sup _{s \leq t}\left|\varepsilon^{-1} \int_{0}^{s} 1_{\left\{a<H_{r} \leq a+\varepsilon\right\}} d r-L_{s}^{a}\right|\right)=0 \tag{1.29}
\end{equation*}
$$

The next corollary will now be a consequence of Proposition 1.1.4. We use the notation introduced before Proposition 1.2.5.

Corollary 1.3.4 For any nonnegative measurable function $F$ on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$, and every $a>0$,

$$
N\left(\int_{0}^{\sigma} d L_{s}^{a} F\left(\widehat{X}_{r \wedge s}^{(s)}, r \geq 0\right)\right)=E\left[1_{\left\{L^{-1}(a)<\infty\right\}} F\left(X_{r \wedge L^{-1}(a)}, r \geq 0\right)\right]
$$

In particular, for any nonnegative measurable function $F$ on $M_{f}\left(\mathbb{R}_{+}\right)$,

$$
N\left(\int_{0}^{\sigma} d L_{s}^{a} F\left(\rho_{s}\right)\right)=e^{-\alpha a} E^{0}\left[F\left(J_{a}\right)\right]
$$

Proof. We may assume that $F$ is bounded and continuous. Then let $h$ be a nonnegative continuous function on $\mathbb{R}_{+}$, which vanishes outside $[\delta, A]$, for some $0<\delta<A<$ $\infty$. For the first identity, it is enough to prove that

$$
N\left(\int_{0}^{\sigma} d L_{s}^{a} h(s) F\left(\widehat{X}_{r \wedge s}^{(s)}, r \geq 0\right)\right)=E\left[1_{\left\{L^{-1}(a)<\infty\right\}} h\left(L^{-1}(a)\right) F\left(X_{r \wedge L^{-1}(a)}, r \geq 0\right)\right]
$$

Notice that the mapping $s \longrightarrow\left(\widehat{X}_{t}^{(s)}, t \geq 0\right)$ is continuous except possibly on a countable set that is not charged by the measure $d L_{s}^{a}$. From (1.29), applied with $V=\{\omega, \sigma(\omega)>\delta\}$, and then Proposition 1.1.4, we get

$$
\begin{aligned}
& N\left(\int_{0}^{\sigma} d L_{s}^{a} h(s) F\left(\widehat{X}_{r \wedge s}^{(s)}, r \geq 0\right)\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} N\left(\frac{1}{\varepsilon} \int_{0}^{\sigma} d s 1_{\left\{a<H_{s} \leq a+\varepsilon\right\}} h(s) F\left(\widehat{X}_{r \wedge s}^{(s)}, r \geq 0\right)\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} E\left[\int_{a \wedge L_{\infty}}^{(a+\varepsilon) \wedge L_{\infty}} d x h\left(L^{-1}(x)\right) F\left(X_{t \wedge L^{-1}(x)}, t \geq 0\right)\right] \\
& \quad=E\left[1_{\left\{L^{-1}(a)<\infty\right\}} h\left(L^{-1}(a)\right) F\left(X_{r \wedge L^{-1}(a)}, r \geq 0\right)\right],
\end{aligned}
$$

which completes the proof of the first assertion. The second assertion follows from the first one in the same way as Proposition 1.2.5 was derived from Proposition 1.1.4.

We conclude this section with some remarks that will be useful in the applications developed below. Let $x>0$ and let $\left(\alpha_{j}, \beta_{j}\right)$, resp. $\omega_{j}, j \in J$, denote the excursion intervals, resp. the excursions of $X-I$ before time $T_{x}$. For every $a>0$, we have $P$ a.s.

$$
\begin{equation*}
L_{T_{x}}^{a}=\sum_{j \in J} L_{\sigma\left(\omega_{j}\right)}^{a}\left(\omega_{j}\right) \tag{1.30}
\end{equation*}
$$

A first inequality is easily derived by writing

$$
L_{T_{x}}^{a} \geq \int_{0}^{T_{x}} d_{s} L_{s}^{a} 1_{\left\{X_{s}>I_{s}\right\}}=\sum_{j \in J}\left(L_{\beta_{j}}^{a}-L_{\alpha_{j}}^{a}\right)=\sum_{j \in J} L_{\sigma\left(\omega_{j}\right)}^{a}\left(\omega_{j}\right)
$$

where the last equality follows from the approximations of local time. The converse inequality seems to require a different argument in our general setting. Observe that, by excursion theory and then Proposition 1.2.5,

$$
\begin{aligned}
E\left[L_{T_{x}}^{a}\right] & \leq \liminf _{k \rightarrow \infty} E\left[\frac{1}{\varepsilon_{k}} \int_{0}^{T_{x}} d s 1_{\left\{a<H_{s}<a+\varepsilon_{k}\right\}}\right] \\
& =\liminf _{k \rightarrow \infty} E\left[\sum_{j \in J} \frac{1}{\varepsilon_{k}} \int_{0}^{\sigma\left(\omega_{j}\right)} d s 1_{\left\{a<H_{s}\left(\omega_{j}\right)<a+\varepsilon_{k}\right\}}\right] \\
& =\liminf _{k \rightarrow \infty} x N\left(\frac{1}{\varepsilon_{k}} \int_{0}^{\sigma} d s 1_{\left\{a<H_{s}<a+\varepsilon_{k}\right\}}\right) \\
& =\liminf _{k \rightarrow \infty} \frac{x}{\varepsilon_{k}} \int_{a}^{a+\varepsilon_{k}} d b e^{-\alpha b} \\
& =x e^{-\alpha a}
\end{aligned}
$$

whereas Corollary 1.3 .4 (with $F=1$ ) gives $E\left[\sum_{j \in J} L_{\sigma\left(\omega_{j}\right)}^{a}\left(\omega_{j}\right)\right]=x N\left(L_{\sigma}^{a}\right)=x e^{-\alpha a}$. This readily yields (1.30).

Let us finally observe that we can extend the definition of the local times $L_{s}^{a}$ to the process $\rho$ started at a general initial value $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$. In view of forthcoming applications consider the case when $\mu$ is supported on $[0, a)$, for $a>0$. Then, the previous method can be used to construct a continuous increasing process ( $L_{s}^{a}\left(\rho^{\mu}\right), s \geq$ $0)$ such that

$$
L_{s}^{a}\left(\rho^{\mu}\right)=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_{0}^{s} d r 1_{\left\{a<H\left(\rho_{r}^{\mu}\right)<a+\varepsilon\right\}}
$$

in probability (or even in the $L^{1}$-norm). Indeed the arguments of the proof of Proposition 1.3.1 remain valid when $\rho$ is replaced by $\rho^{\mu}$, and the construction and approximation of $L_{s}^{a}\left(\rho^{\mu}\right)$ follow. Recall the notation $\tau_{0}=\inf \left\{s \geq 0: \rho_{s}^{\mu}=0\right\}$ and observe that $\tau_{0}=T_{x}$ if $x=\langle\mu, 1\rangle$. Let $\left(\alpha_{j}, \beta_{j}\right), \omega_{j}, j \in J$ be as above and set $r_{j}=H\left(k_{-I_{\alpha_{j}}} \mu\right)$. Then we have

$$
\begin{equation*}
L_{\tau_{0}}^{a}\left(\rho^{\mu}\right)=\sum_{j \in J} L_{\beta_{j}-\alpha_{j}}^{a-r_{j}}\left(\omega_{j}\right) . \tag{1.31}
\end{equation*}
$$

The proof is much similar to that of (1.30): The fact that the left side of (1.31) is greater than the right side is easy from our approximations of local time. The equality is then obtained from a first-moment argument, using Proposition 1.2.6 and Fatou's lemma to handle the left side.

### 1.4 Three applications

### 1.4.1 The Ray-Knight theorem

Recall that the $\psi$-continuous-state branching process (in short the $\psi$-CSBP) is the Markov process $\left(Y_{a}, a \geq 0\right)$ with values in $\mathbb{R}_{+}$whose transition kernels are characterized by their Laplace transform: For $\lambda>0$ and $b>a$,

$$
E\left[\exp -\lambda Y_{b} \mid Y_{a}\right]=\exp \left(-Y_{a} u_{b-a}(\lambda)\right)
$$

where $u_{t}(\lambda), t \geq 0$ is the unique nonnegative solution of the integral equation

$$
\begin{equation*}
u_{t}(\lambda)+\int_{0}^{t} \psi\left(u_{s}(\lambda)\right) d s=\lambda \tag{1.32}
\end{equation*}
$$

Theorem 1.4.1 Let $x>0$. The process $\left(L_{T_{x}}^{a}, a \geq 0\right)$ is a $\psi$-CSBP started at $x$.
Proof. First observe that $L_{T_{x}}^{a}$ is $\mathcal{H}_{a}$-measurable. This is trivial for $a=0$ since $L_{T_{x}}^{0}=x$. For $a>0$, note that, if

$$
T_{x}^{a}=\inf \left\{s \geq 0: X_{\tilde{\tau}_{s}^{a}}=-x\right\}
$$

we have

$$
\int_{0}^{T_{x}} d s 1_{\left\{a-\varepsilon<H_{s} \leq a\right\}}=\int_{0}^{T_{x}^{a}} d s 1_{\left\{a-\varepsilon<H_{\tilde{\tau}_{s}^{a}} \leq a\right\}}
$$

and the right-hand side is measurable with respect to the $\sigma$-field $\mathcal{H}_{a}$. The measurability of $L_{T_{x}}^{a}$ with respect to $\mathcal{H}_{a}$ then follows from the second convergence of Proposition 1.3.3.

We then verify that the function

$$
u_{a}(\lambda)=N\left[1-e^{-\lambda L_{\sigma}^{a}}\right] \quad(a>0), \quad u_{0}(\lambda)=\lambda
$$

solves equation (1.32). From the strong Markov property of $\rho$ under the excursion measure $N$, we get for $a>0$

$$
u_{a}(\lambda)=\lambda N\left(\int_{0}^{\sigma} d L_{s}^{a} e^{-\lambda\left(L_{\sigma}^{a}-L_{s}^{a}\right)}\right)=\lambda N\left(\int_{0}^{\sigma} d L_{s}^{a} F\left(\rho_{s}\right)\right)
$$

where, for $\mu \in M_{f}\left(\mathbb{R}_{+}\right), F(\mu)=E\left[\exp \left(-\lambda L_{\tau_{0}}^{a}\left(\rho^{\mu}\right)\right)\right]$. By Corollary 1.3.4, we can concentrate on the case when $\mu$ is supported on $[0, a)$, and then (1.31) gives

$$
\begin{aligned}
F(\mu) & =\exp \left(-\int_{0}^{\langle\mu, 1\rangle} d u N\left(1-\exp \left(-\lambda L_{\sigma}^{a-H\left(k_{-u} \mu\right)}\right)\right)\right) \\
& =\exp \left(-\int \mu(d r) N\left(1-\exp \left(-\lambda L_{\sigma}^{a-r}\right)\right)\right)
\end{aligned}
$$

Hence, using again Corollary 1.3.4,

$$
\begin{aligned}
u_{a}(\lambda) & =\lambda N\left(\int_{0}^{\sigma} d L_{s}^{a} \exp \left(-\int \rho_{s}(d r) u_{a-r}(\lambda)\right)\right) \\
& =\lambda e^{-\alpha a} E^{0}\left[\exp \left(-\int J_{a}(d r) u_{a-r}(\lambda)\right)\right] \\
& =\lambda \exp \left(-\int_{0}^{a} d r \tilde{\psi}\left(u_{a-r}(\lambda)\right)\right)
\end{aligned}
$$

It is a simple matter to verify that (1.32) follows from this last equality.
By (1.30) and excursion theory, we have

$$
\begin{equation*}
E\left[\exp \left(-\lambda L_{T_{x}}^{a}\right)\right]=\exp \left(-x N\left(1-\exp \left(-\lambda L_{\sigma}^{a}\right)\right)\right)=\exp \left(-x u_{a}(\lambda)\right) \tag{1.33}
\end{equation*}
$$

To complete the proof, it is enough to show that for $0<a<b$,

$$
\begin{equation*}
E\left[\exp \left(-\lambda L_{T_{x}}^{b}\right) \mid \mathcal{H}_{a}\right]=\exp \left(-u_{b-a}(\lambda) L_{T_{x}}^{a}\right) \tag{1.34}
\end{equation*}
$$

Recall the notation $\rho^{a}$ from Proposition 1.3.1, and denote by $\tilde{L}_{s}^{c}$ the local times of $H_{s}^{a}=H\left(\rho_{s}^{a}\right)$. From our approximations of local times, it is straightforward to verify that

$$
L_{T_{x}}^{b}=\tilde{L}_{A_{T_{x}}^{a}}^{b-a}
$$

where $A_{s}^{a}=\int_{0}^{s} d r 1_{\left\{H_{r}>a\right\}}$ as previously. Write $U=L_{T_{x}}^{a}$ to simplify notation. If $T_{r}^{a}=$ $\inf \left\{t \geq 0: l^{a}(t)>r\right\}$, we have $A_{T_{x}}^{a}=T_{U}^{a}$ (note that $l^{a}\left(A_{T_{x}}^{a}\right)=U$ by construction, and that the strong Markov property of $X$ at time $T_{x}$ implies $l^{a}(t)>l^{a}\left(A_{T_{x}}^{a}\right)$ for every $\left.t>A_{T_{x}}^{a}\right)$. Hence,

$$
E\left[\exp \left(-\lambda L_{T_{x}}^{b}\right) \mid \mathcal{H}_{a}\right]=E\left[\exp \left(-\lambda \tilde{L}_{T_{U}^{b}}^{b-a}\right) \mid \mathcal{H}_{a}\right]=E\left[\exp \left(-\lambda \tilde{L}_{T_{u}^{a}}^{b-a}\right)\right]_{u=U}
$$

where in the second equality, we use the fact that the process $\left(\tilde{L}_{T_{u}^{a}}^{b-a}, u \geq 0\right)$ is a functional of $\rho^{a}$, and is thus independent of $\mathcal{H}_{a}$ (Proposition 1.3.1), whereas $U=L_{T_{x}}^{a}$ is $\mathcal{H}_{a}$-measurable. Since $\rho^{a}$ has the same distribution as $\rho, \tilde{L}_{T_{u}^{a}}^{b-a}$ and $L_{T_{u}}^{b-a}$ also have the same law, and the desired result (1.34) follows from (1.33).

Corollary 1.4.2 For every $a \geq 0$, set

$$
v(a)=N\left(\sup _{0 \leq s \leq \sigma} H_{s}>a\right)
$$

Then,
(i) If $\int_{1}^{\infty} \frac{d u}{\psi(u)}=\infty$, we have $v(a)=\infty$ for every $a>0$.
(ii) If $\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty$, the function $(v(a), a>0)$ is determined by

$$
\int_{v(a)}^{\infty} \frac{d u}{\psi(u)}=a
$$

Proof. By the lower semi-continuity of $H$, the condition $\sup _{0 \leq s \leq \sigma} H_{s}>a$ holds iff $A_{\sigma}^{a}>0$, and our construction shows that this is the case iff $L_{\sigma}^{a}>0$. Thus,

$$
v(a)=N\left(L_{\sigma}^{a}>0\right)=\lim _{\lambda \rightarrow \infty} N\left(1-e^{-\lambda L_{\sigma}^{a}}\right)=\lim _{\lambda \rightarrow \infty} u_{a}(\lambda)
$$

with the notation of the proof of Theorem 1.4.1. From (1.32), we have

$$
\int_{u_{a}(\lambda)}^{\lambda} \frac{d u}{\psi(u)}=a
$$

and the desired result follows.

### 1.4.2 The continuity of the height process

We now use Corollary 1.4.2 to give a necessary and sufficient condition for the path continuity of the height process $H$.

Theorem 1.4.3 The process $H$ has continuous sample paths $P$ a.s. iff $\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty$.
Proof. By excursion theory, we have

$$
P\left[\sup _{0 \leq s \leq T_{x}} H_{s}>a\right]=1-\exp (-x v(a))
$$

By Corollary 1.4.2 (i), we see that $H$ cannot have continuous paths if $\int_{1}^{\infty} \frac{d u}{\psi(u)}=\infty$.
Assume that $\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty$. The previous formula and the property $v(a)<\infty$ imply that

$$
\begin{equation*}
\lim _{t \downarrow 0} H_{t}=0 \quad P \text { a.s. } \tag{1.35}
\end{equation*}
$$

Since $v(a) \downarrow 0$ as $a \uparrow \infty$, we also get that the process $H$ is locally bounded, a.s. The path continuity of $H$ will then follow from Lemma 1.2.4 if we can show that for every fixed interval $[a, a+h], h>0$, the number of upcrossings of $H$ along $[a, a+h]$ is a.s. finite on every finite time interval. Set $\gamma_{0}=0$ and define by induction for every $n \geq 1$,

$$
\begin{aligned}
& \delta_{n}=\inf \left\{t \geq \gamma_{n-1}: H_{t} \geq a+h\right\} \\
& \gamma_{n}=\inf \left\{t \geq \delta_{n}: H_{t} \leq a\right\}
\end{aligned}
$$

Both $\delta_{n}$ and $\gamma_{n}$ are $\left(\mathcal{G}_{t}\right)$-stopping times. Note that $H_{\gamma_{n}} \leq a$ by the lower semicontinuity of $H$. On the other hand, as a straightforward consequence of (1.13), we have a.s. for every $t \geq 0$,

$$
H_{\gamma_{n}+t} \leq H_{\gamma_{n}}+H_{t}^{\left(\gamma_{n}\right)}
$$

where $H^{\left(\gamma_{n}\right)}$ stands for the height process associated with $X_{t}^{\left(\gamma_{n}\right)}=X_{\gamma_{n}+t}-X_{\gamma_{n}}$. Therefore $\delta_{n+1}-\gamma_{n} \geq \kappa_{n}$, if $\kappa_{n}=\inf \left\{t \geq 0: H_{t}^{\left(\gamma_{n}\right)} \geq h\right\}$. The strong Markov
property implies that the variables $\kappa_{n}$ are i.i.d. . Furthermore, $\kappa_{n}>0$ a.s. by (1.35). It follows that $\delta_{n} \uparrow \infty$ as $n \uparrow \infty$, which completes the proof.

It is easy to see that the condition $\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty$ is also necessary and sufficient for $H$ to have continuous sample paths $N$ a.e. On the other hand, we may consider the process $\rho$ started at an arbitrary initial value $\mu \in M_{f}\left(\mathbb{R}^{d}\right)$, as defined by formula (1.11), and ask about the sample path continuity of $H\left(\rho_{s}\right)$. Clearly, the answer will be no if the support of $\mu$ is not connected. For this reason, we introduce the set $M_{f}^{0}$ which consists of all measures $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$such that $H(\mu)<\infty$ and $\operatorname{supp} \mu=[0, H(\mu)]$. By convention the zero measure also belongs to $M_{f}^{0}$.

From (1.11) and Lemma 1.2.2, it is easy to verify that the process $\rho$ started at an initial value $\mu \in M_{f}^{0}$ will remain forever in $M_{f}^{0}$, and furthermore $H\left(\rho_{s}\right)$ will have continuous sample paths a.s. Therefore, we may restrict the state space of $\rho$ to $M_{f}^{0}$. This restriction will be needed in Chapter 4.

### 1.4.3 Hölder continuity of the height process

In view of applications in Chapter 4, we now discuss the Hölder continuity properties of the height process. We assume that the condition $\int_{1}^{\infty} d u / \psi(u)<\infty$ holds so that $H$ has continuous sample paths by Theorem 1.4.3. We set

$$
\gamma=\sup \left\{r \geq 0: \lim _{\lambda \rightarrow \infty} \lambda^{-r} \psi(\lambda)=+\infty\right\}
$$

The convexity of $\psi$ implies that $\gamma \geq 1$.
Theorem 1.4.4 The height process $H$ is P-a.s. locally Hölder continuous with exponent $\alpha$ for any $\alpha \in(0,1-1 / \gamma)$, and is $P$-a.s. not locally Hölder continuous with exponent $\alpha$ if $\alpha>1-1 / \gamma$.

Proof. We rely on the following key lemma. Recall the notation $\widehat{L}^{(t)}$ for the local time at 0 of $\widehat{X}^{(t)}-\widehat{S}^{(t)}(\operatorname{cf~Section~1.2).~}$

Lemma 1.4.5 Let $t \geq 0$ and $s>0$. Then $P$ a.s.,

$$
\begin{aligned}
& H_{t+s}-\inf _{r \in[t, t+s]} H_{r}=H\left(\rho_{s}^{(t)}\right), \\
& H_{t}-\inf _{r \in[t, t+s]} H_{r}=\widehat{L}_{t \wedge R}^{(t)}
\end{aligned}
$$

where $R=\inf \left\{r \geq 0: \widehat{X}_{r}^{(t)}>-I_{s}^{(t)}\right\}(\inf \emptyset=\infty)$.
Proof. From (1.13), we get, a.s. for every $r>0$,

$$
\begin{equation*}
H_{t+r}=H\left(k_{-I_{r}^{(t)}} \rho_{t}\right)+H\left(\rho_{r}^{(t)}\right) \tag{1.36}
\end{equation*}
$$

From this it follows that

$$
\inf _{r \in[t, t+s]} H_{r}=H\left(k_{-I_{s}^{(t)}} \rho_{t}\right)
$$

and the minimum is indeed attained at the (a.s. unique) time $v \in[t, t+s]$ such that $X_{v}=I_{t+s}^{t}$. The first assertion of the lemma now follows by combining the last equality with (1.36) written with $r=s$.

Let us turn to the second assertion. If $I_{t} \geq I_{t+s}^{t}$, then on one hand $X_{v}=I_{v}$ and $\inf _{r \in[t, t+s]} H_{r}=H_{v}=0$, on the other hand, $R=\infty$, and the second assertion reduces
to $H_{t}=\widehat{L}_{t}^{(t)}$ which is the definition of $H_{t}$. Therefore we can assume that $I_{t}<I_{t+s}^{t}$. Let

$$
u=\sup \left\{r \in[0, t]: X_{r-}<I_{t+s}^{t}\right\} .
$$

As in the proof of Proposition 1.2.3, we have

$$
H_{u}=H\left(k_{-I_{s}^{(t)}} \rho_{t}\right)=\inf _{r \in[t, t+s]} H_{r} .
$$

On the other hand, the construction of the height process shows that the equality $H_{r}=\widehat{L}_{t}^{(t)}-\widehat{L}_{t-r}^{(t)}$ holds simultaneously for all $r \in[0, t]$ such that $X_{r-} \leq I_{t}^{r}$ (cf Lemma 1.2.1). In particular for $r=u$ we get

$$
H_{t}-\inf _{r \in[t, t+s]} H_{r}=H_{t}-H_{u}=\widehat{L}_{t}^{(t)}-\left(\widehat{L}_{t}^{(t)}-\widehat{L}_{t-u}^{(t)}\right)=\widehat{L}_{t-u}^{(t)} .
$$

To complete the proof, simply note that we have $t-u=R$ on the event $\left\{I_{t}<I_{t+s}^{t}\right\}$
To simplify notation we set $\varphi(\lambda)=\lambda / \psi^{-1}(\lambda)$. The right-continuous inverse $L^{-1}$ of $L$ is a subordinator with Laplace exponent $\varphi$ : See Theorem VII. 4 (ii) in [5], and note that the constant $c$ in this statement is equal to 1 under our normalization of local time (compare with Lemma 1.1.2).

Lemma 1.4.6 For every $t \geq 0, s>0$ and $q>0$,

$$
E\left[\left|H_{t+s}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] \leq C_{q} \varphi(1 / s)^{-q}
$$

and

$$
E\left[\left|H_{t}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] \leq C_{q} \varphi(1 / s)^{-q}
$$

where $C_{q}=e \Gamma(q+1)$ is a finite constant depending only on $q$.
Proof. Recall that $H\left(\rho_{s}\right)=H_{s} \stackrel{(\mathrm{~d})}{=} L_{s}$. From Lemma 1.4.5 we have

$$
E\left[\left|H_{t+s}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right]=E\left[L_{s}^{q}\right]=q \int_{0}^{+\infty} x^{q-1} P\left[L_{s}>x\right] d x
$$

However,

$$
P\left[L_{s}>x\right]=P\left[s>L^{-1}(x)\right] \leq e E\left[\exp \left(-L^{-1}(x) / s\right)\right]=e \exp (-x \varphi(1 / s))
$$

Thus

$$
E\left[\left|H_{t+s}-\inf _{[t, t+s]} H\right|^{q}\right] \leq e q \int_{0}^{+\infty} x^{q-1} \exp (-x \varphi(1 / s)) d x=C_{q} \varphi(1 / s)^{-q}
$$

This completes the proof of the first assertion.
In order to prove the second one, first note that $I_{s}^{(t)}$ is independent of $\mathcal{G}_{t}$ and therefore also of the time-reversed process $\widehat{X}^{(t)}$. Writing $\tau_{a}=\inf \left\{r \geq 0: S_{r}>a\right\}$, we get from the second assertion of Lemma 1.4.5

$$
E\left[\left|H_{t}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] \leq \int_{[0,+\infty)} P\left[-I_{s} \in d a\right] E\left[L_{\tau_{a}}^{q}\right] .
$$

Note that

$$
E\left[L_{\tau_{a}}^{q}\right]=q \int_{0}^{+\infty} x^{q-1} P\left[L_{\tau_{a}}>x\right] d x
$$

and that $P\left[L_{\tau_{a}}>x\right]=P\left[S_{L^{-1}(x)}<a\right]$. It follows that

$$
E\left[\left|H_{t}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] \leq q \int_{[0,+\infty)} P\left[-I_{s} \in d a\right] \int_{0}^{+\infty} d x x^{q-1} P\left[S_{L^{-1}(x)}<a\right] .
$$

An integration by parts leads to

$$
E\left[\left|H_{t}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] \leq q \int_{0}^{+\infty} d x x^{q-1} \int_{[0,+\infty)} P\left[S_{L^{-1}(x)} \in d b\right] P\left[-I_{s}>b\right]
$$

However

$$
P\left[-I_{s}>b\right]=P\left[T_{b}<s\right] \leq e E\left[\exp \left(-T_{b} / s\right)\right]=e \exp \left(-b \psi^{-1}(1 / s)\right)
$$

since we know ([5] Theorem VII.1) that $\left(T_{b}, b \geq 0\right)$ is a subordinator with exponent $\psi^{-1}$. Recalling Lemma 1.1.2, we get

$$
\begin{aligned}
E\left[\left|H_{t}-\inf _{r \in[t, t+s]} H_{r}\right|^{q}\right] & \leq e q \int_{0}^{+\infty} d x x^{q-1} E\left[\exp \left(-\psi^{-1}(1 / s) S_{L^{-1}(x)}\right)\right] \\
& =e q \int_{0}^{+\infty} d x x^{q-1} \exp \left(-x /\left(s \psi^{-1}(1 / s)\right)\right) \\
& =C_{q} \varphi(1 / s)^{-q}
\end{aligned}
$$

This completes the proof of Lemma 1.4.6.
Proof of Theorem 1.4.4. From Lemma 1.4.6 and an elementary inequality, we get for every $t \geq 0, s>0$ and $q>0$

$$
E\left[\left|H_{t+s}-H_{t}\right|^{q}\right] \leq 2^{q+1} C_{q} \varphi(1 / s)^{q}
$$

Let $\alpha \in(0,1-1 / \gamma)$. Then $(1-\alpha)^{-1}<\gamma$ and thus $\lambda^{-(1-\alpha)^{-1}} \psi(\lambda)$ tends to $\infty$ as $\lambda \rightarrow \infty$. It easily follows that $\lambda^{\alpha-1} \psi^{-1}(\lambda)$ tends to 0 and so $\lambda^{-\alpha} \varphi(\lambda)$ tends to $\infty$ as $\lambda \rightarrow \infty$. The previous bound then yields the existence of a constant $C$ depending on $q$ and $\alpha$ such that for every $t \geq 0$ and $s \in(0,1]$,

$$
E\left[\left|H_{t+s}-H_{t}\right|^{q}\right] \leq C s^{q \alpha}
$$

The classical Kolmogorov lemma gives the first assertion of the theorem.
To prove the second assertion, observe that for every $\alpha>0$ and $A>0$,

$$
P\left[H_{s}<A s^{\alpha}\right]=P\left[L_{s}<A s^{\alpha}\right]=P\left[s<L^{-1}\left(A s^{\alpha}\right)\right]
$$

Then use the elementary inequality

$$
P\left[s<L^{-1}\left(A s^{\alpha}\right)\right] \leq \frac{e}{e-1} E\left[1-\exp \left(-L^{-1}\left(A s^{\alpha}\right) / s\right)\right]
$$

which leads to

$$
P\left[H_{s}<A s^{\alpha}\right] \leq \frac{e}{e-1}\left(1-\exp \left(-A s^{\alpha} \varphi(1 / s)\right)\right)
$$

If $\alpha>1-1 / \gamma$, we can find a sequence $\left(s_{n}\right)$ decreasing to zero such $s_{n}^{\alpha} \varphi\left(1 / s_{n}\right)$ tends to 0 . Thus, for any $A>0$

$$
\lim _{n \rightarrow \infty} P\left[H_{s_{n}}<A s_{n}^{\alpha}\right]=0
$$

and it easily follows that $\limsup _{s \rightarrow 0} s^{-\alpha} H_{s}=\infty, P$ a.s., which completes the proof.

## Chapter 2

## Convergence of Galton-Watson trees

### 2.1 Preliminaries

Our goal in this chapter is to study the convergence in distribution of Galton-Watson trees, under the assumption that the associated Galton-Watson processes, suitably rescaled, converge in distribution to a continuous-state branching process. To give a precise meaning to the convergence of trees, we will code Galton-Watson trees by a discrete height process, and we will establish the convergence of these (rescaled) discrete processes to the continuous height process of the previous chapter. We will also prove that similar convergences hold when the discrete height processes are replaced by the contour processes of the trees.

Let us introduce the basic objects considered in this chapter. For every $p \geq 1$, let $\mu_{p}$ be a subcritical or critical offspring distribution. That is, $\mu_{p}$ is a probability distribution on $\mathbb{Z}_{+}=\{0,1, \ldots\}$ such that

$$
\sum_{k=0}^{\infty} k \mu_{p}(k) \leq 1
$$

We systematically exclude the trivial cases where $\mu_{p}(1)=1$ or $\mu_{p}(0)=1$. We also define another probability measure $\nu_{p}$ on $\{-1,0,1, \ldots\}$ by setting $\nu_{p}(k)=\mu_{p}(k+1)$ for every $k \geq-1$.

We denote by $V^{p}=\left(V_{k}^{p}, k=0,1,2, \ldots\right)$ a discrete-time random walk on $\mathbb{Z}$ with jump distribution $\nu_{p}$ and started at 0 . We also denote by $Y^{p}=\left(Y_{k}^{p}, k=0,1,2, \ldots\right)$ a Galton-Watson branching process with offspring distribution $\mu_{p}$ started at $Y_{0}^{p}=p$.

Finally, we consider a Lévy process $X=\left(X_{t}, t \geq 0\right)$ started at the origin and satisfying assumptions (H1) and (H2) of Chapter 1. As in Chapter 1, we write $\psi$ for the Laplace exponent of $X$. We denote by $Y=\left(Y_{t}, t \geq 0\right)$ a $\psi$-continuous-state branching process started at $Y_{0}=1$.

The following variant of a result due to Grimvall [21] plays an important role in our approach. Unless otherwise specified the convergence in distribution of processes is in the functional sense, that is in the sense of the weak convergence of the laws of the processes on the Skorokhod space $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. We will use the notation $\xrightarrow{(\mathrm{fd})}$ to indicate weak convergence of finite-dimensional marginals.

For $a \in \mathbb{R},[a]$ denotes the integer part of $a$.

Theorem 2.1.1 Let $\left(\gamma_{p},, p=1,2, \ldots\right)$ be a nondecreasing sequence of positive integers converging to $\infty$. The convergence in distribution

$$
\begin{equation*}
\left(p^{-1} Y_{\left[\gamma_{p} t\right]}^{p}, t \geq 0\right) \xrightarrow[p \rightarrow \infty]{\stackrel{(\mathrm{d})}{\longrightarrow}}\left(Y_{t}, t \geq 0\right) \tag{2.1}
\end{equation*}
$$

holds if and only if

$$
\begin{equation*}
\left(p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(X_{t}, t \geq 0\right) . \tag{2.2}
\end{equation*}
$$

Proof. By standard results on the convergence of triangular arrays (see e.g. Theorem 2.7 in Skorokhod [46]), the functional convergence (2.2) holds iff

$$
\begin{equation*}
p^{-1} V_{p \gamma_{p}}^{p} \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\rightarrow}} X_{1} . \tag{2.3}
\end{equation*}
$$

Fix any sequence $p_{1}<p_{2}<\cdots<p_{k}<\cdots$ such that $\gamma_{p_{1}}<\gamma_{p_{2}}<\cdots$. If $j=\gamma_{p_{k}}$ for some $k \geq 1$, set $c_{j}=p_{k}, V^{(j)}=V^{p_{k}}$ and let $\theta_{j}$ be the probability measure on $\mathbb{R}$ defined by $\theta_{j}\left(\frac{n}{c_{j}}\right)=\nu_{p_{k}}(n)$ for every integer $n \geq-1$. Then (2.3) is equivalent to saying that

$$
\frac{1}{c_{j}} V_{j c_{j}}^{(j)} \xrightarrow[j \rightarrow \infty]{(\mathrm{d})} X_{1}
$$

for any choice of the sequence $p_{1}<p_{2}<\cdots$. Equivalently the convolutions $\left(\theta_{j}\right)^{* j c_{j}}$ converge weakly to the law of $X_{1}$. By Theorems 3.4 and 3.1 of Grimvall [21], this property holds iff the convergence (2.1) holds along the sequence $\left(p_{k}\right)$. (Note that condition (b) in Theorem 3.4 of [21] is automatically satisfied here since we restrict our attention to the (sub)critical case.) This completes the proof.

### 2.2 The convergence of finite-dimensional marginals

From now on, we suppose that assumption (H3) holds in addition to (H1) and (H2). Thus we can consider the height process $H=\left(H_{t}, t \geq 0\right)$ of Chapter 1.

For every $p \geq 1$, let $H^{p}=\left(H_{k}^{p}, k \geq 0\right)$ be the discrete height process associated with a sequence of independent Galton-Watson trees with offspring distribution $\mu_{p}$ (cf Section 0.2). As was observed in Section 0.2, we may and will assume that the processes $H^{p}$ and $V^{p}$ are related by the formula

$$
\begin{equation*}
H_{k}^{p}=\operatorname{Card}\left\{j \in\{0,1, \ldots, k-1\}: V_{j}^{p}=\inf _{j \leq l \leq k} V_{l}^{p}\right\} \tag{2.4}
\end{equation*}
$$

The following theorem sharpens a result of [32].
Theorem 2.2.1 Under either of the convergences (2.1) or (2.2), we have also

$$
\begin{equation*}
\left(\frac{1}{\gamma_{p}} H_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{(\mathrm{fd})}\left(H_{t}, t \geq 0\right) . \tag{2.5}
\end{equation*}
$$

Proof. Let $f_{0}$ be a truncation function, that is a bounded continuous function from $\mathbb{R}$ into $\mathbb{R}$ such that $f_{0}(x)=x$ for every $x$ belonging to a neighborhood of 0 . By standard results on the convergence of rescaled random walks (see e.g. Theorem II.3.2 in [23]), the convergence (2.2) holds iff the following three conditions are satisfied:

$$
\begin{equation*}
\lim _{p \rightarrow \infty} p \gamma_{p} \sum_{k=-1}^{\infty} f_{0}\left(\frac{k}{p}\right) \nu_{p}(k)=-\alpha+\int_{0}^{\infty}\left(f_{0}(r)-r\right) \pi(d r) \tag{C1}
\end{equation*}
$$

$$
\begin{align*}
& \lim _{p \rightarrow \infty} p \gamma_{p} \sum_{k=-1}^{\infty} f_{0}\left(\frac{k}{p}\right)^{2} \nu_{p}(k)=2 \beta+\int_{0}^{\infty} f_{0}(r)^{2} \pi(d r)  \tag{C2}\\
& \lim _{p \rightarrow \infty} p \gamma_{p} \sum_{k=-1}^{\infty} h\left(\frac{k}{p}\right) \nu_{p}(k)=\int_{0}^{\infty} h(r) \pi(d r) \tag{C3}
\end{align*}
$$

for any bounded continuous function $h$ on $\mathbb{R}$ that vanishes on a neighborhood of 0 .
By (2.4) and time-reversal, $H_{k}^{p}$ has the same distribution as

$$
\Lambda_{k}^{(p)}=\operatorname{Card}\left\{j \in\{1, \ldots, k\}: V_{j}^{p}=\sup _{0 \leq l \leq j} V_{l}^{p}\right\} .
$$

Without loss of generality, the Skorokhod representation theorem allows us to assume that the convergence

$$
\begin{equation*}
\left(p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{\longrightarrow}\left(X_{t}, t \geq 0\right) \tag{2.6}
\end{equation*}
$$

holds a.s. in the sense of Skorokhod's topology. Suppose we can prove that for every $t>0$,

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \gamma_{p}^{-1} \Lambda_{\left[p \gamma_{p} t\right]}^{(p)}=L_{t} \tag{2.7}
\end{equation*}
$$

in probability. (Here $L=\left(L_{t}, t \geq 0\right)$ is the local time process of $X-S$ at level 0 as in Chapter 1.) Then a simple time-reversal argument implies that $\gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}$ also converges in probability to $\widehat{L}_{t}^{(t)}=H_{t}$, with the notation of Chapter 1 . Therefore the proof of Theorem 2.2.1 reduces to showing that (2.7) holds.

We first consider the case where $\int_{(0,1)} r \pi(d r)=\infty$. We introduce the stopping times $\left(\tau_{k}^{p}\right)_{k \geq 0}$ defined recursively as follows:

$$
\begin{aligned}
\tau_{0}^{p} & =0 \\
\tau_{m+1}^{p} & =\inf \left\{n>\tau_{m}^{p}: V_{n}^{p} \geq V_{\tau_{m}^{p}}^{p}\right\}
\end{aligned}
$$

Conditionally on the event $\left\{\tau_{m}^{p}<\infty\right\}$, the random variable $1_{\left\{\tau_{m+1}^{p}<\infty\right\}}\left(V_{\tau_{m+1}^{p}}^{p}-V_{\tau_{m}^{p}}^{p}\right)$ is independent of the past of $V^{p}$ up to time $\tau_{m}^{p}$ and has the same law as $1_{\left\{\tau_{1}^{p}<\infty\right\}} V_{\tau_{1}^{p}}^{p}$. Also recall the classical equality (cf (5.4) in [32]):

$$
\begin{equation*}
P\left[\tau_{1}^{p}<\infty, V_{\tau_{1}^{p}}^{p}=j\right]=\nu_{p}([j, \infty)), \quad j \geq 0 \tag{2.8}
\end{equation*}
$$

For every $u>\delta>0$, set:

$$
\begin{aligned}
\kappa(\delta, u)= & \int_{0}^{\infty} \pi(d r) \int_{0}^{r} d x 1_{(\delta, u]}(x)=\int_{0}^{\infty} \pi(d r)\left((r-\delta)^{+} \wedge(u-\delta)\right) \\
\kappa_{p}(\delta, u)= & \frac{\sum_{p \delta<j \leq p u} \nu_{p}([j, \infty))}{\sum_{j \geq 0} \nu_{p}([j, \infty))}=P\left[p \delta<V_{\tau_{1}^{p}}^{p} \leq p u \mid \tau_{1}^{p}<\infty\right] \\
L_{t}^{\delta, u}= & \operatorname{Card}\left\{s \leq t: \Delta S_{s} \in(\delta, u]\right\} \\
l_{k}^{p, \delta, u}= & \operatorname{Card}\left\{j<k: \bar{V}_{j}^{p}+p \delta<V_{j+1}^{p} \leq \bar{V}_{j}^{p}+p u\right\}
\end{aligned}
$$

where $\bar{V}_{j}^{p}=\sup \left\{V_{i}^{p}, 0 \leq i \leq j\right\}$. Note that $\kappa(\delta, u) \uparrow \infty$ as $\delta \downarrow 0$, by our assumption $\int_{(0,1)} r \pi(d r)=\infty$. From the a.s. convergence of the processes $p^{-1} V_{\left[p \gamma_{p} t\right]}$, we have

$$
\begin{equation*}
\lim _{p \rightarrow \infty} l_{\left[p \gamma_{p} t\right]}^{p, \delta, u}=L_{t}^{\delta, u}, \quad \text { a.s. } \tag{2.9}
\end{equation*}
$$

(Note that $P\left[\Delta S_{s}=a\right.$ for some $\left.s>0\right]=0$ for every fixed $a>0$, by (1.3).) By applying excursion theory to the process $X-S$ and using formula (1.3), one easily gets for every $u>0$

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \kappa(\delta, u)^{-1} L_{t}^{\delta, u}=L_{t}, \quad \text { a.s. } \tag{2.10}
\end{equation*}
$$

We claim that we have also

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \gamma_{p} \kappa_{p}(\delta, u)=\int_{0}^{\infty}\left((r-\delta)^{+} \wedge(u-\delta)\right) \pi(d r)=\kappa(\delta, u) \tag{2.11}
\end{equation*}
$$

To get this convergence, first apply (C3) to the function $h(x)=(x-\delta)^{+} \wedge(u-\delta)$. It follows that

$$
\lim _{p \rightarrow \infty} p \gamma_{p} \sum_{k=-1}^{\infty} \nu_{p}(k)\left(\left(\frac{k}{p}-\delta\right)^{+} \wedge(u-\delta)\right)=\kappa(\delta, u)
$$

On the other hand, it is elementary to verify that

$$
\left|p \gamma_{p} \sum_{k=-1}^{\infty} \nu_{p}(k)\left(\left(\frac{k}{p}-\delta\right)^{+} \wedge(u-\delta)\right)-\gamma_{p} \sum_{p \delta<j \leq p u} \nu_{p}([j, \infty))\right| \leq \gamma_{p} \sum_{k \geq \delta p} \nu_{p}(k)
$$

and the right-hand side tends to 0 by (C3). Thus we get

$$
\lim _{p \rightarrow \infty} \gamma_{p} \sum_{p \delta<j \leq p u} \nu_{p}([j, \infty))=\kappa(\delta, u) .
$$

Furthermore, as a simple consequence of (C1) and the (sub)criticality of $\mu_{p}$, we have also

$$
\begin{equation*}
\sum_{j=0}^{\infty} \nu_{p}([j, \infty))=1+\sum_{k=-1}^{\infty} k \nu_{p}(k) \underset{p \rightarrow \infty}{\longrightarrow} 1 \tag{2.12}
\end{equation*}
$$

(This can also be obtained from (2.8) and the weak convergence (2.2).) Our claim (2.11) now follows.

Finally, we can also obtain a relation between $l_{k}^{p, \delta, u}$ and $\Lambda_{\left[p \gamma_{p} t\right]}^{(p)}$. Simply observe that conditional on $\left\{\tau_{k}^{p}<\infty\right\}$, $l_{\tau_{k}^{p}}^{p, \delta, u}$ is the sum of $k$ independent Bernoulli variables with parameter $\kappa_{p}(\delta, u)$. Fix an integer $A>0$ and set $A_{p}=\gamma_{p} A+1$. From Doob's inequality, we easily get (see [32], p. 249 for a similar estimate)

$$
E\left[\sup _{0 \leq j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{\gamma_{p}}\left(\Lambda_{j}^{(p)}-\kappa_{p}(\delta, u)^{-1} l_{j}^{p, \delta, u}\right)\right|^{2}\right] \leq \frac{8(A+1)}{\gamma_{p}} \kappa_{p}(\delta, u)^{-1}
$$

Hence, using (2.11), we have

$$
\begin{equation*}
\limsup _{p \rightarrow \infty} E\left[\sup _{j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{\gamma_{p}}\left(\Lambda_{j}^{(p)}-\kappa_{p}(\delta, u)^{-1} l_{j}^{p, \delta, u}\right)\right|^{2}\right] \leq \frac{8(A+1)}{\kappa(\delta, u)} \tag{2.13}
\end{equation*}
$$

To complete the proof, let $\varepsilon>0$ and first choose $A$ large enough so that $P\left[L_{t} \geq\right.$ $A-3 \varepsilon]<\varepsilon$. If $u>0$ is fixed, we can use (2.10) and (2.13) to pick $\delta>0$ small enough and then $p_{0}=p_{0}(\delta)$ so that

$$
\begin{equation*}
P\left[\left|\kappa(\delta, u)^{-1} L_{t}^{\delta, u}-L_{t}\right|>\varepsilon\right]<\varepsilon \tag{2.14}
\end{equation*}
$$

and

$$
\begin{equation*}
P\left[\sup _{j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{\gamma_{p}}\left(\Lambda_{j}^{(p)}-\kappa_{p}(\delta, u)^{-1} l_{j}^{p, \delta, u}\right)\right|>\varepsilon\right]<\varepsilon, \quad \text { if } p \geq p_{0} \tag{2.15}
\end{equation*}
$$

From (2.9) and (2.11), we can also find $p_{1}(\delta)$ so that for every $p \geq p_{1}$,

$$
\begin{equation*}
P\left[\left|\frac{1}{\gamma_{p} \kappa_{p}(\delta, u)} l_{\left[p \gamma_{p} t\right]}^{p, \delta, u}-\kappa(\delta, u)^{-1} L_{t}^{\delta, u}\right|>\varepsilon\right]<\varepsilon . \tag{2.16}
\end{equation*}
$$

By combining the previous estimates (2.14), (2.15) and (2.16), we get for $p \geq p_{0} \vee p_{1}$

$$
\begin{equation*}
P\left[\left|\frac{1}{\gamma_{p}} \Lambda_{\left[p \gamma_{p} t\right]}^{(p)}-L_{t}\right|>3 \varepsilon\right] \leq 3 \varepsilon+P\left[\left[p \gamma_{p} t\right]>\tau_{A_{p}}^{p}\right] \tag{2.17}
\end{equation*}
$$

Furthermore, by using (2.15) and then (2.14) and (2.16), we have for $p$ sufficiently large,

$$
\begin{aligned}
P\left[\tau_{A_{p}}^{p}<\left[p \gamma_{p} t\right]\right] & \leq \varepsilon+P\left[\frac{1}{\gamma_{p} \kappa_{p}(\delta, u)} l_{\left[p \gamma_{p} t\right]}^{p, \delta, u} \geq A-\varepsilon\right] \\
& \leq 3 \varepsilon+P\left[L_{t} \geq A-3 \varepsilon\right] \\
& \leq 4 \varepsilon
\end{aligned}
$$

from our choice of $A$. Combining this estimate with (2.17) completes the proof of (2.7) in the case $\int_{0}^{1} r \pi(d r)=\infty$.

It remains to treat the case where $\int_{0}^{1} r \pi(d r)<\infty$. In that case, (H3) implies that $\beta>0$, and we know from (1.4) that

$$
L_{t}=\frac{1}{\beta} m\left(\left\{S_{s} ; s \leq t\right\}\right) .
$$

Furthermore, (1.3) and the assumption $\int_{0}^{1} r \pi(d r)<\infty$ imply that for any $t>0$,

$$
\operatorname{Card}\left\{s \in[0, t] ; \Delta S_{s}>0\right\}<\infty, \quad \text { a.s. }
$$

For every $\delta>0$ and $t \geq 0$, we set

$$
\widetilde{S}_{t}^{\delta}=S_{t}-\sum_{s \in[0, t]} \mathbf{1}_{(\delta, \infty)}\left(\Delta S_{s}\right) \Delta S_{s}
$$

By the previous remarks, we have a.s. for $\delta$ small enough,

$$
\begin{equation*}
\widetilde{S}_{t}^{\delta}=S_{t}-\sum_{s \in[0, t]} \Delta S_{s}=m\left(\left\{S_{s} ; s \leq t\right\}\right)=\beta L_{t} \tag{2.18}
\end{equation*}
$$

Let us use the same notation $\tau_{m}^{p}, \bar{V}_{j}^{p}$ as in the case $\int_{0}^{1} r \pi(d r)=\infty$, and also set for any $m \geq 1$,

$$
d_{m}^{p}=\mathbf{1}_{\left\{\tau_{m}^{p}<\infty\right\}}\left(V_{\tau_{m}^{p}}^{p}-V_{\tau_{m-1}^{p}}^{p}\right)
$$

and

$$
\widetilde{S}_{m}^{\delta, p}=\sum_{n \geq 1} d_{n}^{p} \mathbf{1}_{\left\{d_{n}^{p} \leq p \delta, \tau_{n}^{p} \leq m\right\}}
$$

The convergence (2.6) implies that

$$
\begin{equation*}
\left(\frac{1}{p} \bar{V}_{\left[p \gamma_{p} s\right]}^{p}, s \geq 0\right) \underset{p \rightarrow \infty}{\longrightarrow}\left(S_{s}, s \geq 0\right), \quad \text { a.s. } \tag{2.19}
\end{equation*}
$$

and, for every $t \geq 0$,

$$
\frac{1}{p} \sum_{n \geq 1} d_{n}^{p} \mathbf{1}_{\left\{d_{n}^{p}>p \delta, \tau_{n}^{p} \leq\left[p \gamma_{p} t\right]\right\}} \underset{p \rightarrow \infty}{\longrightarrow} \sum_{s \in[0, t]} \mathbf{1}_{(\delta, \infty)}\left(\Delta S_{s}\right) \Delta S_{s}, \quad \text { a.s. }
$$

Thus we have

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \frac{1}{p} \widetilde{S}_{\left[p \gamma_{p} t\right]}^{\delta, p}=\widetilde{S}_{t}^{\delta} \quad \text { a.s. } \tag{2.20}
\end{equation*}
$$

The desired convergence (2.7) is then a consequence of (2.18), (2.20) and the following result: For every $\varepsilon>0$,

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \limsup _{p \rightarrow \infty} P\left[\left|\frac{1}{p} \widetilde{S}_{\left[p \gamma_{p} t\right]}^{\delta, p}-\frac{\beta}{\gamma_{p}} \Lambda_{\left[p \gamma_{p} t\right]}^{(p)}\right|>\varepsilon\right]=0 \tag{2.21}
\end{equation*}
$$

To prove (2.21), set

$$
\begin{aligned}
& \alpha_{1}(p, \delta)=E\left[d_{1}^{p} \mathbf{1}_{\left\{d_{1}^{p} \leq \delta p\right\}} \mid \tau_{1}^{p}<\infty\right] \\
& \alpha_{2}(p, \delta)=E\left[\left(d_{1}^{p}\right)^{2} \mathbf{1}_{\left\{d_{1}^{p} \leq \delta p\right\}} \mid \tau_{1}^{p}<\infty\right] .
\end{aligned}
$$

Observe that

$$
E\left[\left(d_{1}^{p} \mathbf{1}_{\left\{d_{1}^{p} \leq \delta p\right\}}-\alpha_{1}(p, \delta)\right)^{2} \mid \tau_{1}^{p}<\infty\right] \leq \alpha_{2}(p, \delta)
$$

Let $A>0$ be an integer and let $A_{p}=\gamma_{p} A+1$ as above. By Doob's inequality,

$$
E\left[\sup _{\substack{1 \leq m \leq A_{p} \\ \tau_{m}^{p}<\infty}}\left|\widetilde{S}_{\tau_{m}^{p}}^{\delta, p}-m \alpha_{1}(p, \delta)\right|^{2}\right] \leq 4 A_{p} \alpha_{2}(p, \delta)
$$

Since

$$
\sup _{\substack{1 \leq m \leq A_{p} \\ \tau_{m}^{p}<\infty}}\left|\widetilde{S}_{\tau_{m}^{p}}^{\delta, p}-m \alpha_{1}(p, \delta)\right|=\sup _{1 \leq j \leq \tau_{A_{p}}^{p}}\left|\widetilde{S}_{j}^{\delta, p}-\alpha_{1}(p, \delta) \Lambda_{j}^{(p)}\right|
$$

we have

$$
\begin{equation*}
E\left[\sup _{0 \leq j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{p} \widetilde{S}_{j}^{\delta, p}-\frac{\alpha_{1}(p, \delta)}{p} \Lambda_{j}^{(p)}\right|^{2}\right] \leq \frac{4 A_{p}}{p^{2}} \alpha_{2}(p, \delta) \tag{2.22}
\end{equation*}
$$

We now claim that

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \frac{\gamma_{p}}{p} \alpha_{1}(p, \delta)=\beta+\frac{1}{2} \int_{(0, \infty)}(r \wedge \delta)^{2} \pi(d r) \underset{\delta \rightarrow 0}{\longrightarrow} \beta \tag{2.23}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \limsup _{p \rightarrow \infty} \frac{\gamma_{p}}{p^{2}} \alpha_{2}(p, \delta)=0 \tag{2.24}
\end{equation*}
$$

To verify (2.23), note that, by (2.8),

$$
\begin{equation*}
\frac{\gamma_{p}}{p} \alpha_{1}(p, \delta)=\frac{\gamma_{p}}{p} \frac{\sum_{j=0}^{[\delta p]} j \nu_{p}([j, \infty))}{\sum_{j=0}^{\infty} \nu_{p}([j, \infty))}=\frac{p \gamma_{p}}{2} \frac{\sum_{k=0}^{\infty} \nu_{p}(k)\left(\frac{k}{p} \wedge \frac{[\delta p]}{p}\right)\left(\frac{k}{p} \wedge \frac{[\delta p]}{p}+\frac{1}{p}\right)}{\sum_{j=0}^{\infty} \nu_{p}([j, \infty))} \tag{2.25}
\end{equation*}
$$

We now apply (C1) and (C2) with the truncation function $f_{0}(x)=(x \wedge \delta) \vee(-\delta)$. Multiplying by $p^{-1}$ the convergence in (C1) and adding the one in (C2), we get

$$
\lim _{p \rightarrow \infty} p \gamma_{p} \sum_{k=0}^{\infty} \nu_{p}(k)\left(\frac{k}{p} \wedge \delta\right)\left(\frac{k}{p} \wedge \delta+\frac{1}{p}\right)=2 \beta+\int_{(0, \infty)}(r \wedge \delta)^{2} \pi(d r)
$$

Comparing with (2.25) and using (2.12), we immediately get (2.23). The proof of (2.24) is analogous.

By (2.22) and an elementary inequality, we have

$$
E\left[\sup _{0 \leq j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{p} \widetilde{S}_{j}^{\delta, p}-\frac{\beta}{\gamma_{p}} \Lambda_{j}^{(p)}\right|^{2}\right] \leq \frac{8 A_{p}}{p^{2}} \alpha_{2}(p, \delta)+2\left(\frac{A_{p}}{\gamma_{p}}\right)^{2}\left(\beta-\frac{\gamma_{p}}{p} \alpha_{1}(p, \delta)\right)^{2}
$$

Thus, (2.23) and (2.24) imply that for any $A>0$

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \limsup _{p \rightarrow \infty} E\left[\sup _{0 \leq j \leq \tau_{A_{p}}^{p}}\left|\frac{1}{p} \widetilde{S}_{j}^{\delta, p}-\frac{\beta}{\gamma_{p}} \Lambda_{j}^{(p)}\right|^{2}\right]=0 \tag{2.26}
\end{equation*}
$$

It follows that

$$
\lim _{\delta \rightarrow 0} \limsup _{p \rightarrow \infty} P\left[\left|\frac{1}{p} \widetilde{S}_{\left[p \gamma_{p} t\right]}^{\delta, p}-\frac{\beta}{\gamma_{p}} \Lambda_{\left[p \gamma_{p} t\right]}^{(p)}\right|>\varepsilon\right] \leq \limsup _{p \rightarrow \infty} P\left[\tau_{A_{p}}^{p}<\left[p \gamma_{p} t\right]\right]
$$

However,

$$
P\left[\tau_{A_{p}}^{p}<\left[p \gamma_{p} t\right]\right] \leq P\left[\frac{1}{p} \bar{V}_{\left[p \gamma_{p} t\right]}^{p} \geq \frac{1}{p} \tilde{S}_{\tau_{A_{p}}^{p}}^{\delta, p}, \tau_{A_{p}}^{p}<\infty\right]
$$

and by (2.26) the right side is bounded above for $p$ large by $P\left[\frac{1}{p} \bar{V}_{\left[p \gamma_{p} t\right]}^{p} \geq \frac{\beta}{\gamma_{p}} A_{p}-1\right]+\varepsilon_{\delta}$, where $\varepsilon_{\delta} \rightarrow 0$ as $\delta \rightarrow 0$. In view of (2.19), this is enough to conclude that

$$
\lim _{A \rightarrow \infty} \limsup _{p \rightarrow \infty} P\left[\tau_{A_{p}}^{p}<\left[p \gamma_{p} t\right]\right]=0
$$

and the desired result (2.21) follows. This completes the proof of (2.7) and of Theorem 2.2.1.

### 2.3 The functional convergence

Our goal is now to discuss conditions that ensure that the convergence of Theorem 2.2.1 holds in a functional sense. We assume that the function $\psi$ satisfies the condition

$$
\begin{equation*}
\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty \tag{2.27}
\end{equation*}
$$

By Theorem 1.4.3, this implies that the height process ( $H_{t}, t \geq 0$ ) has continuous sample paths. On the other hand, if this condition does not hold, the paths of the height process do not belong to any of the usual functional spaces.

For every $p \geq 1$, we denote by $g^{(p)}$ the generating function of $\mu_{p}$, and by $g_{n}^{(p)}=$ $g^{(p)} \circ \cdots \circ g^{(p)}$ the $n$-th iterate of $g^{(p)}$.

Theorem 2.3.1 Suppose that the convergences (2.1) and (2.2) hold and that the continuity condition (2.27) is satisfied. Suppose in addition that for every $\delta>0$,

$$
\begin{equation*}
\liminf _{p \rightarrow \infty} g_{\left[\delta \gamma_{p}\right]}^{(p)}(0)^{p}>0 \tag{2.28}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\left(\gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(H_{t}, t \geq 0\right) \tag{2.29}
\end{equation*}
$$

in the sense of weak convergence on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$.
Let us make some important remarks. Condition (2.28) can be restated in probabilistic terms as follows: For every $\delta>0$,

$$
\liminf _{p \rightarrow \infty} P\left[Y_{\left[\delta \gamma_{p}\right]}^{p}=0\right]>0
$$

(As will follow from our results, this implies that the extinction time of $Y^{p}$, scaled by $\gamma_{p}^{-1}$, converges in distribution to the extinction time of $Y$, which is finite a.s. under (2.27).) It is easy to see that the condition (2.28) is necessary for the conclusion (2.29) to hold. Indeed, suppose that (2.28) fails, so that there exists $\delta>0$ such that $P\left[Y_{\left[\delta \gamma_{p}\right]}^{p}=0\right]$ converges to 0 as $p \rightarrow \infty$, at least along a suitable subsequence. Clearly, this convergence also holds (along the same subsequence) if $Y^{p}$ starts at [ap] instead of $p$, for any fixed $a>0$. From the definition of the discrete height process, we get that

$$
P\left[\sup _{k \leq T_{[a p]}^{p}} H_{k}^{p} \geq\left[\delta \gamma_{p}\right]\right] \underset{p \rightarrow \infty}{\longrightarrow} 1
$$

where $T_{j}^{p}=\inf \left\{k \geq 0: V_{k}^{p}=-j\right\}$. From (2.2), we know that $\left(p \gamma_{p}\right)^{-1} T_{[a p]}^{p}$ converges in distribution to $T_{a}$. Since $T_{a} \downarrow 0$ as $a \downarrow 0$, a.s., we easily conclude that, for every $\varepsilon>0$,

$$
P\left[\sup _{t \leq \varepsilon} \gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p} \geq \frac{\left[\delta \gamma_{p}\right]}{\gamma_{p}}\right] \underset{p \rightarrow \infty}{\longrightarrow} 1
$$

and thus (2.29) cannot hold.
On the other hand, one might think that the condition (2.28) is automatically satisfied under (2.1) and (2.27). Let us explain why this is not the case. Suppose for simplicity that $\psi$ is of the type

$$
\psi(\lambda)=\alpha \lambda+\int_{(0, \infty)} \pi(d r)\left(e^{-\lambda r}-1+\lambda r\right),
$$

and for every $\varepsilon>0$ set

$$
\psi_{\varepsilon}(\lambda)=\alpha \lambda+\int_{(\varepsilon, \infty)} \pi(d r)\left(e^{-\lambda r}-1+\lambda r\right)
$$

Note that $\psi_{\varepsilon}(\lambda) \leq C_{\varepsilon} \lambda$ and so $\int_{1}^{\infty} \psi_{\varepsilon}(\lambda)^{-1} d \lambda=\infty$. Thus, if $Y^{\varepsilon}$ is a $\psi_{\varepsilon}$-CSBP started at 1 , we have $Y_{t}^{\varepsilon}>0$ for every $t>0$ a.s. (Grey [20], Theorem 1). It is easy to verify that

$$
\left(Y_{t}^{\varepsilon}, t \geq 0\right) \underset{\varepsilon \rightarrow 0}{\longrightarrow}\left(Y_{t}, t \geq 0\right)
$$

at least in the sense of the weak convergence of finite-dimensional marginals. Let us fix a sequence $\left(\varepsilon_{k}\right)$ decreasing to 0 . Recall from [27] that every continuous-state branching process can be obtained as a weak limit of rescaled Galton-Watson branching
processes. Thus for every $k$, we can find a subcritical or critical offspring distribution $\nu_{k}$, and two positive integers $p_{k} \geq k$ and $\gamma_{k} \geq k$, in such a way that if $Z^{k}=\left(Z_{j}^{k}, j \geq 0\right)$ is a Galton-Watson process with offspring distribution $\nu_{k}$ started at $Z_{0}^{k}=p_{k}$, the law of the rescaled process

$$
Z_{t}^{(k)}=\left(p_{k}\right)^{-1} Z_{\left[\gamma_{k} t\right]}^{k}
$$

is arbitrarily close to that of $Y^{\varepsilon_{k}}$. In particular, we may assume that $P\left[Z_{k}^{(k)}>\right.$ $0]>1-2^{-k}$, and that the rescaled processes $Z^{(k)}$ converge to $Y$ in the sense of weak convergence of finite-dimensional marginals. By Theorem 3.4 of [21], this convergence also holds in the functional sense in the Skorokhod space. However, the extinction time of $Z^{(k)}$ converges in probability to $+\infty$, and so the condition (2.28) cannot hold.

There is however a very important special case where (2.28) holds.
Theorem 2.3.2 Suppose that $\mu_{p}=\mu$ for every $p$ and that the convergence (2.1) holds. Then the condition (2.28) is automatically satisfied and the conclusion of Theorem 2.3 .1 holds.

As we will see in the proof, under the assumption of Theorem 2.3.2, the process $X$ must be stable with index $\alpha \in(1,2]$. Clearly condition (2.27) holds in that case.

Proof of Theorem 2.3.1. To simplify notation, we set $H_{t}^{(p)}=\gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}$ and $V_{t}^{(p)}=p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}$. In view of Theorem 2.2.1, the proof of Theorem 2.3.1 reduces to checking that the laws of the processes $\left(H_{t}^{(p)}, t \geq 0\right)$ are tight in the set of probability measures on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. By standard results (see e.g. Corollary 3.7.4 in [14]), it is enough to verify the following two properties:
(i) For every $t \geq 0$ and $\eta>0$, there exists a constant $K \geq 0$ such that

$$
\liminf _{p \rightarrow \infty} P\left[H_{t}^{(p)} \leq K\right] \geq 1-\eta
$$

(ii) For every $T>0$ and $\delta>0$,

$$
\lim _{n \rightarrow \infty} \limsup _{p \rightarrow \infty} P\left[\sup _{1 \leq i \leq 2^{n}} \sup _{t \in\left[(i-1) 2^{-n} T, i 2^{-n} T\right]}\left|H_{t}^{(p)}-H_{(i-1) 2^{-n} T}^{(p)}\right|>\delta\right]=0
$$

Property (i) is immediate from the convergence of finite-dimensional marginals. Thus the real problem is to prove (ii). We fix $\delta>0$ and $T>0$ and first observe that

$$
\begin{align*}
& P\left[\sup _{1 \leq i \leq 2^{n}} \sup _{t \in\left[(i-1) 2^{-n} T, i 2^{-n} T\right]}\left|H_{t}^{(p)}-H_{(i-1) 2^{-n} T}^{(p)}\right|>\delta\right]  \tag{2.30}\\
& \quad \leq A_{1}(n, p)+A_{2}(n, p)+A_{3}(n, p)
\end{align*}
$$

where

$$
\begin{aligned}
& A_{1}(n, p)=P\left[\sup _{1 \leq i \leq 2^{n}}\left|H_{i 2^{-n} T}^{(p)}-H_{(i-1) 2^{-n} T}^{(p)}\right|>\frac{\delta}{5}\right] \\
& A_{2}(n, p)=P\left[\sup _{t \in\left[(i-1) 2^{-n} T, i 2^{-n} T\right]} H_{t}^{(p)}>H_{(i-1) 2^{-n} T}^{(p)}+\frac{4 \delta}{5} \text { for some } 1 \leq i \leq 2^{n}\right] \\
& A_{3}(n, p)=P\left[\inf _{t \in\left[(i-1) 2^{-n} T, i 2^{-n} T\right]} H_{t}^{(p)}<H_{i 2^{-n} T}^{(p)}-\frac{4 \delta}{5} \text { for some } 1 \leq i \leq 2^{n}\right]
\end{aligned}
$$

The term $A_{1}$ is easy to bound. By the convergence of finite-dimensional marginals, we have

$$
\limsup _{p \rightarrow \infty} A_{1}(n, p) \leq P\left[\sup _{1 \leq i \leq 2^{n}}\left|H_{i 2^{-n} T}-H_{(i-1) 2^{-n} T}\right| \geq \frac{\delta}{5}\right]
$$

and the path continuity of the process $H$ ensures that the right-hand side tends to 0 as $n \rightarrow \infty$.

To bound the terms $A_{2}$ and $A_{3}$, we introduce the stopping times $\tau_{k}^{(p)}, k \geq 0$ defined by induction as follows:

$$
\begin{aligned}
& \tau_{0}^{(p)}=0 \\
& \tau_{k+1}^{(p)}=\inf \left\{t \geq \tau_{k}^{(p)}: H_{t}^{(p)}>\inf _{\tau_{k}^{(p)} \leq r \leq t} H_{r}^{(p)}+\frac{\delta}{5}\right\} .
\end{aligned}
$$

Let $i \in\left\{1, \ldots, 2^{n}\right\}$ be such that

$$
\begin{equation*}
\sup _{t \in\left[(i-1) 2^{-n} T, i 2^{-n} T\right]} H_{t}^{(p)}>H_{(i-1) 2^{-n} T}^{(p)}+\frac{4 \delta}{5} \tag{2.31}
\end{equation*}
$$

Then it is clear that the interval $\left[(i-1) 2^{-n} T, i 2^{-n} T\right]$ must contain at least one of the random times $\tau_{k}^{(p)}, k \geq 0$. Let $\tau_{j}^{(p)}$ be the first such time. By construction we have

$$
\sup _{t \in\left[(i-1) 2^{-n} T, \tau_{j}^{(p)}\right)} H_{t}^{(p)} \leq H_{(i-1) 2^{-n} T}^{(p)}+\frac{\delta}{5},
$$

and since the positive jumps of $H^{(p)}$ are of size $\gamma_{p}^{-1}$, we get also

$$
H_{\tau_{j}^{(p)}}^{(p)} \leq H_{(i-1) 2^{-n} T}^{(p)}+\frac{\delta}{5}+\gamma_{p}^{-1}<H_{(i-1) 2^{-n} T}^{(p)}+\frac{2 \delta}{5}
$$

provided that $\gamma_{p}>5 / \delta$. From (2.31), we have then

$$
\sup _{t \in\left[\tau_{j}^{(p)}, i 2^{-n} T\right]} H_{t}^{(p)}>H_{\tau_{j}^{(p)}}^{(p)}+\frac{\delta}{5},
$$

which implies that $\tau_{j+1}^{(p)} \leq i 2^{-n} T$. Summarizing, we get for $p$ large enough so that $\gamma_{p}>5 / \delta$

$$
\begin{equation*}
A_{2}(n, p) \leq P\left[\tau_{k}^{(p)}<T \text { and } \tau_{k+1}^{(p)}-\tau_{k}^{(p)}<2^{-n} T \text { for some } k \geq 0\right] \tag{2.32}
\end{equation*}
$$

A similar argument gives exactly the same bound for the quantity $A_{3}(n, p)$.
The following lemma is directly inspired from [14] p.134-135.
Lemma 2.3.3 For every $x>0$ and $p \geq 1$, set

$$
G_{p}(x)=P\left[\tau_{k}^{(p)}<T \text { and } \tau_{k+1}^{(p)}-\tau_{k}^{(p)}<x \text { for some } k \geq 0\right]
$$

and

$$
F_{p}(x)=\sup _{k \geq 0} P\left[\tau_{k}^{(p)}<T \text { and } \tau_{k+1}^{(p)}-\tau_{k}^{(p)}<x\right]
$$

Then, for every integer $L \geq 1$,

$$
G_{p}(x) \leq L F_{p}(x)+L e^{T} \int_{0}^{\infty} d y e^{-L y} F_{p}(y)
$$

Proof. For every integer $L \geq 1$, we have

$$
\begin{aligned}
G_{p}(x) & \leq \sum_{k=0}^{L-1} P\left[\tau_{k}^{(p)}<T \text { and } \tau_{k+1}^{(p)}-\tau_{k}^{(p)}<x\right]+P\left[\tau_{L}^{(p)}<T\right] \\
& \leq L F_{p}(x)+e^{T} E\left[1_{\left\{\tau_{L}^{(p)}<T\right\}} \exp \left(-\sum_{k=0}^{L-1}\left(\tau_{k+1}^{(p)}-\tau_{k}^{(p)}\right)\right)\right] \\
& \leq L F_{p}(x)+e^{T} \prod_{k=0}^{L-1} E\left[1_{\left\{\tau_{L}^{(p)}<T\right\}} \exp \left(-L\left(\tau_{k+1}^{(p)}-\tau_{k}^{(p)}\right)\right)\right]^{1 / L} .
\end{aligned}
$$

Then observe that for every $k \in\{0,1, \ldots, L-1\}$,

$$
\begin{aligned}
E\left[1_{\left\{\tau_{L}^{(p)}<T\right\}} \exp \left(-L\left(\tau_{k+1}^{(p)}-\tau_{k}^{(p)}\right)\right)\right] & \leq E\left[1_{\left\{\tau_{k}^{(p)}<T\right\}} \int_{\tau_{k+1}^{(p)}-\tau_{k}^{(p)}}^{\infty} d y L e^{-L y}\right] \\
& \leq \int_{0}^{\infty} d y L e^{-L y} F_{p}(y)
\end{aligned}
$$

The desired result follows.
Thanks to Lemma 2.3.3, the limiting behavior of the right-hand side of (2.32) will be reduced to that of the function $F_{p}(x)$. To handle $F_{p}(x)$, we use the next lemma.

Lemma 2.3.4 The random variables $\tau_{k+1}^{(p)}-\tau_{k}^{(p)}$ are independent and identically distributed. Under the assumptions of Theorem 2.3.1, we have

$$
\lim _{x \downarrow 0}\left(\limsup _{p \rightarrow \infty} P\left[\tau_{1}^{(p)} \leq x\right]\right)=0
$$

We need a simple lemma.
Lemma 2.3.5 Let $V$ be a random walk on $\mathbb{Z}$. For every $n \geq 0$, set

$$
\begin{equation*}
H_{n}^{o}=\operatorname{Card}\left\{k \in\{0,1, \ldots, n-1\}: V_{k}=\inf _{k \leq j \leq n} V_{j}\right\} \tag{2.33}
\end{equation*}
$$

Let $\tau$ be a stopping time of the filtration $\left(\mathcal{F}_{n}^{o}\right)$ generated by $V$. Then the process

$$
\left(H_{\tau+n}^{o}-\inf _{\tau \leq k \leq \tau+n} H_{k}^{o}, n \geq 0\right)
$$

is independent of $\mathcal{F}_{\tau}^{o}$ and has the same distribution as $\left(H_{n}^{o}, n \geq 0\right)$.
Proof. By considering the first time after $\tau$ where the random walk $V$ attains its minimum over $[\tau, \tau+n]$, one easily gets

$$
\inf _{\tau \leq k \leq \tau+n} H_{k}^{o}=\operatorname{Card}\left\{k \in\{0,1, \ldots, \tau-1\}: V_{k}=\inf _{k \leq j \leq \tau+n} V_{j}\right\}
$$

Hence,

$$
\begin{aligned}
H_{\tau+n}^{o}-\inf _{\tau \leq k \leq \tau+n} H_{k}^{o} & =\operatorname{Card}\left\{k \in\{\tau, \ldots, \tau+n-1\}: V_{k}=\inf _{k \leq j \leq \tau+n} V_{j}\right\} \\
& =\operatorname{Card}\left\{k \in\{0, \ldots, n-1\}: V_{k}^{\tau}=\inf _{k \leq j \leq n} V_{j}^{\tau}\right\},
\end{aligned}
$$

where $V^{\tau}$ denotes the shifted random walk $V_{n}^{\tau}=V_{\tau+n}-V_{\tau}$. Since $V^{\tau}$ is independent of $\mathcal{F}_{\tau}$ and has the same distribution as $V$, the desired result follows from the previous formula and (2.33).
Proof of Lemma 2.3.4. Fix $k \geq 1$ and set for every $t \geq 0$,

$$
\tilde{H}_{t}^{(p)}=H_{\tau_{k}^{(p)}+t}^{(p)}-\inf _{\tau_{k}^{(p)} \leq r \leq \tau_{k}^{(p)}+t} H_{r}^{(p)} .
$$

As a consequence of Lemma 2.3.5, the process $\left(\tilde{H}_{t}^{(p)}, t \geq 0\right)$ is independent of the past of $V^{(p)}$ up to the stopping time $\tau_{k}^{(p)}$ and has the same distribution as $\left(H_{t}^{(p)}, t \geq 0\right)$. Since by definition

$$
\tau_{k+1}^{(p)}-\tau_{k}^{(p)}=\inf \left\{t \geq 0: \tilde{H}_{t}^{(p)}>\frac{\delta}{5}\right\}
$$

the first assertion of the lemma follows.
Let us turn to the second assertion. To simplify notation, we write $\delta^{\prime}=\delta / 5$. For every $\eta>0$, set

$$
T_{\eta}^{(p)}=\inf \left\{t \geq 0: V_{t}^{(p)}=-\frac{[p \eta]}{p}\right\}
$$

Then,

$$
P\left[\tau_{1}^{(p)} \leq x\right]=P\left[\sup _{s \leq x} H_{s}^{(p)}>\delta^{\prime}\right] \leq P\left[\sup _{s \leq T_{\eta}^{(p)}} H_{s}^{(p)}>\delta^{\prime}\right]+P\left[T_{\eta}^{(p)}<x\right]
$$

On one hand,

$$
\limsup _{p \rightarrow \infty} P\left[T_{\eta}^{(p)}<x\right] \leq P\left[T_{\eta} \leq x\right]
$$

and for any choice of $\eta>0$, the right-hand side goes to zero as $x \downarrow 0$. On the other hand, the construction of the discrete height process shows that the quantity

$$
\sup _{s \leq T_{\eta}^{(p)}} H_{s}^{(p)}
$$

is distributed as $\gamma_{p}^{-1}\left(M_{p}-1\right)$, where $M_{p}$ is the extinction time of a Galton-Watson process with offspring distribution $\mu_{p}$, started at $[p \eta]$. Hence,

$$
P\left[\sup _{s \leq T_{\eta}^{(p)}} H_{s}^{(p)}>\delta^{\prime}\right]=1-g_{\left[\delta^{\prime} \gamma_{p}\right]+1}^{(p)}(0)^{[p \eta]},
$$

and our assumption (2.28) implies that

$$
\lim _{\eta \rightarrow 0}\left(\limsup _{p \rightarrow \infty} P\left[\sup _{s \leq T_{\eta}^{(p)}} H_{s}^{(p)}>\delta^{\prime}\right]\right)=0
$$

The second assertion of the lemma now follows.
We can now complete the proof of Theorem 2.3.1. Set:

$$
F(x)=\underset{p \rightarrow \infty}{\limsup } F_{p}(x), \quad G(x)=\underset{p \rightarrow \infty}{\limsup } G_{p}(x)
$$

Lemma 2.3.4 immediately shows that $F(x) \downarrow 0$ as $x \downarrow 0$. On the other hand, we get from Lemma 2.3.3 that for every integer $L \geq 1$,

$$
G(x) \leq L F(x)+L e^{T} \int_{0}^{\infty} d y e^{-L y} F(y)
$$

It follows that we have also $G(x) \downarrow 0$ as $x \downarrow 0$. By (2.32), this gives

$$
\lim _{n \rightarrow \infty}\left(\limsup _{p \rightarrow \infty} A_{2}(n, p)\right)=0
$$

and the same property holds for $A_{3}(n, p)$. This completes the proof of (ii) and of Theorem 2.3.1.

Proof of Theorem 2.3.2. We now assume that $\nu_{p}=\nu$ for every $p$ and so $g_{n}^{(p)}=g_{n}$. We first observe that the process $X$ must be stable. This is not immediate, since the convergence (2.2) a priori implies only that $\nu$ belongs to the domain of partial attraction of the law of $X_{1}$, which is not enough to conclude that $\nu$ is stable. However, the conditions $(\mathrm{C} 1)-(\mathrm{C} 3)$, which are equivalent to (2.2), immediately show that the sequence $\gamma_{p} / \gamma_{p+1}$ converges to 1 as $p \rightarrow \infty$. Then Theorem 2.3 in [37] implies that $\nu$ belongs to the domain of attraction of the law of $X_{1}$, and by classical results the law of $X_{1}$ must be stable with index $\alpha \in(0,2]$. We can exclude $\alpha \in(0,1]$ thanks to our assumptions (H2) and (H3) (the latter is only needed to exclude the trivial case $\psi(\lambda)=c \lambda)$. Thus $\alpha \in(1,2]$ and $\psi(\lambda)=c \lambda^{\alpha}$ for some $c>0$. As a consequence of (1.32), we have $E\left[e^{-\lambda Y_{\delta}}\right]=\exp -\left(\lambda^{-\bar{\alpha}}+c \bar{\alpha} \delta\right)^{-1 / \bar{\alpha}}$, where $\bar{\alpha}=\alpha-1$. In particular, $P\left[Y_{\delta}=0\right]=\exp -(c \bar{\alpha} \delta)^{-1 / \bar{\alpha}}>0$.

Let $g=g_{1}$ be the generating function of $\mu$. We have $g^{\prime}(1)=\sum k \mu(k)=1$, because otherwise this would contradict (2.2). From Theorem 2 in [16], p.577, the function

$$
\sum_{k \geq x} \mu(k)
$$

must be regularly varying as $x \rightarrow \infty$, with exponent $-\alpha$. Then note that

$$
g\left(e^{-\lambda}\right)-1+\lambda=\sum_{k=0}^{\infty} \mu(k)\left(e^{-\lambda k}-1+\lambda k\right)=\lambda \int_{0}^{\infty} d x\left(1-e^{-\lambda x}\right) \sum_{k \geq x} \mu(k)
$$

An elementary argument shows that $g\left(e^{-\lambda}\right)-1+\lambda$ is also regularly varying as $\lambda \rightarrow 0$ with exponent $\alpha$. Put differently,

$$
g(r)=r+(1-r)^{\alpha} L(1-r), \quad 0 \leq r<1
$$

where the function $L(x)$ is slowly varying as $x \rightarrow 0$. This is exactly what we need to apply a result of Slack [47].

Let $Z_{1}^{(p)}$ be a random variable distributed as $\left(1-g_{\left[\delta \gamma_{p}\right]}(0)\right)$ times the value at time $\left[\delta \gamma_{p}\right.$ ] of a Galton-Watson process with offspring distribution $\mu$ started with one individual at time 0 and conditioned to be non-extinct at time $\left[\delta \gamma_{p}\right]$. Theorem 1 of [47] implies that

$$
Z_{1}^{(p)} \xrightarrow[p \rightarrow \infty]{(\mathrm{d})} U
$$

where $U>0$ a.s. In particular, we can choose positive constants $c_{0}$ and $c_{1}$ so that $P\left[Z_{1}^{(p)}>c_{0}\right]>c_{1}$ for all $p$ sufficiently large. On the other hand, we have

$$
\frac{1}{p} Y_{\left[\delta \gamma_{p}\right]}^{p} \stackrel{(\mathrm{~d})}{=} \frac{1}{p\left(1-g_{\left[\delta \gamma_{p}\right]}(0)\right)}\left(Z_{1}^{(p)}+\cdots+Z_{M_{p}}^{(p)}\right)
$$

where $Z_{1}^{(p)}, Z_{2}^{(p)}, \ldots$ are i.i.d., and $M_{p}$ is independent of the sequence $\left(Z_{j}^{(p)}\right)$ and has a binomial $B\left(p, 1-g_{\left[\delta \gamma_{p}\right]}(0)\right)$ distribution.

It is now easy to obtain the condition (2.28). Fix $\delta>0$. Clearly it suffices to verify that the sequence $p\left(1-g_{\left[\delta \gamma_{p}\right]}(0)\right)$ is bounded. If not the case, we can choose a sequence $\left(p_{k}\right)$ such that $p_{k}\left(1-g_{\left[\delta \gamma_{p_{k}}\right]}(0)\right)$ converges to $\infty$. From the previous representation for the law of $\frac{1}{p} Y_{\left[\delta \gamma_{p}\right]}^{p}$, it then follows that

$$
P\left[\frac{1}{p_{k}} Y_{\left[\delta \gamma_{p_{k}}\right]}^{p_{k}}>c_{0} c_{1}\right] \underset{k \rightarrow \infty}{\longrightarrow} 1 .
$$

From (2.1), we get that $P\left[Y_{\delta} \geq c_{0} c_{1}\right]=1$, which gives a contradiction since $P\left[Y_{\delta}=\right.$ $0]>0$. This completes the proof of (2.28).

Finally, since (2.27) holds, we can apply Theorem 2.3.1.

### 2.4 Convergence of contour processes

In this section, we show that the limit theorems obtained in the previous section for rescaled discrete height processes can be formulated as well in terms of the contour processes of the Galton-Watson trees. The proof relies on simple connections between the height process and the contour process of a sequence of Galton-Watson trees.

To begin with, we consider a (subcritical or critical) offspring distribution $\mu$, and a sequence of independent $\mu$-Galton-Watson trees. Let $\left(H_{n}, n \geq 0\right)$ and $\left(C_{t}, t \geq 0\right)$ be respectively the height process and the contour process associated with this sequence of trees (see Section 0.2). We also set

$$
K_{n}=2 n-H_{n}
$$

Note that the sequence $K_{n}$ is strictly increasing and $K_{n} \geq n$.
Recall that the value at time $n$ of the height process corresponds to the generation of the individual visited at time $n$, assuming that individuals are visited in lexicographical order one tree after another. It is easily checked by induction on $n$ that [ $K_{n}, K_{n+1}$ ] is exactly the time interval during which the contour process goes from the individual $n$ to the individual $n+1$. From this observation, we get

$$
\sup _{t \in\left[K_{n}, K_{n+1}\right]}\left|C_{t}-H_{n}\right| \leq\left|H_{n+1}-H_{n}\right|+1
$$

A more precise argument for this bound follows from the explicit formula for $C_{t}$ in terms of the height process: For $t \in\left[K_{n}, K_{n+1}\right]$,

$$
\begin{array}{ll}
C_{t}=\left(H_{n}-\left(t-K_{n}\right)\right)^{+} & \text {if } t \in\left[K_{n}, K_{n+1}-1\right] \\
C_{t}=\left(H_{n+1}-\left(K_{n+1}-t\right)\right)^{+} & \text {if } t \in\left[K_{n+1}-1, K_{n+1}\right]
\end{array}
$$

These formulas are easily checked by induction on $n$.
Define a random function $f: \mathbb{R}_{+} \longrightarrow \mathbb{Z}_{+}$by setting $f(t)=n$ iff $t \in\left[K_{n}, K_{n+1}\right)$. From the previous bound, we get for every integer $m \geq 1$,

$$
\begin{equation*}
\sup _{t \in[0, m]}\left|C_{t}-H_{f(t)}\right| \leq \sup _{t \in\left[0, K_{m}\right]}\left|C_{t}-H_{f(t)}\right| \leq 1+\sup _{n \leq m}\left|H_{n+1}-H_{n}\right| \tag{2.34}
\end{equation*}
$$

Similarly, it follows from the definition of $K_{n}$ that

$$
\begin{equation*}
\sup _{t \in[0, m]}\left|f(t)-\frac{t}{2}\right| \leq \sup _{t \in\left[0, K_{m}\right]}\left|f(t)-\frac{t}{2}\right| \leq \frac{1}{2} \sup _{n \leq m} H_{n}+1 \tag{2.35}
\end{equation*}
$$

We now come back to the setting of the previous sections, considering for every $p \geq 1$ a sequence of independent Galton-Watson trees with offspring distribution $\mu_{p}$. For every $p \geq 1$, we denote by ( $C_{t}^{p}, t \geq 0$ ) the corresponding contour process.

Theorem 2.4.1 Suppose that the convergences (2.2) and (2.29) hold. Then,

$$
\begin{equation*}
\left(\gamma_{p}^{-1} C_{p \gamma_{p} t}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{(\mathrm{~d})}\left(H_{t / 2}, t \geq 0\right) . \tag{2.36}
\end{equation*}
$$

In particular, (2.36) holds under the assumptions of Theorem 2.3.1 or those of Theorem 2.3.2.

Proof. For every $p \geq 1$, write $f_{p}$ for the analogue of the function $f$ introduced above. Also set $\varphi_{p}(t)=\left(p \gamma_{p}\right)^{-1} f_{p}\left(p \gamma_{p} t\right)$. By (2.34), we have for every $m \geq 1$,

$$
\begin{equation*}
\sup _{t \leq m}\left|\frac{1}{\gamma_{p}} C_{p \gamma_{p} t}^{p}-\frac{1}{\gamma_{p}} H_{p \gamma_{p} \varphi_{p}(t)}^{p}\right| \leq \frac{1}{\gamma_{p}}+\frac{1}{\gamma_{p}} \sup _{n \leq m p \gamma_{p}}\left|H_{n+1}^{p}-H_{n}^{p}\right| \underset{p \rightarrow \infty}{\longrightarrow} 0 \tag{2.37}
\end{equation*}
$$

in probability, by (2.29).
On the other hand, we get from (2.35)

$$
\begin{equation*}
\sup _{t \leq m}\left|\varphi_{p}(t)-\frac{t}{2}\right| \leq \frac{1}{2 p \gamma_{p}} \sup _{k \leq m p \gamma_{p}} H_{k}^{p}+\frac{1}{p \gamma_{p}} \underset{p \rightarrow \infty}{\longrightarrow} 0 \tag{2.38}
\end{equation*}
$$

in probability, by (2.29).
The statement of the theorem now follows from (2.29), (2.37) and (2.38).

### 2.5 A joint convergence and an application to conditioned trees

The convergences in distribution (2.29) and (2.36) hold jointly with (2.1) and (2.2). This fact is useful in applications and we state it here as a corollary.

As previously, we consider for every $p$ a sequence of independent $\mu_{p}$-Galton-Watson trees and we denote by $\left(H_{n}^{p}, n \geq 0\right)$ the associated height process and by $\left(C_{t}^{p}, t \geq 0\right)$ the associated contour process. The random walk $V^{p}$ with jump distribution $\nu_{p}(k)=$ $\mu_{p}(k+1)$ is related to $H^{p}$ via formula (2.4). Finally, for every integer $k \geq 0$, we denote by $Y_{k}^{p}$ the number of individuals at generation $k$ in the first $p$ trees of the sequence, so that, in agreement with the previous notation, $\left(Y_{n}^{p}, n \geq 0\right)$ is a Galton-Watson process with offspring distribution $\mu_{p}$ started at $Y_{0}^{p}=p$.

Recall that ( $L_{t}^{a}, a \geq 0, t \geq 0$ ) denote the local times of the (continuous-time) height process associated with the Lévy process $X$. From Theorem 1.4.1, we know that ( $L_{T_{1}}^{a}, a \geq 0$ ) is a $\psi$-CSBP and thus has a càdlàg modification.

Corollary 2.5.1 Suppose that the assumptions of Theorem 2.3.1 are satisfied. Then,

$$
\left(p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}, \gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}, \gamma_{p}^{-1} C_{2 p \gamma_{p} t}^{p} ; t \geq 0\right) \underset{p \rightarrow \infty}{(\mathrm{~d})}\left(X_{t}, H_{t}, H_{t} ; t \geq 0\right)
$$

in distribution in $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right)$. We have also

$$
\left(p^{-1} Y_{\left[\gamma_{p} a\right]}^{p}, a \geq 0\right) \underset{p \rightarrow \infty}{(\mathrm{~d})}\left(L_{T_{1}}^{a}, a \geq 0\right)
$$

in distribution in $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. Furthermore, these two convergences hold jointly, in the sense that, for any bounded continuous function $F$ on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right) \times \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$,

$$
\begin{aligned}
& \lim _{p \rightarrow \infty} E\left[F\left(\left(p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}, \gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}, \gamma_{p}^{-1} C_{2 p \gamma_{p} t}^{p}\right)_{t \geq 0},\left(p^{-1} Y_{\left[\gamma_{p} a\right]}^{p}\right)_{a \geq 0}\right)\right] \\
& \quad=E\left[F\left(\left(X_{t}, H_{t}, H_{t}\right)_{t \geq 0},\left(L_{T_{1}}^{a}\right)_{a \geq 0}\right)\right] .
\end{aligned}
$$

Proof. To simplify notation, write $V_{t}^{(p)}=p^{-1} V_{\left[p \gamma_{p} t\right]}^{p}, H_{t}^{(p)}=\gamma_{p}^{-1} H_{\left[p \gamma_{p} t\right]}^{p}, C_{t}^{(p)}=$ $\gamma_{p}^{-1} C_{2 p \gamma_{p} t}^{p}$ and $Y_{a}^{(p)}=p^{-1} Y_{\left[\gamma_{p} a\right]}^{p}$. By (2.2), (2.29) and (2.36), we know that each of the three sequences of the laws of the processes $V^{(p)}, H^{(p)}, C^{(p)}$ is tight, and furthermore $H^{(p)}$ and $C^{(p)}$ converge in distribution towards a continuous process. By a standard result (see e.g. Corollary II.3.33 in [24]), we get that the laws of the triples $\left(V^{(p)}, H^{(p)}, C^{(p)}\right)$ are tight in $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right)$. Let $\left(X, H^{*}, H^{* *}\right)$ be a weak limit point of this sequence of triples (with a slight abuse of notation, we may assume that the first component of the limiting triple is the underlying Lévy process $X$ ). By the Skorokhod representation theorem, we may assume that along a subsequence,

$$
\left(V^{(p)}, H^{(p)}, C^{(p)}\right) \longrightarrow\left(X, H^{*}, H^{* *}\right)
$$

a.s. in $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right)$. However, the convergence (2.6) and a time-reversal argument imply that

$$
\lim _{p \rightarrow \infty} H_{t}^{(p)}=\widehat{L}_{t}^{(t)}=H_{t}
$$

in probability. This is enough to conclude that $H_{t}^{*}=H_{t}$. Similarly, the proof of Theorem 2.4.1 shows that

$$
\lim _{p \rightarrow \infty}\left(C_{t}^{(p)}-H_{t}^{(p)}\right)=0
$$

in probability. This yields $H_{t}^{* *}=H_{t}^{*}=H_{t}$ and we see that the limiting triple is equal to $(X, H, H)$ and does not depend on the choice of the subsequence. The first convergence of the corollary now follows.

By (2.1), we know that

$$
\left(Y_{a}^{(p)}, a \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(Y_{a}, a \geq 0\right)
$$

where $Y$ is a $\psi$-CSBP started at 1 . Since we also know that $\left(L_{T_{1}}^{a}, a \geq 0\right)$ is a $\psi$-CSBP started at 1 , the second convergence in distribution is immediate, and the point is to verify that this convergence holds jointly with the first one. To this end, note that the laws of the pairs $\left(\left(V^{(p)}, H^{(p)}, C^{(p)}\right), Y^{(p)}\right)$ are tight in the space of probability measures on $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right) \times \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. By extracting a subsequence and using the Skorokhod representation theorem, we may assume that

$$
\left(\left(V^{(p)}, H^{(p)}, C^{(p)}\right), Y^{(p)}\right) \underset{p \rightarrow \infty}{\longrightarrow}((X, H, H), Z)
$$

a.s. in $\mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}^{3}\right) \times \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$. The proof will be finished if we can verify that $Z_{a}=L_{T_{1}}^{a}$, the local time of $H$ at level $a$ and time $T_{1}$. To this end, let $g$ be a Lipschitz continuous function from $\mathbb{R}_{+}$into $\mathbb{R}_{+}$with compact support. The preceding convergence implies

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \int_{0}^{\infty} g(a) Y_{a}^{(p)} d a=\int_{0}^{\infty} g(a) Z_{a} d a, \quad \text { a.s. } \tag{2.39}
\end{equation*}
$$

On the other hand, let $T_{p}^{p}$ be the hitting time of $-p$ by $V^{p}$. The convergence of $V^{(p)}$ towards $X$ easily implies

$$
\begin{equation*}
\lim _{p \rightarrow \infty} \frac{1}{p \gamma_{p}} T_{p}^{p}=\inf \left\{t \geq 0: X_{t}=-1\right\}=T_{1}, \quad \text { a.s. } \tag{2.40}
\end{equation*}
$$

Then, from the definition of the height process of a sequence of trees, we have

$$
\int_{0}^{\infty} g(a) Y_{a}^{(p)} d a=\int_{0}^{\infty} g(a) \frac{1}{p} Y_{\left[\gamma_{p} a\right]}^{p} d a
$$

$$
\begin{aligned}
& =\frac{1}{p} \sum_{k=0}^{\infty} \int_{\gamma_{p}^{-1} k}^{\gamma_{p}^{-1}(k+1)} g(a)\left(\sum_{j=0}^{T_{p}^{p}-1} 1_{\left\{H_{j}^{p}=k\right\}}\right) d a \\
& =\frac{1}{p} \sum_{j=0}^{T_{p}^{p}-1} \int_{\gamma_{p}^{-1} H_{j}^{p}}^{\gamma_{p}^{-1}\left(H_{j}^{p}+1\right)} g(a) d a \\
& =\frac{1}{p \gamma_{p}} \sum_{j=0}^{T_{p}^{p}-1} g\left(\gamma_{p}^{-1} H_{j}^{p}\right)+O\left(\frac{1}{p \gamma_{p}^{2}} T_{p}^{p}\right) \\
& =\int_{0}^{\left(p \gamma_{p}\right)^{-1} T_{p}^{p}} g\left(\gamma_{p}^{-1} H_{\left[p \gamma_{p} s\right]}^{p}\right) d s+O\left(\frac{1}{p \gamma_{p}^{2}} T_{p}^{p}\right)
\end{aligned}
$$

and in view of (2.40) this converges to

$$
\int_{0}^{T_{1}} g\left(H_{s}\right) d s=\int_{0}^{\infty} g(a) L_{T_{1}}^{a} d a
$$

Comparing with (2.39), we conclude that

$$
\int_{0}^{\infty} g(a) Z_{a} d a=\int_{0}^{\infty} g(a) L_{T_{1}}^{a} d a .
$$

This implies that $Z_{a}=L_{T_{1}}^{a}$ and completes the proof.
As an application, we now discuss conditioned trees. Fix $T>0$ and on some probability space, consider a $\mu_{p}$-Galton-Watson tree conditioned on non-extinction at generation $\left[\gamma_{p} T\right]$, which is denoted by $\tilde{T}^{p}$. Let $\tilde{H}^{p}=\left(\tilde{H}_{n}^{p}, n \geq 0\right)$ be the associated height process, with the convention that $\tilde{H}_{n}^{p}=0$ for $n \geq \operatorname{Card}\left(\tilde{T}^{p}\right)$.

Proposition 2.5.2 Under the assumptions of Theorem 2.3.1, we have

$$
\left(\gamma_{p}^{-1} \tilde{H}_{\left[p \gamma_{p} t\right]}^{p}, t \geq 0\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}\left(\tilde{H}_{t}, t \geq 0\right),
$$

where the limiting process $\tilde{H}$ is distributed as $H$ under $N\left(\cdot \mid \sup H_{s} \geq T\right)$.
Remark. We could have stated a similar result for the contour process instead of the discrete height process.
Proof. Write $\tilde{H}_{s}^{(p)}=\gamma_{p}^{-1} \tilde{H}_{\left[p \gamma_{p} s\right]}^{p}$ to simplify notation. Also let $H_{s}^{(p)}=\gamma_{p}^{-1} H_{\left[p \gamma_{p} s\right]}^{p}$ be as above the rescaled height process for a sequence of independent $\mu_{p}$-Galton-Watson trees. Set

$$
\begin{aligned}
& R_{T}^{(p)}=\inf \left\{s \geq 0: H_{s}^{(p)}=\frac{\left[\gamma_{p} T\right]}{\gamma_{p}}\right\} \\
& G_{T}^{(p)}=\sup \left\{s \leq R_{T}^{(p)}: H_{s}^{(p)}=0\right\} \\
& D_{T}^{(p)}=\inf \left\{s \geq R_{T}^{(p)}: H_{s}^{(p)}=0\right\}
\end{aligned}
$$

Then without loss of generality we may assume that

$$
\tilde{H}_{s}^{(p)}=H_{\left(G_{T}^{(p)}+s\right) \wedge D_{T}^{(p)}}^{(p)}, \quad s \geq 0 .
$$

This is simply saying that the first tree with height at least $\left[\gamma_{p} T\right]$ in a sequence of independent $\mu_{p}$-Galton-Watson trees is a $\mu_{p}$-Galton-Watson tree conditioned on non-extinction at generation $\left[\gamma_{p} T\right]$.

Set

$$
\begin{aligned}
& R_{T}=\inf \left\{s \geq 0: H_{s}=T\right\}, \\
& G_{T}=\sup \left\{s \leq R_{T}: H_{s}=0\right\}, \\
& D_{T}=\inf \left\{s \geq R_{T}: H_{s}=0\right\},
\end{aligned}
$$

and note that we may take $\tilde{H}_{s}=H_{\left(G_{T}+s\right) \wedge D_{T}}$, by excursion theory for $X-I$.
We now claim that the convergence in distribution of $\tilde{H}^{(p)}$ towards $\tilde{H}$ follows from the previous corollary, and more precisely from the joint convergence

$$
\left(V^{(p)}, H^{(p)}\right) \underset{p \rightarrow \infty}{\stackrel{(\mathrm{~d})}{\longrightarrow}}(X, H)
$$

It is again convenient to use the Skorokhod representation theorem and to assume that the latter convergence holds a.s. We can then prove that $\tilde{H}^{(p)}$ converges a.s. towards $\tilde{H}$.

To this end we need a technical lemma about the height process. We state it in greater generality than needed here in view of other applications.

Lemma 2.5.3 Let $b>0$. Then $P$ a.s. or $N$ a.e. $b$ is not $a$ local maximum nor $a$ local minimum of the height process.

Proof. Let

$$
D=\left\{b>0: P\left[\sup _{p \leq s \leq q} H_{s}=b\right]>0 \text { for some rationals } q>p \geq 0\right\} .
$$

Clearly $D$ is at most countable. However, from Proposition 1.3.1 and the relation between the height process and the exploration process, it immediately follows that if $b \in D$ then $b-a \in D$ for every $a \in[0, b)$. This is only possible if $D=\emptyset$. The case of local minima is treated in the same way.

It follows from the lemma that we have also $R_{T}=\inf \left\{s \geq 0: H_{s}>T\right\}$. Then the a.s. convergence of $H^{(p)}$ towards $H$ easily implies that $R_{T}^{(p)}$ converges to $R_{T}$ a.s., and that

$$
\limsup _{p \rightarrow \infty} G_{T}^{(p)} \leq G_{T}, \quad \liminf _{p \rightarrow \infty} D_{T}^{(p)} \geq D_{T}
$$

To get reverse inequalities, we may argue as follows. Recall that the support of the random measure $d I_{s}$ is exactly the set $\left\{s: H_{s}=0\right\}$, so that for every fixed $s \geq 0$, we have $I_{s}>I_{R_{T}}$ a.s. on the set $\left\{s<G_{T}\right\}$. If $I_{s}^{(p)}=\inf \left\{V_{r}^{(p)}, r \leq s\right\}$, the a.s. convergence of $V^{(p)}$ to $X$ implies that $I_{s}^{(p)}$ converges to $I_{s}$ uniformly on compact sets, a.s. It readily follows that a.s. on the set $\left\{s<G_{T}\right\}$ we have $I_{s}^{(p)}>I_{R_{T}^{(p)}}^{(p)}$ for all $p$ sufficiently large. Hence a.s. for $p$ large, we have $s<G_{T}^{(p)}$ on the set $\left\{s<G_{T}\right\}$. We conclude that $G_{T}^{(p)} \rightarrow G_{T}$ a.s., and a similar argument gives $D_{T}^{(p)} \rightarrow D_{T}$. From the preceding formulas for $\tilde{H}^{(p)}$ and $\tilde{H}$, it follows that $\tilde{H}^{(p)} \rightarrow \tilde{H}$ a.s. This completes the proof of the proposition.
Remark. The methodology of proof of Proposition 2.5.2 could be applied to other conditioned limit theorems. For instance, we could consider the rescaled height (or contour) process of the $\mu_{p}$-Galton-Watson tree conditioned to have at least $p \gamma_{p}$ vertices and derive a convergence towards the excursion of the height process $H$ conditioned to have length greater than 1 . We will leave such extensions to the reader. We point out here that it is much harder to handle degenerate conditionings. To give
an important example, consider the case where $\mu_{p}=\mu$ for every $p$. It is natural to ask for a limit theorem for the (rescaled) height or contour process of a $\mu$-GaltonWatson tree conditioned to have a large fixed number of vertices. The previous results strongly suggest that the limiting process should be a normalized (i.e. conditioned to have length equal to 1) excursion of the height process $H$. This is indeed true under suitable assumptions: When $\mu$ is critical with finite variance, this was proved by Aldous [3] in the case of the contour process and the limit is a normalized Brownian excursion as expected. Aldous' result has been extended by Duquesne [10] to the case when $\mu$ is in the domain of attraction of a stable law of index $\gamma \in(1,2]$. In this more general setting, the limit is the normalized excursion of the stable height process, which is discussed in Section 3.5 below.

### 2.6 The convergence of reduced trees

Consider a $\mu$-Galton-Watson tree, which describes the genealogy of a Galton-Watson process with offspring distribution $\mu$ starting with one individual at time 0 . For every integer $n \geq 1$, denote by $P^{(n)}$ the conditional probability knowing that the process is not extinct at time $n$, or equivalently the height of the tree is at least $n$. Under $P^{(n)}$, we can consider the reduced tree that consists only of those individuals in the generations up to time $n$ that have descendants at generation $n$. The results of the previous sections can be used to investigate the limiting behavior of these reduced trees when $n$ tends to $\infty$, even in the more general setting where the offspring distribution depends on $n$.

Here, we will concentrate on the population of the reduced tree at every generation. For every $k \in\{0,1, \ldots, n\}$, we denote by $Z_{k}^{n}$ the number of individuals in the tree at generation $k$ which have descendants at generation $n$. Obviously, $k \rightarrow Z_{k}^{n}$ is nondecreasing, $Z_{0}^{n}=1$ and $Z_{n}^{n}$ is equal to the number of individuals in the original tree at generation $n$. If $g$ denotes the generating function of $\mu$ and $g_{n}, n \geq 0$ are the iterates of $g$, it is easy to verify that $\left(Z_{k}^{n}, 0 \leq k \leq n\right)$ is a time-inhomogeneous Markov chain whose transition kernels are characterized by:

$$
E^{(n)}\left[r^{Z_{k+1}^{n}} \mid Z_{k}^{n}\right]=\left(\frac{g\left(r\left(1-g_{n-k-1}(0)\right)+g_{n-k-1}(0)\right)-g_{n-k}(0)}{1-g_{n-k}(0)}\right)^{Z_{k}^{n}}, \quad 0 \leq k<n
$$

The process $\left(Z_{k}^{n}, 0 \leq k \leq n\right)$ (under the probability measure $P^{(n)}$ ) will be called the reduced process of the $\mu$-Galton-Watson tree at generation $n$. It is easy to see that for every $k \in\{0,1, \ldots, n-1\}, Z_{k}^{n}$ can be written as a simple functional of the height process of the tree: $Z_{k}^{n}$ counts the number of excursions of the height process above level $k$ that hit level $n$.

Consider as in the previous sections a sequence ( $\mu_{p}, p=1,2, \ldots$ ) of (sub)critical offspring distributions, and for every integer $n \geq 1$ let $Z^{(p), n}=\left(Z_{k}^{(p), n}, 0 \leq k \leq n\right)$ be the reduced process of the $\mu_{p}$-Galton-Watson tree at generation $n$. For every $T>0$, we denote by $N_{(T)}$ the conditional probability $N\left(\cdot \mid \sup \left\{H_{s}, s \geq 0\right\} \geq T\right)$ (this makes sense provided that the condition (2.27) holds, cf Corollary 1.4.2).

Theorem 2.6.1 Suppose that the assumptions of Theorem 2.3.1 hold and let $T>0$. Then,

$$
\left(Z_{\left[\gamma_{p} t\right]}^{(p),\left[\gamma_{p} T\right]}, 0 \leq t<T\right) \underset{p \rightarrow \infty}{\stackrel{(f d)}{\longrightarrow}}\left(Z_{t}^{T}, 0 \leq t<T\right)
$$

where the limiting process $\left(Z_{t}^{T}, 0 \leq t<T\right)$ is defined under $N_{(T)}$ as follows: For every $t \in[0, T), Z_{t}^{T}$ is the number of excursions of $H$ above level $t$ that hit level $T$.

A more explicit description of the limiting process and of the associated tree will be given in the next section.
Proof. We use the notation of the proof of Proposition 2.5.2. In particular, the height process of the $\mu_{p}$-Galton-Watson tree conditioned on non-extinction at generation $\left[\gamma_{p} T\right]$ is $\left(\tilde{H}_{k}^{p}, k \geq 0\right)$ and the associated rescaled process is $\tilde{H}_{s}^{(p)}=\gamma_{p}^{-1} \tilde{H}_{\left[p \gamma_{p} s\right]}^{p}$. We may and will assume that $\tilde{H}_{s}^{(p)}$ is given by the formula

$$
\tilde{H}_{s}^{(p)}=H_{\left(G_{T}^{(p)}+s\right) \wedge D_{T}^{(p)}}^{(p)}
$$

and that $\left(\tilde{H}_{s}^{(p)}, s \geq 0\right)$ converges a.s. in the sense of the Skorokhod topology, towards the process $\tilde{H}_{s}=H_{\left(G_{T}+s\right) \wedge D_{T}}$ whose law is the distribution of $H$ under $N_{(T)}$.

Now we observe that the reduced process $Z_{\left[\gamma_{p} t\right]}^{(p),\left[\gamma_{p} T\right]}$ can be expressed in terms of $\tilde{H}^{(p)}$. More precisely, it is clear by construction that for every $k \in\left\{0,1, \ldots,\left[\gamma_{p} T\right]-\right.$ $1\}, Z_{k}^{(p),\left[\gamma_{p} T\right]}$ is the number of excursions of $\tilde{H}^{p}$ above level $k$ that hit level $\left[\gamma_{p} T\right]$. Equivalently, for every $t$ such that $\left[\gamma_{p} t\right]<\left[\gamma_{p} T\right]$,

$$
\tilde{Z}_{t}^{(p)}:=Z_{\left[\gamma_{p} t\right]}^{(p),\left[\gamma_{p} T\right]}
$$

is the number of excursions of $\tilde{H}^{(p)}$ above level $\left[\gamma_{p} t\right] / \gamma_{p}$ that hit level $\left[\gamma_{p} T\right] / \gamma_{p}$.
Let $t>0$. Using the fact that $t$, resp. $T$, is a.s. not a local minimum, resp maximum, of $H$ (Lemma 2.5.3), it is easy to deduce from the convergence $\tilde{H}^{(p)} \rightarrow \tilde{H}$ that the number of excursions of $\tilde{H}^{(p)}$ above level $\left[\gamma_{p} t\right] / \gamma_{p}$ that hit level $\left[\gamma_{p} T\right] / \gamma_{p}$ converges a.s. to the number of excursions of $\tilde{H}$ above level $t$ that hit level $T$. In other words, $\tilde{Z}_{t}^{(p)}$ converges a.s. to $Z_{t}^{T}$. This completes the proof.

### 2.7 The law of the limiting reduced tree

In this section, we will describe the law of the process $\left(Z_{t}^{T}, 0 \leq t<T\right)$ of the previous section, and more precisely the law of the underlying branching tree. We suppose that the Lévy process $X$ satisfies (2.27) in addition to (H1) - (H3). The random variable $Z_{t}^{T}$ (considered under the probability measure $N_{(T)}$ ) counts the number of excursions of $H$ above level $t$ that hit level $T$.

Before stating our result, we recall the notation of Section 1.4. For every $\lambda>0$ and $t>0$,

$$
u_{t}(\lambda)=N\left(1-\exp \left(-\lambda L_{\sigma}^{t}\right)\right)
$$

solves the integral equation

$$
u_{t}(\lambda)+\int_{0}^{t} \psi\left(u_{s}(\lambda)\right) d s=\lambda
$$

and

$$
v(t)=u_{t}(\infty)=N\left(L_{\sigma}^{t}>0\right)=N\left(\sup _{s \geq 0} H_{s}>t\right)
$$

is determined by

$$
\int_{v(t)}^{\infty} \frac{d x}{\psi(x)}=t
$$

Note the composition property $u_{t} \circ u_{s}=u_{t+s}$, and in particular $u_{t}(v(r))=v(t+r)$.

Theorem 2.7.1 Under $N_{(T)}$, the process $\left(Z_{t}^{T}, 0 \leq t<T\right)$ is a time-inhomogeneous Markov process whose law is characterized by the following identities: For every $\lambda>0$,

$$
\begin{equation*}
N_{(T)}\left[\exp -\lambda Z_{t}^{T}\right]=1-\frac{u_{t}\left(\left(1-e^{-\lambda}\right) v(T-t)\right)}{v(T)} \tag{2.41}
\end{equation*}
$$

and if $0 \leq t<t^{\prime}<T$,

$$
\begin{equation*}
N_{(T)}\left[\exp -\lambda Z_{t^{\prime}}^{T} \mid Z_{t}^{T}\right]=\left(N_{(T-t)}\left[\exp -\lambda Z_{t^{\prime}-t}^{T-t}\right]\right)^{Z_{t}^{T}} \tag{2.42}
\end{equation*}
$$

Alternatively, we can describe the law of the process $\left(Z_{t}^{T}, 0 \leq t<T\right)$ under $N_{(T)}$ by the following properties.

- $Z_{r}^{T}=1$ if and only if $r \in\left[0, \gamma_{T}\right)$, where the law of $\gamma_{T}$ is given by

$$
\begin{equation*}
N_{(T)}\left[\gamma_{T}>t\right]=\frac{\tilde{\psi}(v(T))}{\tilde{\psi}(v(T-t))}, \quad 0 \leq t<T \tag{2.43}
\end{equation*}
$$

where $\tilde{\psi}(x)=x^{-1} \psi(x)$.

- The conditional distribution of $Z_{\gamma_{T}}^{T}$ knowing $\gamma_{T}$ is characterized by

$$
\begin{equation*}
N_{(T)}\left[r^{Z_{\gamma_{T}}^{T}} \mid \gamma_{T}=t\right]=r \frac{\psi^{\prime}(U)-\gamma_{\psi}(U,(1-r) U)}{\psi^{\prime}(U)-\gamma_{\psi}(U, 0)}, \quad 0 \leq r \leq 1 \tag{2.44}
\end{equation*}
$$

where $U=v(T-t)$ and for every $a, b \geq 0$,

$$
\gamma_{\psi}(a, b)= \begin{cases}(\psi(a)-\psi(b)) /(a-b) & \text { if } a \neq b \\ \psi^{\prime}(a) & \text { if } a=b\end{cases}
$$

- Conditionally on $\gamma_{T}=t$ and $Z_{\gamma_{T}}^{T}=k$, the process $\left(Z_{t+r}^{T}, 0 \leq r<T-t\right)$ is distributed as the sum of $k$ independent copies of the process $\left(Z_{r}^{T-t}, 0 \leq r<\right.$ $T-t)$ under $N_{(T-t)}$.

Proof. One can give several approaches to Theorem 2.7.1. In particular, the timeinhomogeneous Markov property could be deduced from the analogous result for discrete reduced trees by using Theorem 2.6.1. We will prefer to give a direct approach relying on the properties of the height process.

Before stating a key lemma, we introduce some notation. We fix $t \in(0, T)$. Note that the definition of $Z_{t}^{T}$ also makes sense under the conditional probability $N_{(t)}$. We denote by $\left(e_{i}^{t}, i=1, \ldots, Z_{t}^{T}\right)$ the successive excursions of $H$ above level $t$ that hit level $T-t$, shifted in space and time so that each starts from 0 at time 0 . Recall the notation $L_{s}^{a}$ for the local times of the height process. We also write $L_{(i)}^{t}$ for the local time of $H$ at level $t$ at the beginning of excursion $e_{i}^{t}$.

Lemma 2.7.2 Under $N_{(t)}$, conditionally on the local time $L_{\sigma}^{t}$, the point measure

$$
\sum_{i=1}^{Z_{t}^{T}} \delta_{\left(L_{(i)}^{t}, e_{i}^{t}\right)}
$$

is Poisson with intensity $1_{\left[0, L_{\sigma}^{t}\right]}(\ell) d \ell N\left(d e \cap\left\{\sup H_{s}>T-t\right\}\right)$. In particular, under $N_{(t)}$ or under $N_{(T)}$, conditionally on $Z_{t}^{T}$, the excursions $\left(e_{i}^{t}, i=1, \ldots, Z_{t}^{T}\right)$ are independent with distribution $N_{(T-t)}$.

Proof. We rely on Proposition 1.3.1 and use the notation of Chapter 1. Under the probability measure $P$, denote by $f_{i}^{t}, i=1,2, \ldots$ the successive excursions of $H$ above level $t$ that hit $T$, and let $\ell_{i}^{t}$ be the local time of $H$ at level $t$ at the beginning (or the end) of excursion $f_{i}^{t}$. Then the $f_{i}^{t}$ 's are also the successive excursions of the process $H_{s}^{t}=H\left(\rho_{s}^{t}\right)$ that hit level $T-t$, and the numbers $\ell_{i}^{t}$ are the corresponding local times (of $H^{t}$ ) at level 0. By Proposition 1.3.1 and excursion theory, the point measure

$$
\sum_{i=1}^{\infty} \delta_{\left(\ell_{i}^{t}, f_{i}^{t}\right)}
$$

is Poisson with intensity $d \ell N\left(d f \cap\left\{\sup H_{s}>T-t\right\}\right)$ and is independent of the $\sigma$-field $\mathcal{H}_{t}$.

On the other hand, let $\lambda_{1}$ be the local time of $H$ at level $t$ at the end of the first excursion of $H$ away from 0 that hits level $t$. From the approximation of local time provided by Proposition 1.3.3, it is easy to see that $\lambda_{1}$ is $\mathcal{H}_{t}$-measurable. By excursion theory for $X-I$, the law under $N_{(t)}$ of the pair

$$
\left(L_{\sigma}^{t}, \sum_{i=1}^{Z_{t}^{T}} \delta_{\left(L_{(i)}^{t}, e_{i}^{t}\right)}\right)
$$

is the same as the law under $P$ of

$$
\left(\lambda_{1}, \sum_{\left\{i: \ell_{i}^{t} \leq \lambda_{1}\right\}} \delta_{\left(\ell_{i}^{t}, f_{i}^{t}\right)}\right) .
$$

The first assertion of the lemma now follows from the preceding considerations.
The second assertion stated under $N_{(t)}$ is an immediate consequence of the first one. The statement under $N_{(T)}$ follows since $N_{(T)}=N_{(t)}\left(\cdot \mid Z_{t}^{T} \geq 1\right)$.

We return to the proof of Theorem 2.7.1. Note that (2.42) is an immediate consequence of the second assertion of the lemma. Let us prove (2.41). By the first assertion of the lemma, $Z_{t}^{T}$ is Poisson with intensity $v(T-t) L_{\sigma}^{t}$, conditionally on $L_{\sigma}^{t}$, under $N_{(t)}$. Hence,

$$
\begin{aligned}
N_{(t)}\left[e^{-\lambda Z_{t}^{T}}\right] & =N_{(t)}\left[e^{-L_{\sigma}^{t} v(T-t)\left(1-e^{-\lambda}\right)}\right] \\
& =1-\frac{1}{v(t)} N\left(1-e^{-L_{\sigma}^{t} v(T-t)\left(1-e^{-\lambda}\right)}\right) \\
& =1-\frac{1}{v(t)} u_{t}\left(\left(1-e^{-\lambda}\right) v(T-t)\right)
\end{aligned}
$$

Then observe that

$$
N_{(t)}\left[1-e^{-\lambda Z_{t}^{T}}\right]=\frac{1}{v(t)} N\left(1-e^{-\lambda Z_{t}^{T}}\right)=\frac{v(T)}{v(t)} N_{(T)}\left[1-e^{-\lambda Z_{t}^{T}}\right] .
$$

Formula (2.41) follows immediately.
It is clear that there exists a random variable $\gamma_{T}$ such that $Z_{t}^{T}=1$ iff $0 \leq t<\gamma_{T}$, $N_{(T)}$ a.s. $\left(\gamma_{T}\right.$ is the minimum of the height process between the first and the last hitting time of $T$ ). Let us prove (2.43). By (2.41), we have,

$$
N_{(T)}\left[\gamma_{T}>t\right]=\lim _{\lambda \rightarrow \infty} e^{\lambda} N_{(T)}\left[e^{-\lambda Z_{t}^{T}}\right]=\lim _{\lambda \rightarrow \infty} e^{\lambda}\left(1-\frac{u_{t}\left(\left(1-e^{-\lambda}\right) v(T-t)\right)}{v(T)}\right) .
$$

Recalling that $u_{t}(v(T-t))=v(T)$, we have as $\varepsilon \rightarrow 0$,

$$
u_{t}((1-\varepsilon) v(T-t))=v(T)-\varepsilon v(T-t) \frac{\partial u_{t}}{\partial \lambda}(v(T-t))+o(\varepsilon)
$$

and it follows that

$$
N_{(T)}\left[\gamma_{T}>t\right]=\frac{v(T-t)}{v(T)} \frac{\partial u_{t}}{\partial \lambda}(v(T-t)) .
$$

Formula (2.43) follows from that identity and the fact that, for $\lambda>0$,

$$
\begin{equation*}
\frac{\partial u_{t}}{\partial \lambda}(\lambda)=\frac{\psi\left(u_{t}(\lambda)\right)}{\psi(\lambda)} \tag{2.45}
\end{equation*}
$$

To verify (2.45), differentiate the integral equation for $u_{t}(\lambda)$ :

$$
\frac{\partial u_{t}}{\partial \lambda}(\lambda)=1-\int_{0}^{t} \frac{\partial u_{s}}{\partial \lambda}(\lambda) \psi^{\prime}\left(u_{s}(\lambda)\right) d s
$$

which implies

$$
\frac{\partial u_{t}}{\partial \lambda}(\lambda)=\exp \left(-\int_{0}^{t} \psi^{\prime}\left(u_{s}(\lambda)\right) d s\right)
$$

Then note that $\frac{\partial}{\partial t} \log \psi\left(u_{t}(\lambda)\right)=-\psi^{\prime}\left(u_{t}(\lambda)\right)$ and thus

$$
\int_{0}^{t} \psi^{\prime}\left(u_{s}(\lambda)\right) d s=\log \psi\left(u_{t}(\lambda)\right)-\log \psi(\lambda)
$$

This completes the proof of (2.45) and (2.43).
We now prove the last assertion of the theorem. Recall the notation introduced before Lemma 2.7.2. Clearly it suffices to prove that the following property holds:
(P) Under $N_{(T)}$, conditionally on $\gamma_{T}=t$ and $Z_{\gamma_{T}}^{T}=n$, the excursions $e_{1}^{\gamma_{T}}, \ldots, e_{n}^{\gamma_{T}}$ are i.i.d. according to the distribution $N_{(T-t)}$.

We can deduce property $(\mathrm{P})$ from Lemma 2.7 .2 via an approximation procedure. Let us sketch the argument. For any $p \geq 2$ and any bounded continuous functional $F$ on $\mathbb{R}_{+} \times C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)^{p}$,

$$
\begin{align*}
& N_{(T)}\left[1_{\left\{Z_{\gamma_{T}}^{T}=p\right\}} F\left(\gamma_{T}, e_{1}^{\gamma_{T}}, \ldots, e_{p}^{\gamma_{T}}\right)\right] \\
& =\lim _{n \rightarrow \infty} \sum_{j=1}^{n-1} N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ;(j-1) T / n<\gamma_{T} \leq j T / n\right\}} F\left(\frac{j T}{n}, e_{1}^{j T / n}, \ldots, e_{p}^{j T / n}\right)\right] \tag{2.46}
\end{align*}
$$

Note that the event $\left\{\gamma_{T} \leq j T / n\right\}$ contains $\left\{Z_{j T / n}^{T}=p\right\}$. As a consequence of the second part of Lemma 2.7.2 (applied with $t=j T / n$ ) we have

$$
\begin{aligned}
& N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ; \gamma_{T} \leq j T / n\right\}} F\left(\frac{j T}{n}, e_{1}^{j T / n}, \ldots, e_{p}^{j T / n}\right)\right] \\
& =N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ; \gamma_{T} \leq j T / n\right\}}\right. \\
& \left.\quad \int N_{(T-j T / n)}\left(d f_{1}\right) \ldots N_{(T-j T / n)}\left(d f_{p}\right) F\left(\frac{j T}{n}, f_{1}, \ldots, f_{p}\right)\right] .
\end{aligned}
$$

We want to get a similar identity where the event $\left\{\gamma_{T} \leq j T / n\right\}$ is replaced by $\left\{\gamma_{T} \leq(j-1) T / n\right\}=\left\{Z_{(j-1) T / n}^{T} \geq 2\right\}$. A slightly more complicated argument (relying
on two applications of Lemma 2.7.2, the first one with $t=(j-1) T / n$ and then with $t=T / n)$ shows similarly that

$$
\begin{aligned}
& N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ; \gamma_{T} \leq(j-1) T / n\right\}} F\left(\frac{j T}{n}, e_{1}^{j T / n}, \ldots, e_{p}^{j T / n}\right)\right] \\
& =N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ; \gamma_{T} \leq(j-1) T / n\right\}}\right. \\
& \left.\quad \times \int N_{(T-j T / n)}\left(d f_{1}\right) \ldots N_{(T-j T / n)}\left(d f_{p}\right) F\left(\frac{j T}{n}, f_{1}, \ldots, f_{p}\right)\right] .
\end{aligned}
$$

By making the difference between the last two displays, we see that the sum in the right side of (2.46) exactly equals

$$
\begin{aligned}
& \sum_{j=1}^{n-1} N_{(T)}\left[1_{\left\{Z_{j T / n}^{T}=p ;(j-1) T / n<\gamma_{T} \leq j T / n\right\}}\right. \\
& \left.\quad \times \int N_{(T-j T / n)}\left(d f_{1}\right) \ldots N_{(T-j T / n)}\left(d f_{p}\right) F\left(\frac{j T}{n}, f_{1}, \ldots, f_{p}\right)\right]
\end{aligned}
$$

Using an easy continuity property of the mapping $r \rightarrow N_{(r)}$, we get from this and (2.46) that

$$
\begin{aligned}
& N_{(T)}\left[1_{\left\{Z_{\gamma_{T}}^{T}=p\right\}} F\left(\gamma_{T}, e_{1}^{\gamma_{T}}, \ldots, e_{p}^{\gamma_{T}}\right)\right] \\
& =N_{(T)}\left[1_{\left\{Z_{\gamma_{T}}^{T}=p\right\}} \int N_{\left(T-\gamma_{T}\right)}\left(d f_{1}\right) \ldots N_{\left(T-\gamma_{T}\right)}\left(d f_{p}\right) F\left(\gamma_{T}, f_{1}, \ldots, f_{p}\right)\right]
\end{aligned}
$$

which completes the proof of property ( P ) and of the last assertion of the theorem.
We finally verify (2.44). First observe from (2.43) that the density of the law of $\gamma_{T}$ under $N_{(T)}$ is given by

$$
h_{T}(t)=\tilde{\psi}(v(T)) h(T-t)
$$

where

$$
h(t)=\frac{v(t) \psi^{\prime}(v(t))}{\psi(v(t))}-1 .
$$

On the other hand, fix $\delta \in(0, T)$, and note that $\left\{\gamma_{T}>\delta\right\}=\left\{Z_{\delta}^{T}=1\right\}$. By the last assertion of Lemma 2.7.2 we have for any nonnegative function $f$,

$$
N_{(T)}\left[f\left(\gamma_{T}, Z_{\gamma_{T}}^{T}\right) 1_{\left\{\gamma_{T}>\delta\right\}} \mid \gamma_{T}>\delta\right]=N_{(T-\delta)}\left[f\left(\gamma_{T-\delta}+\delta, Z_{\gamma_{T-\delta}}^{T-\delta}\right)\right] .
$$

Hence, if $\left(\theta_{t}^{T}(k), k=2,3, \ldots\right), 0<t<T$ denotes a regular version of the conditional law of $Z_{\gamma_{T}}^{T}$ knowing that $\gamma_{T}=t$, we have

$$
\begin{aligned}
\int_{\delta}^{T} d t h(T-t) \sum_{k=2}^{\infty} \theta_{t}^{T}(k) f(t, k) & =\int_{0}^{T-\delta} d t h(T-\delta-t) \sum_{k=2}^{\infty} \theta_{t}^{T-\delta}(k) f(t+\delta, k) \\
& =\int_{\delta}^{T} d t h(T-t) \sum_{k=2}^{\infty} \theta_{t-\delta}^{T-\delta}(k) f(t, k)
\end{aligned}
$$

This shows that we must have $\theta_{t}^{T}=\theta_{t-\delta}^{T-\delta}$ for a.a. $t \in(\delta, T)$. By simple arguments, we can choose the regular versions $\theta_{t}^{T}(k)$ in such a way that $\theta_{t}^{T}(k)=\theta_{T-t}(k)$ for every $k \geq 2, T>0$ and $t \in(0, T)$.

We can then compute $N_{(T)}\left[e^{-\lambda L_{\sigma}^{T}}\right]$ in two different ways. First,

$$
N_{(T)}\left[e^{-\lambda L_{\sigma}^{T}}\right]=1-\frac{N\left(1-e^{-\lambda L_{\sigma}^{T}}\right)}{v(T)}=1-\frac{u_{T}(\lambda)}{v(T)} .
$$

Then, using property (P) once again,

$$
\begin{aligned}
N_{(T)}\left[e^{-\lambda L_{\sigma}^{T}}\right] & =N_{(T)}\left[\left(N_{(T-t)}\left[e^{-\lambda L_{\sigma}^{T-t}}\right]_{t=\gamma_{T}}\right)^{Z_{\gamma_{T}}^{T}}\right] \\
& =\int_{0}^{T} d t h_{T}(t) \sum_{k=2}^{\infty} \theta_{T-t}(k)\left(1-\frac{u_{T-t}(\lambda)}{v(T-t)}\right)^{k} .
\end{aligned}
$$

By comparing with the previous display and using the formula for $h_{T}(t)$, we get

$$
\int_{0}^{T} d t\left(\frac{\psi^{\prime}(v(t)) v(t)}{\psi(v(t))}-1\right) \sum_{k=2}^{\infty} \theta_{t}(k)\left(1-\frac{u_{t}(\lambda)}{v(t)}\right)^{k}=\frac{v(T)-u_{T}(\lambda)}{\psi(v(T))}
$$

We can now differentiate with respect to $T$ (for a proper justification we should argue that the mapping $t \rightarrow \theta_{t}$ is continuous, but we omit details). It follows that

$$
\begin{aligned}
& \left(\frac{\psi^{\prime}(v(T)) v(T)}{\psi(v(T))}-1\right) \sum_{k=2}^{\infty} \theta_{T}(k)\left(1-\frac{u_{T}(\lambda)}{v(T)}\right)^{k} \\
& \quad=-1+\frac{\psi^{\prime}(v(T)) v(T)}{\psi(v(T))}+\frac{\psi\left(u_{T}(\lambda)\right)-u_{T}(\lambda) \psi^{\prime}(v(T))}{\psi(v(T))} .
\end{aligned}
$$

Hence,

$$
\sum_{k=2}^{\infty} \theta_{T}(k)\left(1-\frac{u_{T}(\lambda)}{v(T)}\right)^{k}=1-\frac{\psi\left(u_{T}(\lambda)\right)-u_{T}(\lambda) \psi^{\prime}(v(T))}{\psi(v(T))-v(T) \psi^{\prime}(v(T))}
$$

If we substitute $r=1-\frac{u_{T}(\lambda)}{v(T)}$ in this last identity we get

$$
\sum_{k=2}^{\infty} \theta_{T}(k) r^{k}=1-\frac{\psi((1-r) v(T))-(1-r) v(T) \psi^{\prime}(v(T))}{\psi(v(T))-v(T) \psi(v(T)}
$$

Formula (2.44) follows after straightforward transformations of the last expression.
The proof of Theorem 2.7.1 is now complete. Observe that the (time-inhomogeneous) Markov property of the process $\left(Z_{t}^{T}, 0 \leq t<T\right)$ is a consequence of the description provided in the second part of the theorem, and in particular of the special form of the law of $\gamma_{T}$ and the fact that the law of $Z_{\gamma_{T}}^{T}$ under $N_{(T)}\left[\cdot \mid \gamma_{T}>\delta\right]$ coincides with the law of $Z_{\gamma_{T-\delta}}^{T-\delta}$ under $N_{(T-\delta)}$.

Let us discuss special cases of the theorem. When $\psi(u)=c u^{\alpha}$, with $c>0$ and $1<\alpha \leq 2$, we have $v(t)=(c(\alpha-1) t)^{-1 /(\alpha-1)}$, and formula (2.43) shows that the law of $\gamma_{T}$ is uniform over $[0, T]$. This is the only case where this property holds: If we assume that $\gamma_{T}$ is uniform over $[0, T]$, (2.43) implies that $\tilde{\psi}(v(t))=C / t$ for some $C>0$. By differentiating $\log v(t)$, we then get that $v(t)=C^{\prime} t^{-C}$ and it follows that $\psi$ is of the desired form.

Also in the stable case $\psi(u)=c u^{\alpha}$, formula (2.44) implies that $Z_{\gamma_{T}}^{T}$ is independent of $\gamma_{T}$, and that its distribution is characterized by

$$
N_{(T)}\left[r^{Z_{\gamma_{T}}^{T}}\right]=\frac{(1-r)^{\alpha}-1+\alpha r}{\alpha-1} .
$$

Of course when $\alpha=2$, we recover the well known fact that $Z_{\gamma_{T}}^{T}=2$. When $\alpha \in(1,2)$, we get

$$
N_{(T)}\left[Z_{\gamma_{T}}^{T}=k\right]=\frac{\alpha(2-\alpha)(3-\alpha) \cdots(k-1-\alpha)}{k!}, \quad k \geq 2 .
$$

To conclude let us mention that limiting reduced trees have been studied extensively in the literature. In the finite variance case, the uniform distribution for $\gamma_{T}$ appears in Zubkov [50], and the full structure of the reduced tree is derived by Fleischmann and Siegmund-Schultze [17]. Analogous results in the stable case (and in the more general setting of multitype branching processes) can be found in Vatutin [48] and Yakymiv [49]

## Chapter 3

## Marginals of continuous trees

### 3.1 Duality properties of the exploration process

In this section, we study certain duality properties of the process $\rho$. In view of forthcoming applications, the main result is the time-reversal property stated in Corollary 3.1.6 below. However the intermediate results needed to derive this property are of independent interest.

We work in the general setting of Chapter 1. In particular, the Lévy process $X$ satisfies assumptions (H1) - (H3), and starts at 0 under the probability measure $P$. Since the subordinator $S_{L^{-1}(t)}$ has drift $\beta$ (Lemma 1.1.2), it readily follows from formula (1.10) that the continuous part of $\rho_{t}$ is $\beta 1_{\left[0, H_{t}\right]}(r) d r$. We can thus rewrite Definition 1.2.2 in an equivalent way as follows:

$$
\begin{equation*}
\rho_{t}(d r)=\beta 1_{\left[0, H_{t}\right]}(r) d r+\sum_{\substack{0<s \leq t \\ X_{s-}<I_{t}^{s}}}\left(I_{t}^{s}-X_{s-}\right) \delta_{H_{s}}(d r) \tag{3.1}
\end{equation*}
$$

We then introduce another measure-valued process $\left(\eta_{t}, t \geq 0\right)$ by setting

$$
\begin{equation*}
\eta_{t}(d r)=\beta 1_{\left[0, H_{t}\right]}(r) d r+\sum_{\substack{0<s \leq t \\ X_{s-}<I_{t}^{s}}}\left(X_{s}-I_{t}^{s}\right) \delta_{H_{s}}(d r) \tag{3.2}
\end{equation*}
$$

In the same way as $\rho_{t}$, the measure $\eta_{t}$ is supported on $\left[0, H_{t}\right]$. We will see below that $\eta_{t}$ is a.s. a finite measure, a fact that is not obvious from the previous formula. In the queueing system interpretation of [32], the measure $\rho_{t}$ accounts for the remaining service times for all customers present in the queue at time $t$. In this interpretation, $\eta_{t}$ describes the services already accomplished for these customers.

We will see that in some sense, the process $\left(\eta_{t}, t \geq 0\right)$ is the dual of ( $\rho_{t}, t \geq 0$ ). It turns out that the study of $\left(\eta_{t}, t \geq 0\right)$ is significantly more difficult than that of $\left(\rho_{t}, t \geq 0\right)$. We start with a basic lemma.

Lemma 3.1.1 For each fixed value of $t>0$, we have $\left\langle\eta_{t}, 1\right\rangle<\infty$, P a.s. or $N$ a.e. The process $\left(\eta_{t}, t \geq 0\right)$, which takes values in $M_{f}\left(\mathbb{R}_{+}\right)$, is right-continuous in probability under $P$. Similarly, $\left(\eta_{t}, t>0\right)$ is right-continuous in measure under $N$.

Proof. Let us first prove that $\left\langle\eta_{t}, 1\right\rangle<\infty, P$ a.s. It is enough to verify that

$$
\sum_{0<s \leq t} \Delta X_{s} 1_{\left\{X_{s-}<I_{t}^{s}\right\}}<\infty
$$

$P$ a.s. By time-reversal, this is equivalent to

$$
\begin{equation*}
\sum_{0<s \leq t} \Delta X_{s} 1_{\left\{X_{s}>S_{s-\}}\right\}}<\infty \tag{3.3}
\end{equation*}
$$

$P$ a.s. However, for every $a>0$,

$$
\begin{aligned}
E\left[\sum_{0<s \leq L^{-1}(a)}\left(\Delta X_{s} \wedge 1\right) 1_{\left\{X_{s}>S_{s-}\right\}}\right] & \left.=a N^{*}\left(\left(\Delta X_{\sigma}\right) \wedge 1\right) 1_{\left\{X_{\sigma}>0\right\}}\right) \\
& =a \int \pi(d x) \int_{0}^{x} d z(z \wedge 1) \\
& \leq a \int \pi(d x)\left(x \wedge x^{2}\right) \\
& <\infty
\end{aligned}
$$

using (1.3) in the second equality. This gives our claim (3.3) and the first assertion of the lemma under $P$. The property $\left\langle\eta_{t}, 1\right\rangle<\infty, N$ a.e., then follows from arguments of excursion theory, using in particular the Markov property of $X$ under $N$.

The preceding considerations also imply that

$$
\lim _{t \downarrow 0} \sum_{0<s \leq t} \Delta X_{s} 1_{\left\{X_{s}>S_{s-}\right\}}=0
$$

in $P$-probability. Via time-reversal, it follows that the process $\eta_{t}$ is right-continuous at $t=0$ in probability under $P$. Then let $t_{0}>0$. We first observe that $\eta_{t_{0}}\left(\left\{H_{t_{0}}\right\}\right)=0 P$ a.s. This follows from the fact that there is a.s. no value of $s \in\left(0, t_{0}\right]$ with $X_{s}>S_{s-}$ and $L_{s}=0$. Then, for $t>t_{0}$, write $u=u(t)$ for the (first) time of the minimum of $X$ over $\left[t_{0}, t\right]$. Formula (3.2) implies that $\eta_{t}$ is bounded below by the restriction of $\eta_{t_{0}}$ to $\left[0, H_{u}\right)$, and bounded above by $\eta_{t_{0}}+\widetilde{\eta}_{t-t_{0}}^{(t)}$, where $\left\langle\widetilde{\eta}_{t-t_{0}}^{(t)}, 1\right\rangle$ has the same distribution as $\left\langle\eta_{t-t_{0}}, 1\right\rangle$ (more precisely, $\widetilde{\eta}_{t-t_{0}}^{(t)}$ is distributed as $\eta_{t-t_{0}}$, up to a translation by $H_{u}$ ). The right-continuity in $P$-probability of the mapping $t \rightarrow \eta_{t}$ at $t=t_{0}$ follows from this observation, the property $\eta_{t_{0}}\left(\left\{H_{t_{0}}\right\}\right)=0$, the a.s. lower semi-continuity of $H_{t}$, and the case $t_{0}=0$.

The right-continuity in measure under $N$ follows from the same arguments.
Rather than investigating the Markovian properties of $\left(\eta_{t}, t \geq 0\right)$ we will consider the pair $\left(\rho_{t}, \eta_{t}\right)$. We first introduce some notation. Let $(\mu, \nu) \in M_{f}\left(\mathbb{R}_{+}\right)^{2}$, and let $a>0$. Recall the notation of Proposition 1.2.3. In a way analogous to Chapter 1, we define $k_{a}(\mu, \nu) \in M_{f}\left(\mathbb{R}_{+}\right)^{2}$ by setting

$$
k_{a}(\mu, \nu)=(\bar{\mu}, \bar{\nu})
$$

where $\bar{\mu}=k_{a} \mu$ and the measure $\bar{\nu}$ is the unique element of $M_{f}\left(\mathbb{R}_{+}\right)$such that

$$
(\mu+\nu)_{\mid\left[0, H\left(k_{a} \mu\right)\right]}=k_{a} \mu+\bar{\nu} .
$$

Note that the difference $\mu_{\left[0, H\left(k_{a} \mu\right)\right]}-k_{a} \mu$ is a nonnegative multiple of the Dirac measure at $H\left(k_{a} \mu\right)$, so that $\bar{\nu}$ and $\nu_{\left[0, H\left(k_{a} \mu\right)\right]}$ may only differ at the point $H\left(k_{a} \mu\right)$.

Then, if $\theta_{1}=\left(\mu_{1}, \nu_{1}\right) \in M_{f}\left(\mathbb{R}_{+}\right)^{2}$ and $\theta_{2}=\left(\mu_{2}, \nu_{2}\right) \in M_{f}\left(\mathbb{R}_{+}\right)^{2}$, and if $H\left(\mu_{1}\right)<$ $\infty$, we define the concatenation $\left[\theta_{1}, \theta_{2}\right]$ by

$$
\left[\theta_{1}, \theta_{2}\right]=\left(\left[\mu_{1}, \mu_{2}\right], \nu\right)
$$

where $\langle\nu, f\rangle=\int \nu_{1}(d s) 1_{\left[0, H\left(\mu_{1}\right)\right]}(s) f(s)+\int \nu_{2}(d s) f\left(H\left(\mu_{1}\right)+s\right)$.

Proposition 3.1.2 (i) Let $s \geq 0$ and $t>0$. Then, for every nonnegative measurable function $f$ on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$,

$$
E\left[f\left(\rho_{s+t}, \eta_{s+t}\right) \mid \mathcal{G}_{s}\right]=\Pi_{t}^{0} f\left(\rho_{s}, \eta_{s}\right)
$$

where $\Pi_{t}^{0}\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)$ is the distribution of the pair

$$
\left[k_{-I_{t}}(\mu, \nu),\left(\rho_{t}, \eta_{t}\right)\right]
$$

under $P$. The collection $\left(\Pi_{t}^{0}, t>0\right)$ is a Markovian semigroup on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$.
(ii) Let $s>0$ and $t>0$. Then, for every nonnegative measurable function $f$ on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$,

$$
N\left(f\left(\rho_{s+t}, \eta_{s+t}\right) 1_{\{s+t<\sigma\}} \mid \mathcal{G}_{s}\right)=1_{\{s<\sigma\}} \Pi_{t} f\left(\rho_{s}, \eta_{s}\right)
$$

where $\Pi_{t}\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)$ is the distribution of the pair

$$
\left[k_{-I_{t}}(\mu, \nu),\left(\rho_{t}, \eta_{t}\right)\right]
$$

under $P\left(\cdot \cap\left\{T_{<\mu, 1>}>t\right\}\right)$. The collection $\left(\Pi_{t}, t>0\right)$ is a submarkovian semigroup on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$.

Proof. (i) Recall the notation of the proof of Proposition 1.2.3, and in particular formula (1.13). According to this formula, we have

$$
\begin{equation*}
\rho_{s+t}=\left[k_{-I_{t}^{(s)}} \rho_{s}, \rho_{t}^{(s)}\right] \tag{3.4}
\end{equation*}
$$

where the pair $\left(I_{t}^{(s)}, \rho_{t}^{(s)}\right)$ is defined in terms of the shifted process $X^{(s)}$, which is independent of $\mathcal{G}_{s}$. We then want to get an analogous expression for $\eta_{t}$. Precisely, we claim that

$$
\begin{equation*}
\left(\rho_{s+t}, \eta_{s+t}\right)=\left[k_{-I_{t}^{(s)}}\left(\rho_{s}, \eta_{s}\right),\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)\right] \tag{3.5}
\end{equation*}
$$

with an obvious notation. Note that (3.4) is the equality of the first components in (3.5).

To deal with the second components, recall the definition of $\eta_{s+t}$

$$
\eta_{s+t}(d u)=\beta 1_{\left[0, H_{s+t}\right]}(u) d u+\sum_{\substack{0<r \leq s+t \\ X_{r-}<I_{s+t}^{r}}}\left(X_{r}-I_{s+t}^{r}\right) \delta_{H_{r}}(d u) .
$$

First consider the absolutely continuous part. By (3.4), we have

$$
H_{s+t}=H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)+H\left(\rho_{t}^{(s)}\right)=H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)+H_{t}^{(s)}
$$

and thus

$$
\begin{aligned}
& \int d u 1_{\left[0, H_{s+t]}\right.}(u) f(u) \\
& \quad=\int d u 1_{\left[0, H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)\right]}(u) f(u)+\int d u 1_{\left[0, H_{t}^{(s)}\right]}(u) f\left(H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)+u\right) .
\end{aligned}
$$

This shows that the absolutely continuous part of $\eta_{s+t}$ is the same as that of the second component of the right side of (3.5).

Then the singular part of $\eta_{s+t}$ is equal to

$$
\begin{equation*}
\sum_{\substack{0<r \leq s \\ X_{r-}<I_{s+t}^{r}}}\left(X_{r}-I_{s+t}^{r}\right) \delta_{H_{r}}+\sum_{\substack{s<r \leq s+t \\ X_{r-}<I_{s+t}^{r}}}\left(X_{r}-I_{s+t}^{r}\right) \delta_{H_{r}} . \tag{3.6}
\end{equation*}
$$

Note that, if $r \in(s, s+t]$ is such that $X_{r-}<I_{s+t}^{r}$, we have $H_{r}=H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)+H_{r-s}^{(s)}$ (see the proof of Proposition 1.2.3). Thanks to this remark, we see that the second term of the sum in (3.6) is the image of the singular part of $\eta_{t}^{(s)}$ under the mapping $u \rightarrow H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)+u$.

To handle the first term of (3.6), we consider two cases. Suppose first that $I_{s}<$ $I_{s+t}^{s}$. Then set

$$
v=\sup \left\{r \in(0, s]: X_{r-}<I_{s+t}^{s}\right\} .
$$

In the first term of (3.6), we need only consider values $r \in(0, v]$. Note that $H_{v}=$ $H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)$ and that the measures $\rho_{v}$ and $k_{-I_{t}^{(s)}} \rho_{s}$ are equal except possibly at the point $H_{v}$ (see again the proof of Proposition 1.2.3). Then,

$$
\sum_{\substack{0<r<v \\ X_{r-}<I_{s+t}^{r}}}\left(X_{r}-I_{s+t}^{r}\right) \delta_{H_{r}}=\sum_{\substack{0<r<v \\ X_{r-}<I_{s}^{r}}}\left(X_{r}-I_{s}^{r}\right) \delta_{H_{r}}
$$

coincides with the restriction of the singular part of $\eta_{s}$ to $\left[0, H_{v}\right)=\left[0, H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)\right)$. On the other hand, $\eta_{s+t}\left(\left\{H_{v}\right\}\right)$ is equal to

$$
X_{v}-I_{s+t}^{v}=\eta_{s}\left(\left\{H_{v}\right\}\right)+\rho_{s}\left(\left[0, H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)\right]\right)-\left\langle k_{-I_{t}^{(s)}} \rho_{s}, 1\right\rangle
$$

since by construction

$$
\begin{aligned}
& \eta_{s}\left(\left\{H_{v}\right\}\right)=X_{v}-I_{s}^{v} \\
& \rho_{s}\left(\left[0, H\left(k_{-I_{t}^{(s)}} \rho_{s}\right)\right]\right)=\rho_{s}\left(\left[0, H_{v}\right]\right)=I_{s}^{v}-I_{s} \\
& \left\langle k_{-I_{t}^{(s)}} \rho_{s}, 1\right\rangle=X_{s}-I_{s}+I_{t}^{(s)}=I_{s+t}^{s}-I_{s}=I_{s+t}^{v}-I_{s}
\end{aligned}
$$

By comparing with the definition of $k_{a}(\mu, \nu)$, we see that the proof of (3.5) is complete in the case $I_{s}<I_{s+t}^{s}$.

The case $I_{s} \geq I_{s+t}^{s}$ is easier. In that case $k_{-I_{t}^{(s)}} \rho_{s}=0$, and even $k_{-I_{t}^{(s)}}\left(\rho_{s}, \eta_{s}\right)=$ $(0,0)$ (note that $\eta_{s}$ gives no mass to 0 , a.s.). Furthermore, the first sum in (3.6) vanishes, and it immediately follows that (3.5) holds.

The first assertion in (i) is a consequence of (3.5) and the fact that $X^{(s)}$ is independent of $\mathcal{G}_{s}$.

As for the second assertion, it is enough to verify that, for every $s, t>0$ we have

$$
\begin{equation*}
\left[k_{-I_{t}^{(s)}}\left[k_{-I_{s}}(\mu, \nu),\left(\rho_{s}, \eta_{s}\right)\right],\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)\right]=\left[k_{-I_{s+t}}(\mu, \nu),\left(\rho_{s+t}, \eta_{s+t}\right)\right] \tag{3.7}
\end{equation*}
$$

Note that the case $\mu=\nu=0$ is just (3.5). To prove (3.7), we consider the same two cases as previously.

If $I_{s+t}^{s}>I_{s}$, or equivalently $-I_{t}^{(s)}<\left\langle\rho_{s}, 1\right\rangle$, then $I_{s}=I_{s+t}$ and so $k_{-I_{s}}(\mu, \nu)=$ $k_{-I_{s+t}}(\mu, \nu)$. Furthermore, it is easy to verify that a.s.

$$
k_{-I_{t}^{(s)}}\left[k_{-I_{s}}(\mu, \nu),\left(\rho_{s}, \eta_{s}\right)\right]=\left[k_{-I_{s}}(\mu, \nu), k_{-I_{t}^{(s)}}\left(\rho_{s}, \eta_{s}\right)\right] .
$$

Hence

$$
\begin{aligned}
{\left[k_{-I_{t}^{(s)}}\left[k_{-I_{s}}(\mu, \nu),\left(\rho_{s}, \eta_{s}\right)\right],\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)\right] } & =\left[\left[k_{-I_{s}}(\mu, \nu), k_{-I_{t}^{(s)}}\left(\rho_{s}, \eta_{s}\right)\right],\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)\right] \\
& =\left[k_{-I_{s}}(\mu, \nu),\left[k_{-I_{t}^{(s)}}\left(\rho_{s}, \eta_{s}\right),\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)\right]\right]
\end{aligned}
$$

and (3.7) follows from (3.5).
Finally, if $I_{s+t}^{s} \leq I_{s}$, or equivalently $-I_{t}^{(s)}>\left\langle\rho_{s}, 1\right\rangle$, it easily follows from our definitions (and from the fact that $\eta_{s}(\{0\})=0$ a.s.) that

$$
k_{-I_{t}^{(s)}}\left[k_{-I_{s}}(\mu, \nu),\left(\rho_{s}, \eta_{s}\right)\right]=k_{-I_{s+t}}(\mu, \nu), \quad \text { a.s. }
$$

Furthermore, the property $I_{s+t}^{s} \leq I_{s}$ also implies that $\left(\rho_{t}^{(s)}, \eta_{t}^{(s)}\right)=\left(\rho_{s+t}, \eta_{s+t}\right)$, and this completes the proof of (3.7).
(ii) First note that, for $s, t>0$, the identity (3.5) also holds $N$ a.e. on $\{s+t<\sigma\}$ with the same proof (the argument is even simpler as we do not need to consider the case $I_{s+t}^{s} \leq I_{s}$ ). Also observe that $N$ a.e. on $\{s<\sigma\}$, the condition $s+t<\sigma$ holds iff $-I_{t}^{(s)}<X_{s}=\left\langle\rho_{s}, 1\right\rangle$, or equivalently $t<T_{\left\langle\rho_{s}, 1\right\rangle}^{(s)}=\inf \left\{r \geq 0: X_{r}^{(s)}=-\left\langle\rho_{s}, 1\right\rangle\right\}$. The first assertion in (ii) follows from these observations and the Markov property under $N$.

The second assertion in (ii) follows from (3.7) and the fact that

$$
\begin{aligned}
\left\{T_{<\mu, 1>}>s+t\right\} & =\left\{I_{s+t}>-\langle\mu, 1\rangle\right\} \\
& =\left\{I_{s}>-\langle\mu, 1\rangle\right\} \cap\left\{I_{t}^{(s)}>-\langle\mu, 1\rangle-X_{s}\right\} \\
& =\left\{I_{s}>-\langle\mu, 1\rangle\right\} \cap\left\{T_{<k_{-I_{s}} \mu+\rho_{s}, 1>}^{(s)}>t\right\}
\end{aligned}
$$

The previous proposition shows that the process $\left(\rho_{s}, \eta_{s}\right)$ is Markovian under $P$. We now proceed to investigate its invariant measure.

Let $\mathcal{N}(d s d \ell d x)$ be a Poisson point measure on $\left(\mathbb{R}_{+}\right)^{3}$ with intensity

$$
d s \pi(d \ell) 1_{[0, \ell]}(x) d x
$$

For every $a>0$, we denote by $\mathbb{M}_{a}$ the law on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$ of the pair $\left(\mu_{a}, \nu_{a}\right)$ defined by

$$
\begin{aligned}
& \left\langle\mu_{a}, f\right\rangle=\int \mathcal{N}(d s d \ell d x) 1_{[0, a]}(s) x f(s)+\beta \int_{0}^{a} d s f(s) \\
& \left\langle\nu_{a}, f\right\rangle=\int \mathcal{N}(d s d \ell d x) 1_{[0, a]}(s)(\ell-x) f(s)+\beta \int_{0}^{a} d s f(s)
\end{aligned}
$$

Note that $\mathbb{M}_{a}$ is invariant under the symmetry $(\mu, \nu) \rightarrow(\nu, \mu)$. We also set

$$
\mathbb{M}=\int_{0}^{\infty} d a e^{-\alpha a} \mathbb{M}_{a}
$$

The marginals of $\mathbb{M}$ coincide with the measure $\mathbf{M}$ of Chapter 1 .
Proposition 3.1.3 Let $\Phi$ be a nonnegative measurable function on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$. Then,

$$
N\left(\int_{0}^{\sigma} d t \Phi\left(\rho_{t}, \eta_{t}\right)\right)=\int \mathbb{M}(d \mu d \nu) \Phi(\mu, \nu)
$$

Proof. This is an extension of Proposition 1.2.5 and the proof is much analogous. Consider (under $P$ ) the countable collection of instants $s_{i}, i \in I$ such that $X_{s_{i}}>S_{s_{i}}$. It follows from (1.3) that

$$
\begin{equation*}
\left(L_{\infty}, \sum_{i \in I} \delta_{\left(L_{s_{i}}, \Delta X_{s_{i}}, X_{s_{i}}-S_{s_{i}-}\right)}(d s d \ell d x)\right) \stackrel{(\mathrm{d})}{=}\left(\zeta, 1_{[0, \zeta]}(s) \mathcal{N}(d s d \ell d x)\right) \tag{3.8}
\end{equation*}
$$

where $\zeta$ is an exponential variable with parameter $\alpha$ independent of $\mathcal{N}(\zeta=\infty$ if $\alpha=0)$. Recall from Chapter 1 the definition of the time-reversed process $\widehat{X}^{(t)}$. As in (1.10), we can rewrite the definition of $\rho_{t}$ and $\eta_{t}$ in terms of the reversed process $\widehat{X}^{(t)}$ :

$$
\begin{gathered}
\left\langle\rho_{t}, f\right\rangle=\beta \int_{0}^{\hat{L}_{t}^{(t)}} d r f(r)+\sum_{\substack{0<s \leq t \\
\widehat{X}_{s}^{(t)}>\widehat{S}_{s-}^{(t)}}}\left(\widehat{X}_{s}^{(t)}-\widehat{S}_{s-}^{(t)}\right) f\left(\widehat{L}_{t}^{(t)}-\widehat{L}_{s}^{(t)}\right), \\
\left\langle\eta_{t}, f\right\rangle=\beta \int_{0}^{\hat{L}_{t}^{(t)}} d r f(r)+\sum_{\substack{0<s \leq t \\
\widehat{X}_{s}^{(t)}>\widehat{S}_{s-}^{(t)}}}\left(\widehat{S}_{s-}^{(t)}-\widehat{X}_{s-}^{(t)}\right) f\left(\widehat{L}_{t}^{(t)}-\widehat{L}_{s}^{(t)}\right) .
\end{gathered}
$$

Hence we can write $\left(\rho_{t}, \eta_{t}\right)=\Gamma\left(\widehat{X}_{s \wedge t}^{(t)}, s \geq 0\right)$ with a measurable functional $\Gamma$ that is made explicit in the previous formulas. Proposition 1.1.4 now gives

$$
N\left(\int_{0}^{\sigma} d t \Phi\left(\rho_{t}, \eta_{t}\right)\right)=E\left[\int_{0}^{L_{\infty}} d a \Phi \circ \Gamma\left(X_{s \wedge L^{-1}(a)}, s \geq 0\right)\right] .
$$

However, $\Gamma\left(X_{s \wedge L^{-1}(a)}, s \geq 0\right)=\left(\bar{\mu}_{a}, \bar{\nu}_{a}\right)$, with

$$
\begin{aligned}
& \left\langle\bar{\mu}_{a}, f\right\rangle=\beta \int_{0}^{a} d r f(r)+\sum_{i \in I} 1_{[0, a]}\left(L_{s_{i}}\right)\left(X_{s_{i}}-S_{s_{i}-}\right) f\left(a-L_{s_{i}}\right) \\
& \left\langle\bar{\nu}_{a}, f\right\rangle=\beta \int_{0}^{a} d r f(r)+\sum_{i \in I} 1_{[0, a]}\left(L_{s_{i}}\right)\left(\Delta X_{s_{i}}-\left(X_{s_{i}}-S_{s_{i}-}\right)\right) f\left(a-L_{s_{i}}\right)
\end{aligned}
$$

Now use (3.8) to complete the proof.
For every $t>0$, we denote by $\widehat{\Pi}_{t}$ the image of the kernel $\Pi_{t}$ under the symmetry $(\mu, \nu) \rightarrow(\nu, \mu)$, that is

$$
\widehat{\Pi}_{t} \Phi(\mu, \nu)=\int \Pi_{t}\left((\nu, \mu), d \nu^{\prime} d \mu^{\prime}\right) \Phi\left(\mu^{\prime}, \nu^{\prime}\right)
$$

Theorem 3.1.4 The kernels $\Pi_{t}$ and $\widehat{\Pi}_{t}$ are in duality under $\mathbb{M}$.
This means that for any nonnegative measurable functions $\Phi$ and $\Psi$ on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$,

$$
\mathbb{M}\left(\Phi \Pi_{t} \Psi\right)=\mathbb{M}\left(\Psi \widehat{\Pi}_{t} \Phi\right)
$$

Proof. We first consider the potential kernels

$$
U=\int_{0}^{\infty} d t \Pi_{t}, \widehat{U}=\int_{0}^{\infty} d t \widehat{\Pi}_{t}
$$

and we prove that

$$
\begin{equation*}
\mathbb{M}(\Phi U \Psi)=\mathbb{M}(\Psi \widehat{U} \Phi) \tag{3.9}
\end{equation*}
$$

This is equivalent to saying that the measure

$$
\mathbb{M}(d \mu d \nu) U\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)
$$

is invariant under the transformation $\left(\mu, \nu, \mu^{\prime}, \nu^{\prime}\right) \rightarrow\left(\nu^{\prime}, \mu^{\prime}, \nu, \mu\right)$.
To this end, we first derive an explicit expression for the kernel $U$. By the definition of the kernels $\Pi_{t}$, we have

$$
U \Phi(\mu, \nu)=E\left[\int_{0}^{T_{<\mu, 1>}} d t \Phi\left(\left[k_{-I_{t}}(\mu, \nu),\left(\rho_{t}, \eta_{t}\right)\right]\right)\right]
$$

This is computed in a way similar to the proof of Proposition 1.2.6, using Proposition 3.1.3 in place of Proposition 1.2.5. It follows that

$$
\begin{equation*}
U \Phi(\mu, \nu)=\int_{0}^{<\mu, 1>} d r \int \mathbb{M}\left(d \mu^{\prime} d \nu^{\prime}\right) \Phi\left(\left[k_{r}(\mu, \nu),\left(\mu^{\prime}, \nu^{\prime}\right)\right]\right) \tag{3.10}
\end{equation*}
$$

We then need to get more information about the joint distribution of $\left((\mu, \nu), k_{r}(\mu, \nu)\right)$ under $\mathbb{M}(d \mu d \nu) 1_{[0,<\mu, 1>]}(r) d r$. Recall the notation $\mathcal{N}, \mu_{a}, \nu_{a}$ introduced before the statement of Proposition 3.1.3. Write

$$
\mathcal{N}=\sum_{i \in I} \delta_{\left(s_{i}, \ell_{i}, x_{i}\right)}
$$

for definiteness, in such a way that

$$
\left(\mu_{a}, \nu_{a}\right)=\left(\beta m_{a}+\sum_{s_{i} \leq a} x_{i} \delta_{s_{i}}, \beta m_{a}+\sum_{s_{i} \leq a}\left(\ell_{i}-x_{i}\right) \delta_{s_{i}}\right),
$$

where $m_{a}$ denotes Lebesgue measure on $[0, a]$. Since $\mathbb{M}_{a}$ is the law of $\left(\mu_{a}, \nu_{a}\right)$, we get

$$
\begin{align*}
& \int \mathbb{M}_{a}(d \mu d \nu) \int_{0}^{<\mu, 1>} d r F\left((\mu, \nu), k_{r}(\mu, \nu)\right)  \tag{3.11}\\
& \quad=E\left[\int_{0}^{<\mu_{a}, 1>} d r F\left(\left(\mu_{a}, \nu_{a}\right), k_{r}\left(\mu_{a}, \nu_{a}\right)\right)\right] \\
& \quad=E\left[\beta \int_{0}^{a} d s F\left(\left(\mu_{a}, \nu_{a}\right),\left(\mu_{a \mid[0, s]}, \nu_{a \mid[0, s]}\right)\right)\right] \\
& \quad+E\left[\sum_{s_{i} \leq a} \int_{0}^{x_{i}} d y F\left(\left(\mu_{a}, \nu_{a}\right),\left(\mu_{a \mid\left[0, s_{i}\right)}+y \delta_{s_{i}}, \nu_{a \mid\left[0, s_{i}\right)}+\left(\ell_{i}-y\right) \delta_{s_{i}}\right)\right)\right]
\end{align*}
$$

using the definition of $k_{r}$.
At this point, we recall the following well-known lemma about Poisson measures.
Lemma 3.1.5 Let $E$ be a measurable space and let $\Delta$ be a $\sigma$-finite measure on $E$. Let $\mathcal{M}$ be a Poisson point measure on $[0, a] \times E$ with intensity $d s \Delta(d e)$. Then, for any nonnegative measurable function $\Phi$,

$$
E\left[\int \mathcal{M}(d s d e) \Phi((s, e), \mathcal{M})\right]=E\left[\int_{0}^{a} d s \int_{E} \Delta(d e) \Phi\left((s, e), \mathcal{M}+\delta_{(s, e)}\right)\right]
$$

Thanks to this lemma, the second term in the right side of (3.11) can be written as

$$
\begin{aligned}
& E\left[\int_{0}^{a} d s \int \pi(d \ell) \int_{0}^{\ell} d x \int_{0}^{x} d y\right. \\
& \left.\quad F\left(\left(\mu_{a}+x \delta_{s}, \nu_{a}+(\ell-x) \delta_{s}\right),\left(\mu_{a \mid[0, s)}+y \delta_{s}, \nu_{a \mid[0, s)}+(\ell-y) \delta_{s}\right)\right)\right] .
\end{aligned}
$$

We now integrate (3.11) with respect to $e^{-\alpha a} d a$. After some easy transformations, we get

$$
\begin{aligned}
& \int \mathbb{M}(d \mu d \nu) \int_{0}^{<\mu, 1>} d r F\left((\mu, \nu), k_{r}(\mu, \nu)\right) \\
& \quad=\beta \int \mathbb{M}\left(d \mu_{1} d \nu_{1}\right) \mathbb{M}\left(d \mu_{2} d \nu_{2}\right) F\left(\left[\left(\mu_{1}, \nu_{1}\right),\left(\mu_{2}, \nu_{2}\right)\right],\left(\mu_{1}, \nu_{1}\right)\right) \\
& +\int \mathbb{M}\left(d \mu_{1} d \nu_{1}\right) \mathbb{M}\left(d \mu_{2} d \nu_{2}\right) \int \pi(d \ell) \int_{0}^{\ell} d x \int_{0}^{x} d y \\
& \quad F\left(\left[\left(\mu_{1}, \nu_{1}\right),\left(x \delta_{0}+\mu_{2},(\ell-x) \delta_{0}+\nu_{2}\right)\right],\left[\left(\mu_{1}, \nu_{1}\right),\left(y \delta_{0},(\ell-y) \delta_{0}\right)\right]\right) .
\end{aligned}
$$

Recalling formula (3.10) for the potential kernel $U$, we see that the measure

$$
\mathbb{M}(d \mu d \nu) U\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)
$$

is the sum of two terms. The first one is the distribution under

$$
\beta \mathbb{M}\left(d \mu_{1} d \nu_{1}\right) \mathbb{M}\left(d \mu_{2} d \nu_{2}\right) \mathbb{M}\left(d \mu_{3} d \nu_{3}\right)
$$

of the pair

$$
(\mu, \nu)=\left[\left(\mu_{1}, \nu_{1}\right),\left(\mu_{2}, \nu_{2}\right)\right],\left(\mu^{\prime}, \nu^{\prime}\right)=\left[\left(\mu_{1}, \nu_{1}\right),\left(\mu_{3}, \nu_{3}\right)\right] .
$$

The second one is the distribution under

$$
\mathbb{M}\left(d \mu_{1} d \nu_{1}\right) \mathbb{M}\left(d \mu_{2}, d \nu_{2}\right) \mathbb{M}\left(d \mu_{3} d \nu_{3}\right) \pi(d \ell) 1_{\{0<y<x<\ell\}} d x d y
$$

of the pair

$$
(\mu, \nu)=\left[\left(\mu_{1}, \nu_{1}\right),\left(x \delta_{0}+\mu_{2},(\ell-x) \delta_{0}+\nu_{2}\right)\right],\left(\mu^{\prime}, \nu^{\prime}\right)=\left[\left(\mu_{1}, \nu_{1}\right),\left(y \delta_{0}+\mu_{3},(\ell-y) \delta_{0}+\nu_{3}\right)\right] .
$$

In this form, it is clear that $\mathbb{M}(d \mu d \nu) U\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)$ has the desired invariance property. This completes the proof of (3.9).

Consider now the resolvent kernels

$$
U_{p}\left((\mu, \nu), d \mu^{\prime} d \nu^{\prime}\right)=\int_{0}^{\infty} d t e^{-p t} \Pi_{t}\left((\mu, \nu), d \mu^{\prime}, d \nu^{\prime}\right)
$$

By a standard argument (see e.g. [9], p.54), (3.9) also implies that, for every $p>0$, $\mathbb{M}\left(\Phi U_{p} \Psi\right)=\mathbb{M}\left(\Psi \widehat{U}_{p} \Phi\right)$, or equivalently

$$
\begin{equation*}
\int_{0}^{\infty} d t e^{-p t} \mathbb{M}\left(\Phi \Pi_{t} \Psi\right)=\int_{0}^{\infty} d t e^{-p t} \mathbb{M}\left(\Psi \Pi_{t} \Phi\right) \tag{3.12}
\end{equation*}
$$

Recall that our goal is to prove the identity $\mathbb{M}\left(\Phi \Pi_{t} \Psi\right)=\mathbb{M}\left(\Psi \Pi_{t} \Phi\right)$ for every $t>0$. We may assume that the functions $\Phi$ and $\Psi$ are continuous and both dominated
by $e^{-a<\mu, 1>}$ for some $a>0$. The latter condition guarantees that $\mathbb{M}(\Phi)<\infty$ and $\mathbb{M}(\Psi)<\infty$. From the definition of $\Pi_{t}$ and the right-continuity in probability of the mapping $t \rightarrow\left(\rho_{t}, \eta_{t}\right)$ (Lemma 3.1.1), it is easy to verify that $t \rightarrow \Pi_{t} \Psi(\mu, \nu)$ is rightcontinuous over $(0, \infty)$. The same holds for the mapping $t \rightarrow \mathbb{M}\left(\Phi \Pi_{t} \Psi\right)$, and the statement of the theorem follows from (3.12).

For notational reasons, we make the convention that $\rho_{s}=\eta_{s}=0$ if $s<0$.
Corollary 3.1.6 The process $\left(\eta_{s}, s \geq 0\right)$ has a càdlàg modification under $N$ or under P. Furthermore, the processes $\left(\rho_{s}, \eta_{s} ; s \geq 0\right)$ and $\left(\eta_{(\sigma-s)-}, \rho_{(\sigma-s)-} ; s \geq 0\right)$ have the same distribution under $N$.

A consequence of the corollary is the fact that the processes $\left(H_{t}, t \geq 0\right)$ and ( $H_{(\sigma-t) \vee 0}, t \geq 0$ ) have the same distribution (say in the sense of finite-dimensional marginals when $H$ is not continuous) under $N$. In view of the results of Chapter 2 , this is not surprising, as the same time-reversal property obviously holds for the discrete contour process. The more precise statement of the corollary will be useful in the next sections.
Proof. The second part of the corollary is essentially a consequence of the duality property stated in the previous theorem. Since we have still little information about regularity properties of the process $\eta$, we will proceed with some care. We first introduce the Kuznetsov measure $\mathbb{K}$, which is the $\sigma$-finite measure on $\mathbb{R} \times \mathbb{D}\left(\mathbb{R}_{+}, \mathbb{R}\right)$ defined by

$$
\mathbb{K}(d r d \omega)=d r N(d \omega)
$$

We then define $\gamma(r, \omega)=r, \delta(r, \omega)=r+\sigma(\omega)$ and, for every $t \in \mathbb{R}$,

$$
\bar{\rho}_{t}(r, \omega)=\rho_{t-r}(\omega), \bar{\eta}_{t}(r, \omega)=\eta_{t-r}(\omega)
$$

with the convention explained before the statement of the corollary. Note that $\left(\bar{\rho}_{t}, \bar{\eta}_{t}\right) \neq(0,0)$ iff $\gamma<t<\delta$.

It readily follows from Proposition 3.1.3 that, for every $t \in \mathbb{R}$, the distribution of $\left(\bar{\rho}_{t}, \bar{\eta}_{t}\right)$ under $\mathbb{K}\left(\cdot \cap\left\{\left(\bar{\rho}_{t}, \bar{\eta}_{t}\right) \neq(0,0)\right\}\right)$ is $\mathbb{M}$. Let $t_{1}, \ldots, t_{p} \in \mathbb{R}$ with $t_{1}<t_{2}<\cdots<t_{p}$. Using Proposition 3.1.2 and induction on $p$, we easily get that the restriction to $\left(M_{f}\left(\mathbb{R}_{+}\right)^{2} \backslash\{(0,0)\}\right)^{p}$ of the distribution of the $p$-tuple $\left(\left(\bar{\rho}_{t_{1}}, \bar{\eta}_{t_{1}}\right), \ldots,\left(\bar{\rho}_{t_{p}}, \bar{\eta}_{t_{p}}\right)\right)$ is

$$
\mathbb{M}\left(d \mu_{1} d \nu_{1}\right) \Pi_{t_{2}-t_{1}}\left(\left(\mu_{1}, \nu_{1}\right), d \mu_{2} d \nu_{2}\right) \ldots \Pi_{t_{p}-t_{p-1}}\left(\left(\mu_{p-1}, \nu_{p-1}\right), d \mu_{p} d \nu_{p}\right)
$$

By Theorem 3.1.4, this measure is equal to

$$
\mathbb{M}\left(d \mu_{p} d \nu_{p}\right) \widehat{\Pi}_{t_{p}-t_{p-1}}\left(\left(\mu_{p}, \nu_{p}\right), d \mu_{p-1} d \nu_{p-1}\right) \ldots \widehat{\Pi}_{t_{2}-t_{1}}\left(\left(\mu_{2}, \nu_{2}\right), d \mu_{1} d \nu_{1}\right)
$$

Hence the two $p$-tuples $\left(\left(\bar{\rho}_{t_{1}}, \bar{\eta}_{t_{1}}\right), \ldots,\left(\bar{\rho}_{t_{p}}, \bar{\eta}_{t_{p}}\right)\right)$ and $\left(\left(\bar{\eta}_{-t_{1}}, \bar{\rho}_{-t_{1}}\right), \ldots,\left(\bar{\eta}_{-t_{p}}, \bar{\rho}_{-t_{p}}\right)\right)$ have the same distribution, in restriction to $\left(M_{f}\left(\mathbb{R}_{+}\right)^{2} \backslash\{(0,0)\}\right)^{p}$, under $\mathbb{K}$. Since $\left(\bar{\rho}_{t}, \bar{\eta}_{t}\right) \neq(0,0)$ iff $\gamma<t<\delta$, a simple argument shows that we can remove the restriction and conclude that these two $p$-tuples have the same distribution under $\mathbb{K}$. (This distribution is $\sigma$-finite except for an infinite mass at the point $(0,0)^{p}$.)

In particular, $\left(\bar{\rho}_{t_{1}}, \ldots, \bar{\rho}_{t_{p}}\right)$ and $\left(\bar{\eta}_{-t_{1}}, \ldots, \bar{\eta}_{-t_{p}}\right)$ have the same distribution under $\mathbb{K}$. Let $F$ be a bounded continuous function on $M_{f}\left(\mathbb{R}_{+}\right)^{p}$, such that $F(0, \ldots, 0)=0$. Suppose that $0<t_{1}<t_{2}<\ldots<t_{p}$ and let $u<v$. Then we have

$$
\begin{aligned}
& \mathbb{K}\left(1_{[u, v]}(\gamma) F\left(\bar{\rho}_{\gamma+t_{1}}, \ldots, \bar{\rho}_{\gamma+t_{p}}\right)\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} \sum_{k \in \mathbb{Z},} \mathbb{K}\left(1_{\left\{\bar{\rho}_{k \varepsilon}=0, \bar{\rho}_{(k+1) \varepsilon} \neq 0\right\}} F\left(\bar{\rho}_{k \varepsilon+t_{1}}, \ldots, \bar{\rho}_{k \varepsilon+t_{p}}\right)\right)
\end{aligned}
$$

and the similar formula

$$
\begin{aligned}
& \mathbb{K}\left(1_{[u, v]}(-\delta) F\left(\bar{\eta}_{\delta-t_{1}}, \ldots, \bar{\eta}_{\delta-t_{p}}\right)\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} \sum_{k \in \mathbb{Z}, k \varepsilon \in[u, v]} \mathbb{K}\left(1_{\left\{\bar{\eta}_{-k \varepsilon}=0, \bar{\eta}_{-(k+1) \varepsilon} \neq 0\right\}} F\left(\bar{\eta}_{-k \varepsilon-t_{1}}, \ldots, \bar{\eta}_{-k \varepsilon-t_{p}}\right)\right),
\end{aligned}
$$

using the right-continuity in $N$-measure of $\eta_{t}$. Hence the vectors $\left(\gamma, \bar{\rho}_{\gamma+t_{1}}, \ldots, \bar{\rho}_{\gamma+t_{p}}\right)$ and $\left(-\delta, \bar{\eta}_{\delta-t_{1}}, \ldots, \bar{\eta}_{\delta-t_{p}}\right)$ have the same distribution under $\mathbb{K}$. It follows that the processes $\left(\rho_{t}, t \geq 0\right)$ and ( $\left.\eta_{\sigma-t}, t \geq 0\right)$ have the same finite-dimensional marginals under $N$. Since we already know that $\left(\rho_{t}, t \geq 0\right)$ is càdlàg, we obtain that $\left(\eta_{t}, t \geq 0\right)$ has a càdlàg modification under $N$. The time-reversal property of the corollary follows immediately from the previous identification of finite-dimensional marginals. This property implies in particular that $\eta_{0+}=\eta_{\sigma-}=0 \mathrm{~N}$ a.e.

It remains to verify that $\left(\eta_{t}, t \geq 0\right)$ has a càdlàg modification under $P$. On each excursion interval of $X-I$ away from 0 , we can apply the result derived above under the excursion measure $N$. It remains to deal with instants $t$ such that $X_{t}=I_{t}$, for which $\eta_{t}=0$. To this end, we note that, for every $\varepsilon>0$,

$$
N\left(\sup _{s \in[0, \sigma]}\left\langle\eta_{s}, 1\right\rangle>\varepsilon\right)=N\left(\sup _{s \in[0, \sigma]}\left\langle\rho_{s}, 1\right\rangle>\varepsilon\right)<\infty .
$$

Hence, for any fixed $x>0$, we will have $\left\langle\eta_{s}, 1\right\rangle \leq \varepsilon$ for all $s \in\left[0, T_{x}\right]$ except possibly for $s$ belonging to finitely many excursion intervals of $X-I$. Together with the continuity of $\eta$ at times 0 and $\sigma$ under $N$, this implies that $P$ a.s. for every $t$ such that $X_{t}=I_{t}$, the right and left limits of $\eta_{s}$ both exist at time $t$ and vanish.

### 3.2 The tree associated with Poissonnian marks

### 3.2.1 Trees embedded in an excursion

We first give the definition of the tree associated with a continuous function $e:[a, b] \rightarrow$ $\mathbb{R}_{+}$and $p$ instants $t_{1}, \ldots, t_{p}$ with $a \leq t_{1} \leq t_{2} \leq \cdots \leq t_{p} \leq b$.

Recall from Section 0.1 the definition of a (finite) rooted ordered tree, and the notation $\mathbf{T}$ for the collection of these trees. If $v$ is an individual (a vertex) in the tree $\mathcal{T} \in \mathbf{T}$, the notation $k_{v}(\mathcal{T})$ stands for the number of children of $v$. Individuals $v$ without children, i.e. such that $k_{v}(\mathcal{T})=0$, are called leaves. For every $p \geq 1$, we denote by $\mathbf{T}_{p}$ the set of all (rooted ordered) trees with $p$ leaves.

If $\mathcal{T}^{1}, \mathcal{T}^{2}, \ldots, \mathcal{T}^{k}$ are $k$ trees, the concatenation of $\mathcal{T}^{1}, \ldots, \mathcal{T}^{k}$, which is denoted by $\left[\mathcal{T}^{1}, \mathcal{T}^{2}, \ldots, \mathcal{T}^{k}\right]$, is defined in the obvious way: For $n \geq 1,\left(i_{1}, \ldots, i_{n}\right)$ belongs to $\left[\mathcal{T}^{1}, \mathcal{T}^{2}, \ldots, \mathcal{T}^{k}\right]$ if and only if $1 \leq i_{1} \leq k$ and $\left(i_{2}, \ldots, i_{n}\right)$ belongs to $\mathcal{T}^{i_{1}}$.

A marked tree is a pair $\theta=\left(\mathcal{T},\left\{h_{v}, v \in \mathcal{T}\right\}\right)$, where $h_{v} \geq 0$ for every $v \in \mathcal{T}$. The number $h_{v}$ is interpreted as the lifetime of individual $v$, and $\mathcal{T}$ is called the skeleton of $\theta$. We denote by $\mathbb{T}_{p}$ the set of all marked trees with $p$ leaves.

Let $\theta^{1}=\left(\mathcal{T}^{1},\left\{h_{v}^{1}, v \in \mathcal{T}\right\}\right) \in \mathbb{T}_{p_{1}}, \ldots, \theta^{k}=\left(\mathcal{T}^{k},\left\{h_{v}^{k}, v \in \mathcal{T}^{k}\right\}\right) \in \mathbb{T}_{p_{k}}$, and $h \geq 0$. The concatenation $\left[\theta^{1}, \theta^{2}, \ldots, \theta^{k}\right]_{h}$ is the element of $\mathbb{T}_{p_{1}+\ldots+p_{k}}$ whose skeleton is $\left[\mathcal{T}^{1}, \mathcal{T}^{2}, \ldots, \mathcal{T}^{k}\right]$ and such that the lifetimes of vertices in $\mathcal{T}^{i}, 1 \leq i \leq k$ become the lifetimes of the corresponding vertices in $\left[\mathcal{T}^{1}, \mathcal{T}^{2}, \ldots, \mathcal{T}^{k}\right]$, and finally the lifetime of $\emptyset$ in $\left[\theta^{1}, \theta^{2}, \ldots, \theta^{k}\right]_{h}$ is $h$.

Let $a, b \in \mathbb{R}_{+}$with $a \leq b$ and let $e:[a, b] \longrightarrow \mathbb{R}_{+}$be a continuous function. For every $a \leq u \leq v \leq b$, we set

$$
m(u, v)=\inf _{u \leq t \leq v} e(t)
$$

Let $t_{1}, \ldots, t_{p} \in \mathbb{R}_{+}$be such that $a \leq t_{1} \leq t_{2} \leq \cdots \leq t_{p} \leq b$. We will now construct a marked tree

$$
\theta\left(e, t_{1}, \ldots, t_{p}\right)=\left(\mathcal{T}\left(e, t_{1}, \ldots, t_{p}\right),\left\{h_{v}\left(e, t_{1}, \ldots, t_{p}\right), v \in \mathcal{T}\right\}\right) \in \mathbb{T}_{p}
$$

associated with the function $e$ and the times $t_{1}, \ldots, t_{p}$. We proceed by induction on $p$. If $p=1, \mathcal{T}\left(e, t_{1}\right)=\{\emptyset\}$ and $h_{\emptyset}\left(e, t_{1}\right)=e\left(t_{1}\right)$.

Let $p \geq 2$ and suppose that the tree has been constructed up to order $p-1$. Then there exists an integer $k \in\{1, \ldots, p-1\}$ and $k$ integers $1 \leq i_{1}<i_{2}<\cdots<i_{k} \leq p-1$ such that $m\left(t_{i}, t_{i+1}\right)=m\left(t_{1}, t_{p}\right)$ iff $i \in\left\{i_{1}, \ldots, i_{k}\right\}$. For every $\ell \in\{0,1, \ldots, k\}$, define $e^{\ell}$ by the formulas

$$
\begin{array}{lll}
e^{0}(t)=e(t)-m\left(t_{1}, t_{p}\right), & t \in\left[t_{1}, t_{i_{1}}\right], & \\
e^{\ell}(t)=e(t)-m\left(t_{1}, t_{p}\right), & t \in\left[t_{i_{\ell}+1}, t_{i_{\ell+1}}\right], & 1 \leq \ell \leq k-1 \\
e^{k}(t)=e(t)-m\left(t_{1}, t_{p}\right), & t \in\left[t_{i_{k}+1}, t_{p}\right] . &
\end{array}
$$

We then set:

$$
\theta\left(e, t_{1}, \ldots, t_{p}\right)=\left[\theta\left(e^{0}, t_{1}, \ldots, t_{i_{1}}\right), \theta\left(e^{1}, t_{i_{1}+1}, \ldots, t_{i_{2}}\right), \ldots, \theta\left(e^{k}, t_{i_{k}+1}, \ldots, t_{p}\right)\right]_{m\left(t_{1}, t_{p}\right)}
$$

This completes the construction of the tree by induction. Note that $k+1$ is the number of children of $\emptyset$ in the tree $\theta\left(e, t_{1}, \ldots, t_{p}\right)$, and $m\left(t_{1}, t_{p}\right)$ is the lifetime of $\emptyset$.

### 3.2.2 Poissonnian marks

We consider a standard Poisson process with parameter $\lambda$ defined under the probability measure $Q_{\lambda}$. We denote by $\tau_{1} \leq \tau_{2} \leq \cdots$ the jump times of this Poisson process. Throughout this section, we argue under the measure $Q_{\lambda} \otimes N$, which means that we consider the excursion measure of $X-I$ together with independent Poissonnian marks with intensity $\lambda$ on $\mathbb{R}_{+}$. To simplify notation however, we will systematically write $N$ instead of $Q_{\lambda} \otimes N$.

Set $M=\sup \left\{i \geq 1: \tau_{i} \leq \sigma\right\}$, which represents the number of marks that fall in the excursion interval (by convention, $\sup \emptyset=0$ ). Then,

$$
N(M \geq 1)=N\left(1-e^{-\lambda \sigma}\right)=\psi^{-1}(\lambda)
$$

where the second equality follows from the fact that the Laplace exponent of the subordinator $T_{x}$ is $\psi^{-1}(\lambda)$ (see [5], Theorem VII.1).

From now on, we assume that the condition $\int_{1}^{\infty} \frac{d u}{\psi(u)}<\infty$ holds, so that $H$ has continuous sample paths (Theorem 1.4.3). We can then use subsection 3.2.1 to define the embedded tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$ under $N(\cdot \mid M \geq 1)$. Our main goal is to determine the law of this tree.

Theorem 3.2.1 Under the probability measure $N(\cdot \mid M \geq 1)$, the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$ is distributed as the family tree of a continuous-time Galton-Watson process starting with one individual at time 0 and such that:

- Lifetimes of individuals have exponential distributions with parameter $\psi^{\prime}\left(\psi^{-1}(\lambda)\right)$;
- The offspring distribution is the law of the variable $\xi$ with generating function

$$
E\left[r^{\xi}\right]=r+\frac{\psi\left((1-r) \psi^{-1}(\lambda)\right)}{\psi^{-1}(\lambda) \psi^{\prime}\left(\psi^{-1}(\lambda)\right)}
$$

Remark. As the proof will show, the theorem remains valid without the assumption that $H$ has continuous paths. We will leave this extension to the reader. Apart from some technical details, it simply requires the straightforward extension of the construction of subsection 3.2 .1 to the case when the function $e$ is only lower semicontinuous.

The proof of Theorem 3.2.1 requires a few intermediate results. To simplify notation, we will write $\tau=\tau_{1}$. We start with an important application of Corollary 3.1.6.

Lemma 3.2.2 For any nonnegative measurable function $f$ on $M_{f}\left(\mathbb{R}_{+}\right)$,

$$
N\left(f\left(\rho_{\tau}\right) 1_{\{M \geq 1\}}\right)=\lambda \int \mathbb{M}(d \mu d \nu) f(\mu) e^{-\psi^{-1}(\lambda)\langle\nu, 1\rangle}
$$

Proof. We have

$$
N\left(f\left(\rho_{\tau}\right) 1_{\{M \geq 1\}}\right)=\lambda N\left(\int_{0}^{\sigma} d t e^{-\lambda t} f\left(\rho_{t}\right)\right)=\lambda N\left(\int_{0}^{\sigma} d t e^{-\lambda(\sigma-t)} f\left(\eta_{t}\right)\right)
$$

using the time-reversal property of Corollary 3.1.6. At this point, we use the Markov property of $X$ under $N$ :

$$
N\left(\int_{0}^{\sigma} d t e^{-\lambda(\sigma-t)} f\left(\eta_{t}\right)\right)=N\left(\int_{0}^{\sigma} d t f\left(\eta_{t}\right) E_{X_{t}}\left[e^{-\lambda T_{0}}\right]\right)
$$

We have already noticed that for $x \geq 0$,

$$
E_{x}\left[e^{-\lambda T_{0}}\right]=E_{0}\left[e^{-\lambda T_{x}}\right]=e^{-x \psi^{-1}(\lambda)} .
$$

Since $X_{t}=\left\langle\rho_{t}, 1\right\rangle$ under $N$, it follows that
$N\left(f\left(\rho_{\tau}\right) 1_{\{M \geq 1\}}\right)=\lambda N\left(\int_{0}^{\sigma} d t f\left(\eta_{t}\right) e^{-\left\langle\rho_{t}, 1\right\rangle \psi^{-1}(\lambda)}\right)=\lambda \int \mathbb{M}(d \mu d \nu) f(\nu) e^{-\langle\mu, 1\rangle \psi^{-1}(\lambda)}$,
using Proposition 3.1.3. Since $\mathbb{M}$ is invariant under the mapping $(\mu, \nu) \rightarrow(\nu, \mu)$, this completes the proof.

We now set

$$
K= \begin{cases}\inf \left\{H_{s}: \tau_{1} \leq s \leq \tau_{M}\right\} & \text { if } M \geq 2 \\ \infty & \text { if } M \leq 1\end{cases}
$$

Then $K$ represents the lifetime of the ancestor in the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$ (assuming that the event $\{M \geq 2\}$ holds). To give a formula for the number of children $\xi$ of the ancestor, set

$$
\tau_{(K)}=\inf \left\{t \geq \tau: H_{t} \leq K\right\}, \tau_{(K)}^{\prime}=\inf \left\{t \geq \tau: H_{t}<K\right\}
$$

Then, again on the event $\{M \geq 2\}, \xi-1$ is the number of excursions of $H$ above level $K$, on the time interval $\left[\tau_{(K)}, \tau_{(K)}^{\prime}\right]$, which contain at least one of the Poissonnian marks. This identification follows readily from the construction of subsection 3.2.1.

The next proposition gives the joint distribution of the pair $(K, \xi)$ under $N(\cdot \cap$ $\{M \geq 2\}$ ).

Proposition 3.2.3 Let $r \in[0,1]$ and let $h$ be a nonnegative measurable function on $[0, \infty]$, with $h(\infty)=0$. Then,

$$
\begin{aligned}
& N\left(r^{\xi} h(K) \mid M \geq 1\right) \\
& \quad=\left(r \psi^{\prime}\left(\psi^{-1}(\lambda)\right)+\frac{\psi\left((1-r) \psi^{-1}(\lambda)\right)-\lambda}{\psi^{-1}(\lambda)}\right) \int_{0}^{\infty} d b h(b) e^{-b \psi^{\prime}\left(\psi^{-1}(\lambda)\right)} .
\end{aligned}
$$

The basic idea of the proof is to apply the Markov property to the process $\rho$ at time $\tau$. To this end, we need some notation. We write $\mathbf{P}_{\mu}^{*}$ for the probability measure under which $\rho$ starts at an arbitrary measure $\mu \in \mathbb{M}_{f}\left(\mathbb{R}^{d}\right)$ and is stopped when it hits 0 . As usual, $H_{t}=H\left(\rho_{t}\right)$. Under $\mathbf{P}_{\mu}^{*}$, the process $X_{t}=\left\langle\rho_{t}, 1\right\rangle$ is the underlying Lévy process started at $\langle\mu, 1\rangle$, stopped at $T_{0}=\inf \left\{t \geq 0: X_{t}=0\right\}$. We keep the notation $I_{t}$ for the minimum process of $X$. We let $\left(a_{j}, b_{j}\right), j \in J$ be the collection of excursion intervals of $X-I$ away from 0 and before time $T_{0}$. For every $j \in J$ we define the corresponding excursion by

$$
\omega_{j}(t)=X_{\left(a_{j}+t\right) \wedge b_{j}}-I_{a_{j}}, \quad t \geq 0
$$

From excursion theory, we know that the point measure

$$
\sum_{j \in J} \delta_{\left(I_{a_{j}}, \omega_{j}\right)}
$$

is Poisson under $\mathbf{P}_{\mu}^{*}$, with intensity $1_{[0,<\mu, 1>]}(u) d u N(d \omega)$ (cf the proof of Proposition 1.2.6). On the other hand, by properties of the exploration process derived in Chapter 1 , we know that $\mathbf{P}_{\mu}^{*}$ a.s. for every $s \in\left[0, T_{0}\right]$ such that $X_{s}-I_{s}=0$ (and in particular for $\left.s=a_{j}, j \in J\right)$ we have $\rho_{s}=k_{<\mu, 1>-I_{s}} \mu$ and thus $H_{s}=H\left(k_{<\mu, 1>-I_{s}} \mu\right)$. Observe also that the image of the measure $1_{[0,<\mu, 1>]}(u) d u$ under the mapping $u \rightarrow H\left(k_{<\mu, 1>-u} \mu\right)$ is exactly $\mu(d h)$. By combining these observations, we get:
(P) The point measure $\sum_{j \in J} \delta_{\left(H_{a_{j}}, \omega_{j}\right)}$ is Poisson under $\mathbf{P}_{\mu}^{*}$, with intensity $\mu(d h) N(d \omega)$.

Finally, assume that we are also given a collection $\mathcal{P}_{\lambda}$ of Poisson marks with intensity $\lambda$, independently of $\rho$ under $\mathbf{P}_{\mu}^{*}$, and set

$$
\begin{aligned}
& L=\inf \left\{H_{a_{j}}: j \in J,\left(a_{j}, b_{j}\right) \cap \mathcal{P}_{\lambda} \neq \emptyset\right\}, \quad(\inf \emptyset=\infty), \\
& \zeta=\operatorname{Card}\left\{j \in J: H_{a_{j}}=L \text { and }\left(a_{j}, b_{j}\right) \cap \mathcal{P}_{\lambda} \neq \emptyset\right\}
\end{aligned}
$$

Then the Markov property of the exploration process at time $\tau$ shows that, for any nonnegative measurable function $h$ on $[0, \infty]$ such that $h(\infty)=0$,

$$
\begin{equation*}
N\left(1_{\{M \geq 1\}} r^{\xi-1} h(K)\right)=N\left(1_{\{M \geq 1\}} \mathbf{E}_{\rho_{\tau}}^{*}\left[r^{\zeta} h(L)\right]\right) . \tag{3.13}
\end{equation*}
$$

To verify this equality, simply observe that those excursions of $H$ above level $K$ on the time interval $\left[\tau_{(K)}, \tau_{(K)}^{\prime}\right]$ that contain one Poissonnian mark, exactly correspond to those excursions of the shifted process $\left\langle\rho_{\tau+t}, 1\right\rangle$ above its minimum that start from the height $K$ and contain one mark.

The next lemma is the key step towards the proof of Proposition 3.2.3.
Lemma 3.2.4 Let $a \geq 0$ and let $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$be such that $\operatorname{supp} \mu=[0, a]$ and $\mu(d t)=\beta 1_{[0, a]}(t) d t+\mu_{s}(d t)$, where $\mu_{s}$ is a countable sum of multiples of Dirac point masses at elements of $[0, a]$. Then, if $r \in[0,1]$ and $h$ is a nonnegative measurable function on $[0, \infty]$ such that $h(\infty)=0$,

$$
\begin{align*}
& \mathbf{E}_{\mu}^{*}\left[r^{\zeta} h(L)\right]=\beta r \psi^{-1}(\lambda) \int_{0}^{a} d b e^{-\psi^{-1}(\lambda) \mu([0, b])} h(b) \\
& +\sum_{\mu(\{s\})>0}\left(e^{-(1-r) \mu(\{s\}) \psi^{-1}(\lambda)}-e^{-\mu(\{s\}) \psi^{-1}(\lambda)}\right) e^{-\mu([0, s)) \psi^{-1}(\lambda)} h(s) . \tag{3.14}
\end{align*}
$$

Proof. First note that it is easy to derive the law of $L$ under $\mathbf{P}_{\mu}^{*}$. Let $b \in[0, a]$. We have by property ( P )

$$
\begin{aligned}
\mathbf{P}_{\mu}^{*}[L>b] & =\mathbf{P}_{\mu}^{*}\left[\left(a_{j}, b_{j}\right) \cap \mathcal{P}_{\lambda}=\emptyset \text { for every } j \in J \text { s.t. } H_{a_{j}} \leq b\right] \\
& =\mathbf{E}_{\mu}^{*}\left[\exp \left(-\mu([0, b]) N\left(1-e^{-\lambda \sigma}\right)\right)\right] \\
& =\exp \left(-\mu([0, b]) \psi^{-1}(\lambda)\right) .
\end{aligned}
$$

In particular, atoms of the distribution of $L$ in $[0, \infty)$ exactly correspond to atoms of $\mu$, and the continuous part of the distribution of $L$ is the measure

$$
\beta \psi^{-1}(\lambda) \exp \left(-\mu([0, b]) \psi^{-1}(\lambda)\right) 1_{[0, a]}(b) d b
$$

We then need to distinguish two cases:
(1) Let $s \in[0, a]$ be an atom of $\mu$. By the preceding formula,

$$
\mathbf{P}_{\mu}^{*}[L=s]=\left(1-e^{\left.-\mu(\{s\}) \psi^{-1}(\lambda)\right)}\right) e^{-\mu([0, s)) \psi^{-1}(\lambda)}
$$

Note that the excursions $\omega_{j}$ that start at height $s$ are the atoms of a Poisson measure with intensity $\mu(\{s\}) N$. Using also the independence properties of Poisson measures, we get that, conditionally on $\{L=s\}, \xi$ is distributed as a Poisson random variable with intensity $\mu(\{s\}) \psi^{-1}(\lambda)$, conditioned to be greater than or equal to 1 :

$$
\mathbf{E}_{\mu}^{*}\left[r^{\xi} \mid L=s\right]=\frac{e^{\left.-(1-r) \mu(\{s\}) \psi^{-1}(\lambda)\right)}-e^{\left.-\mu(\{s\}) \psi^{-1}(\lambda)\right)}}{1-e^{\left.-\mu(\{s\}) \psi^{-1}(\lambda)\right)}}
$$

(2) If $L$ is not an atom of $\mu$, then automatically $\xi=1$. This is so because the values $H_{a_{j}}$ corresponding to indices $j$ such that $\mu\left(\left\{H_{a_{j}}\right\}\right)=0$ must be distinct, by (P) and standard properties of Poisson measures.

The lemma follows by combining these two cases with the distribution of $L$.
Proof of Proposition 3.2.3. By combining (3.13), Lemma 3.2.2 and (3.14), we obtain that

$$
N\left(1_{\{M \geq 1\}} r^{\xi-1} h(K)\right)=A_{1}+A_{2}
$$

where

$$
A_{1}=\beta r \lambda \psi^{-1}(\lambda) \int_{0}^{\infty} d a e^{-\alpha a} \int \mathbb{M}_{a}(d \mu d \nu) e^{-\langle\nu, 1\rangle \psi^{-1}(\lambda)} \int_{0}^{a} d b e^{-\psi^{-1}(\lambda) \mu([0, b])} h(b),
$$

and

$$
\begin{aligned}
A_{2}=\lambda & \int_{0}^{\infty} d a e^{-\alpha a} \int \mathbb{M}_{a}(d \mu d \nu) e^{-\langle\nu, 1\rangle \psi^{-1}(\lambda)} \\
& \quad \times \sum_{\mu(\{s\})>0}\left(e^{-(1-r) \mu(\{s\}) \psi^{-1}(\lambda)}-e^{-\mu(\{s\}) \psi^{-1}(\lambda)}\right) e^{-\mu([0, s)) \psi^{-1}(\lambda)} h(s) .
\end{aligned}
$$

To compute $A_{1}$, we observe that for $u>0$ and $0 \leq b \leq a$,

$$
\begin{aligned}
\mathbb{M}_{a}\left(e^{-u(\mu([0, b])+\nu([0, a])}\right)= & \mathbb{M}_{a}\left(e^{-u(\mu+\nu)([0, b])}\right) \mathbb{M}_{a}\left(e^{-u \nu((b, a])}\right) \\
= & e^{-\beta u(a+b)} \exp \left(-b \int \pi(d \ell) \ell\left(1-e^{-u \ell}\right)\right) \\
& \times \exp \left(-(a-b) \int \pi(d \ell) \int_{0}^{\ell} d x\left(1-e^{-u x}\right)\right) \\
= & e^{\alpha a} \exp \left(-b \psi^{\prime}(u)-(a-b) \frac{\psi(u)}{u}\right),
\end{aligned}
$$

using the easy formulas

$$
\begin{aligned}
& \int \pi(d \ell) \ell\left(1-e^{-u \ell}\right)=\psi^{\prime}(u)-\alpha-2 \beta u \\
& \int \pi(d \ell) \int_{0}^{\ell} d x\left(1-e^{-u x}\right)=\frac{1}{u}\left(\psi(u)-\alpha u-\beta u^{2}\right)
\end{aligned}
$$

It follows that

$$
\begin{aligned}
A_{1} & =\beta r \lambda \psi^{-1}(\lambda) \int_{0}^{\infty} d a \int_{0}^{a} d b h(b) e^{-b \psi^{\prime}\left(\psi^{-1}(\lambda)\right)-(a-b)\left(\lambda / \psi^{-1}(\lambda)\right)} \\
& =\beta r \psi^{-1}(\lambda)^{2} \int_{0}^{\infty} d b h(b) e^{-b \psi^{\prime}\left(\psi^{-1}(\lambda)\right)} .
\end{aligned}
$$

To evaluate $A_{2}$, first observe that, with the notation preceding Proposition 3.1.3, we have

$$
\begin{aligned}
A_{2}= & \lambda \int_{0}^{\infty} d a e^{-\alpha a} E\left[\int_{\{s \leq a\}} \mathcal{N}(d s d \ell d x) h(s)\left(e^{-(1-r) x \psi^{-1}(\lambda)}-e^{-x \psi^{-1}(\lambda)}\right)\right. \\
& \left.\times \exp \left(-\psi^{-1}(\lambda)\left(\int_{\left\{s^{\prime} \leq a\right\}} \mathcal{N}\left(d s^{\prime} d \ell^{\prime} d x^{\prime}\right)\left(\ell^{\prime}-x^{\prime}\right)+\int_{\left\{s^{\prime}<s\right\}} \mathcal{N}\left(d s^{\prime} d \ell^{\prime} d x^{\prime}\right) x^{\prime}\right)\right)\right] .
\end{aligned}
$$

From Lemma 3.1.5, it follows that

$$
\begin{aligned}
A_{2}= & \lambda \int_{0}^{\infty} d a e^{-\alpha a} \int_{0}^{a} d b h(b) \mathbb{M}_{a}\left(e^{-(\mu([0, b])+\nu([0, a])) \psi^{-1}(\lambda)}\right) \\
& \times \int \pi(d \ell) \int_{0}^{\ell} d x\left(e^{-(1-r) x \psi^{-1}(\lambda)}-e^{-x \psi^{-1}(\lambda)}\right) e^{-(\ell-x) \psi^{-1}(\lambda)} \\
= & \psi^{-1}(\lambda) \int_{0}^{\infty} d b h(b) e^{-b \psi^{\prime}\left(\psi^{-1}(\lambda)\right)} \\
& \times \int \pi(d \ell) \int_{0}^{\ell} d x\left(e^{-(1-r) x \psi^{-1}(\lambda)}-e^{-x \psi^{-1}(\lambda)}\right) e^{-(\ell-x) \psi^{-1}(\lambda)}
\end{aligned}
$$

where the last equality is obtained from the same calculations as those made in evaluating $A_{1}$. Furthermore, straightforward calculations give

$$
\begin{aligned}
& \int \pi(d \ell) \int_{0}^{\ell} d x\left(e^{-(1-r) x \psi^{-1}(\lambda)}-e^{-x \psi^{-1}(\lambda)}\right) e^{-(\ell-x) \psi^{-1}(\lambda)} \\
& \quad=\psi^{\prime}\left(\psi^{-1}(\lambda)\right)-\beta r \psi^{-1}(\lambda)+\frac{1}{r \psi^{-1}(\lambda)}\left(\psi\left((1-r) \psi^{-1}(\lambda)\right)-\lambda\right)
\end{aligned}
$$

By substituting this in the previous display and combining with the formula for $A_{1}$, we arrive at the result of the proposition.

Proof of Theorem 3.2.1. It is convenient to introduce the random variable $\Lambda$ defined by

$$
\Lambda= \begin{cases}K & \text { if } M \geq 2 \\ H_{\tau} & \text { if } M=1\end{cases}
$$

On the event $\{M=1\}$ we also set $\xi=0$. We can easily compute the law of the pair $(\Lambda, \xi)$. Indeed, by applying the Markov property at $\tau$ as previously, we easily get

$$
N\left(h(\Lambda) 1_{\{M=1\}}\right)=N\left(h\left(H_{\tau}\right) 1_{\{M=1\}}\right)
$$

$$
\begin{aligned}
& =N\left(h\left(H_{\tau}\right) e^{-<\rho_{\tau}, 1>\psi^{-1}(\lambda)}\right) \\
& \left.=\lambda \int_{0}^{\infty} d a e^{-\alpha a} h(a) \int \mathbb{M}(d \mu d \nu) e^{-(<\mu, 1>+<\nu, 1>) \psi^{-1}(\lambda)}\right) \\
& =\lambda \int_{0}^{\infty} d a h(a) e^{-\psi^{\prime}\left(\psi^{-1}(\lambda)\right) a}
\end{aligned}
$$

By combining with Proposition 3.2.3, we get

$$
\begin{aligned}
& N\left(r^{\xi} h(\Lambda) 1_{\{M \geq 1\}}\right) \\
& \quad=\left(r \psi^{-1}(\lambda) \psi^{\prime}\left(\psi^{-1}(\lambda)\right)+\psi\left((1-r) \psi^{-1}(\lambda)\right)\right) \int_{0}^{\infty} d b h(b) e^{-b \psi^{\prime}\left(\psi^{-1}(\lambda)\right)} .
\end{aligned}
$$

This formula entails that we have the following properties under $N(\cdot \mid M \geq 1)$ : The variables $\Lambda$ and $\xi$ are independent, $\Lambda$ is exponentially distributed with parameter $\psi^{\prime}\left(\psi^{-1}(\lambda)\right)$, and the generating function of $\xi$ is as stated in Theorem 3.2.1. To complete the proof, it remains to verify the "recursivity property" of the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$, that is to verify that under $N(\cdot \mid M \geq 2)$, the shifted trees corresponding to each individual in the first generation are independent and distributed as the whole tree under $N(\cdot \mid M \geq 1)$. This is a consequence of the following claim.
Claim. Let $\left(\alpha_{j}, \beta_{j}\right), j=1, \ldots, \xi$ be the excursion intervals of $H$ above level $\Lambda$ that contain at least one mark, ranked in chronological order, and for every $j=1, \ldots, \xi$ let $h_{j}(s)=H_{\left(\alpha_{j}+s\right) \wedge \beta_{j}}-\Lambda$ be the corresponding excursion. Then, conditionally on the pair $(\Lambda, \xi)$, the excursions $h_{1}, \ldots, h_{\xi}$ are independent and distributed according to the law of $\left(H_{s}, s \geq 0\right)$ under $N(\cdot \mid M \geq 1)$.

To verify this property, we first argue under $\mathbf{P}_{\mu}^{*}$ as previously. Precisely, we consider the excursions $\omega_{j}$ for all $j \in J$ such that $H_{a_{j}}=L$ and $\left(a_{j}, b_{j}\right) \cap \mathcal{P}_{\lambda} \neq \emptyset$. We denote by $\tilde{\omega}_{1}, \ldots, \tilde{\omega}_{\zeta}$ these excursions, ranked in chronological order. Then property (P) and familiar properties of Poisson measures give the following fact. For every $k \geq 1$, under the measure $\mathbf{P}_{\mu}^{*}(\cdot \mid \zeta=k)$, the excursions $\tilde{\omega}_{1}, \ldots, \tilde{\omega}_{k}$ are independent, distributed according to $N(\cdot \mid M \geq 1)$, and these excursions are also independent of the measure

$$
\sum_{H_{a_{j}}>L} \delta_{\left(I_{a_{j}}, \omega_{j}\right)}
$$

Let $\sigma_{L}:=\inf \left\{s \geq 0: H_{s}=L\right\}$. Excursion theory for $X-I$ allows us to reconstruct the process $\left(X_{s \wedge \sigma_{L}}, s \geq 0\right)$ as a measurable function of the point measure in the last display. Hence we can also assert that, under $\mathbf{P}_{\mu}^{*}(\cdot \mid \zeta=k), \tilde{\omega}_{1}, \ldots, \tilde{\omega}_{k}$ are independent of $\left(X_{s \wedge \sigma_{L}}, s \geq 0\right)$. In particular, they are independent of $L=H\left(k_{<\mu, 1>-I_{\sigma_{L}}} \mu\right)$.

We now apply these properties to the shifted process $X_{\tau+}$. under $N\left(\cdot \mid M^{L} \geq 1\right)$. We slightly abuse notation and keep denoting by $\tilde{\omega}_{1}, \ldots, \tilde{\omega}_{\zeta}$ the excursions of $X_{\tau+} .-I_{\tau+}$. that contain a mark (so that $\zeta=\xi-1$ on the event $M \geq 2$ ). By construction, for every $j \in\{2, \ldots, \xi\}$, the function $h_{j}$ is the height process of $\tilde{\omega}_{j-1}$. Hence it follows from the previous properties under $\mathbf{P}_{\mu}^{*}$ that under $N(\cdot \mid \xi=k)$ (for a fixed $k \geq 2$ ), the processes $h_{2}, \ldots, h_{k}$ are independent, have the distribution given in the claim, and are also independent of the pair $\left(h_{1}, \Lambda\right)$. Hence, for any test functions $F_{1}, \ldots, F_{k}, G$, we get

$$
\begin{aligned}
& N\left(F_{1}\left(h_{1}\right) \ldots F_{k}\left(h_{k}\right) G(\Lambda) \mid \xi=k\right) \\
& \quad=N\left(F_{2}(H) \mid M \geq 1\right) \cdots N\left(F_{k}(H) \mid M \geq 1\right) N\left(F_{1}\left(h_{1}\right) G(\Lambda) \mid \xi=k\right)
\end{aligned}
$$

Now from Corollary 3.1.6, we know that the time-reversed process $\left(H_{(\sigma-s)_{+}}, s \geq 0\right)$ has the same distribution under $N$ as the process $\left(H_{s}, s \geq 0\right)$. Furthermore, this
time-reversal operation will leave $\Lambda$ and $\xi$ invariant and transform the excursion $h_{1}$ into the time-reversal of $h_{\xi}$, denoted by $\widehat{h}_{\xi}$ (provided we do simultaneously the similar transformation on the underlying Poissonnian marks). It follows that

$$
\begin{aligned}
N\left(F_{1}\left(h_{1}\right) G(\Lambda) \mid \xi=k\right) & =N\left(F_{1}\left(\widehat{h}_{k}\right) G(\Lambda) \mid \xi=k\right) \\
& =N\left(F_{1}\left(\widehat{h}_{k}\right) \mid \xi=k\right) N(G(\Lambda) \mid \xi=k) \\
& =N\left(F_{1}(H) \mid M \geq 1\right) N(G(\Lambda) \mid M \geq 1)
\end{aligned}
$$

By substituting this equality in the previous displayed formula, we obtain the claim. This completes the proof of Theorem 3.2.1.

### 3.3 Marginals of stable trees

We first reformulate Theorem 3.2.1 in a way more suitable for our applications. Recall that $\mathbf{T}_{p}$ is the set of all (rooted ordered) trees with $p$ leaves. If $\mathcal{T} \in \mathbf{T}_{p}$ we denote by $\mathcal{L}_{\mathcal{T}}$ the set of all leaves of $\mathcal{T}$, and set $\mathcal{N}_{\mathcal{T}}=\mathcal{T} \backslash \mathcal{L}_{\mathcal{T}}$. Recall the notation $k_{v}=k_{v}(\mathcal{T})$ for the number of children of an element $v$ of $\mathcal{T}$. We write $\mathbf{T}_{p}^{*}$ for the subset of $\mathbf{T}_{p}$ composed of all trees $\mathcal{T}$ such that $k_{v}(\mathcal{T}) \geq 2$ for every $v \in \mathcal{N}_{\mathcal{T}}$. By construction, the skeleton of the marked trees $\theta\left(e, t_{1}, \ldots, t_{p}\right)$ always belongs to $\mathbf{T}_{p}^{*}$.

Theorem 3.3.1 Let $p \geq 1$. Then, for any nonnegative measurable function $\Phi$ on $\mathbb{T}_{p}$, and every $\lambda>0$,

$$
\begin{aligned}
& N\left(e^{-\lambda \sigma} \int_{\left\{t_{1}<\cdots<t_{p}<\sigma\right\}} d t_{1} \ldots d t_{p} \Phi\left(\theta\left(H, t_{1}, \ldots, t_{p}\right)\right)\right) \\
& =\sum_{\mathcal{T} \in \mathbf{T}_{p}^{*}}\left(\prod_{v \in \mathcal{N}_{\mathcal{T}}} \frac{\left|\psi^{\left(k_{v}\right)}\left(\psi^{-1}(\lambda)\right)\right|}{k_{v}!}\right) \\
& \quad \times \int \prod_{v \in \mathcal{T}} d h_{v} \exp \left(-\psi^{\prime}\left(\psi^{-1}(\lambda)\right) \sum_{v \in \mathcal{T}} h_{v}\right) \Phi\left(\mathcal{T},\left(h_{v}\right)_{v \in \mathcal{T}}\right)
\end{aligned}
$$

Proof. By elementary properties of the standard Poisson process, the left side is equal to

$$
\lambda^{-p} N\left(\Phi\left(\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)\right) 1_{\{M=p\}}\right)
$$

with the notation of the previous section. This quantity can be evaluated thanks to Theorem 3.2.1: From the generating function of the offspring distribution, we get

$$
\begin{aligned}
& P[\xi=0]=\frac{\lambda}{\psi^{-1}(\lambda) \psi^{\prime}\left(\psi^{-1}(\lambda)\right)} \\
& P[\xi=1]=0 \\
& P[\xi=k]=\frac{1}{k!} \frac{\psi^{-1}(\lambda)^{k-1}\left|\psi^{(k)}\left(\psi^{-1}(\lambda)\right)\right|}{\psi^{\prime}\left(\psi^{-1}(\lambda)\right)}, \quad \text { for every } k \geq 2
\end{aligned}
$$

Hence the probability under $N(\cdot \mid M \geq 1)$ that the skeleton of the tree $\theta\left(H, \tau_{1}, \ldots, \tau_{M}\right)$ is equal to a given tree $\mathcal{T} \in \mathbf{T}_{p}^{*}$ is

$$
\begin{aligned}
& \left(\prod_{v \in \mathcal{N}_{\mathcal{T}}} \frac{1}{k_{v}!} \frac{\psi^{-1}(\lambda)^{k_{v}-1}\left|\psi^{\left(k_{v}\right)}\left(\psi^{-1}(\lambda)\right)\right|}{\psi^{\prime}\left(\psi^{-1}(\lambda)\right)}\left(\frac{\lambda}{\psi^{-1}(\lambda) \psi^{\prime}\left(\psi^{-1}(\lambda)\right)}\right)^{p}\right. \\
& =\frac{1}{\psi^{-1}(\lambda)} \frac{1}{\psi^{\prime}\left(\psi^{-1}(\lambda)\right)^{|\mathcal{T}|}} \lambda^{p} \prod_{v \in \mathcal{N}_{\mathcal{T}}}\left(\frac{1}{k_{v}!}\left|\psi^{\left(k_{v}\right)}\left(\psi^{-1}(\lambda)\right)\right|\right)
\end{aligned}
$$

Recalling that $N(M \geq 1)=\psi^{-1}(\lambda)$, and using the fact that the lifetimes $h_{v}, v \in \mathcal{T}$ are independently distributed according to the exponential distribution with parameter $\psi^{\prime}\left(\psi^{-1}(\lambda)\right)$, we easily arrive at the formula of the theorem.

By letting $\lambda \rightarrow 0$ in the preceding theorem, we get the following corollary, which is closely related to Proposition 3.2 of [33].

Corollary 3.3.2 Suppose that $\int \pi(d r) r^{p}<\infty$. Then, for any nonnegative measurable function $\Phi$ on $\mathcal{T}_{p}$,

$$
\begin{aligned}
& N\left(\int_{\left\{t_{1}<\cdots<t_{p}<\sigma\right\}} d t_{1} \ldots t_{p} \Phi\left(\theta\left(H, t_{1}, \ldots, t_{p}\right)\right)\right) \\
& \quad=\sum_{\mathcal{T} \in \mathbf{T}_{p}^{*}}\left(\prod_{v \in \mathcal{N}_{\mathcal{T}}} \beta_{k_{v}}\right) \int \prod_{v \in \mathcal{T}} d h_{v} \exp \left(-\alpha \sum_{v \in \mathcal{T}} h_{v}\right) \Phi\left(\mathcal{T},\left(h_{v}\right)_{v \in \mathcal{T}}\right)
\end{aligned}
$$

where, for every $k=2, \ldots, p$,

$$
\beta_{k}=\frac{\left|\psi^{(k)}(0)\right|}{k!}=\beta 1_{\{k=2\}}+\frac{1}{k!} \int r^{k} \pi(d r)
$$

Remark. The formula of the corollary still holds without the assumption $\int \pi(d r) r^{p}<$ $\infty$ but it has to be interpreted properly since some of the numbers $\beta_{k}$ may be infinite.

From now on, we concentrate on the stable case $\psi(u)=u^{\gamma}$ for $1<\gamma<2$. Then the Lévy process $X$ satisfies the scaling property

$$
\left(X_{\lambda t}, t \geq 0\right) \stackrel{(\mathrm{d})}{=}\left(\lambda^{1 / \gamma} X_{t}, t \geq 0\right)
$$

under $P$. Thanks to this property, it is possible to choose a regular version of the conditional probabilities $N_{(u)}:=N(\cdot \mid \sigma=u)$ in such a way that for every $u>0$ and $\lambda>0$, the law of $\left(\lambda^{-1 / \gamma} X_{\lambda t}, t \geq 0\right)$ under $N_{(\lambda u)}$ is $N_{(u)}$. Standard arguments then show that the height process $\left(H_{s}, s \geq 0\right)$ is well defined as a continuous process under the probability measures $N_{(u)}$. Furthermore, it follows from the approximations of $H_{t}$ (see Lemma 1.1.3) that the law of $\left(H_{\lambda s}, s \geq 0\right)$ under $N_{(\lambda u)}$ is equal to the law of ( $\lambda^{1-\frac{1}{\gamma}} H_{s}, s \geq 0$ ) under $N_{(u)}$.

The probability measure $N_{(1)}$ is called the law of the normalized excursion. The tree coded by the process $\left(H_{t}, 0 \leq t \leq 1\right)$ under $N_{(1)}$ (in the sense of Section 0.6) is called the stable continuum random tree. The next theorem identifies its finitedimensional marginals, in the sense of Aldous [3].

Theorem 3.3.3 Suppose that $\psi(u)=u^{\gamma}$ for some $\gamma \in(1,2)$. Then the law of the tree $\theta\left(H, t_{1}, \ldots, t_{p}\right)$ under the probability measure

$$
p!1_{\left\{0<t_{1}<t_{2}<\ldots<t_{p}<1\right\}} d t_{1} \ldots d t_{p} N_{(1)}(d \omega)
$$

is characterized by the following properties:
(i) The probability of a given skeleton $\mathcal{T} \in \mathbf{T}_{p}^{*}$ is

$$
\frac{p!}{\prod_{v \in \mathcal{N}_{\mathcal{T}}} k_{v}!} \frac{\prod_{v \in \mathcal{N}_{\mathcal{T}}}\left|(\gamma-1)(\gamma-2) \ldots\left(\gamma-k_{v}+1\right)\right|}{(\gamma-1)(2 \gamma-1) \ldots((p-1) \gamma-1)}
$$

(ii) If $p \geq 2$, then conditionally on the skeleton $\mathcal{T}$, the lifetimes $\left(h_{v}\right)_{v \in \mathcal{T}}$ have a density with respect to Lebesgue measure on $\mathbb{R}_{+}^{\mathcal{T}}$ given by

$$
\frac{\Gamma\left(p-\frac{1}{\gamma}\right)}{\Gamma\left(\delta_{\mathcal{T}}\right)} \gamma^{|\mathcal{T}|} \int_{0}^{1} d u u^{\delta_{\mathcal{T}}-1} q\left(\gamma \sum_{v \in \mathcal{T}} h_{v}, 1-u\right)
$$

where $\delta_{\mathcal{T}}=p-\left(1-\frac{1}{\gamma}\right)|\mathcal{T}|-\frac{1}{\gamma}>0$, and $q(s, u)$ is the continuous density at time $s$ of the stable subordinator with exponent $1-\frac{1}{\gamma}$, which is characterized by

$$
\int_{0}^{\infty} d u e^{-\lambda u} q(s, u)=\exp \left(-s \lambda^{1-\frac{1}{\gamma}}\right) .
$$

If $p=1$, then $\mathcal{T}=\{\emptyset\}$ and the law of $h_{\emptyset}$ has density

$$
\gamma \Gamma\left(1-\frac{1}{\gamma}\right) q(\gamma h, 1)
$$

with respect to Lebesgue measure on $\mathbb{R}_{+}$.
Proof. For every $u>0$, let $\Theta_{(u)}$ be the law of the tree $\theta\left(H, s_{1}, \ldots, s_{p}\right)$ under the probability measure

$$
p!u^{-p} 1_{\left\{s_{1}<s_{2}<\cdots<s_{p}<u\right\}} d s_{1} \ldots d s_{p} N_{(u)}(d \omega) .
$$

By the scaling properties of the height process (see the remarks before Theorem 3.3.3), we have, for every $u>0$ and every $\mathcal{T} \in \mathbf{T}_{p}^{*}$,

$$
\Theta_{(u)}\left(\{\mathcal{T}\} \times \mathbb{R}_{+}^{\mathcal{T}}\right)=\Theta_{(1)}\left(\{\mathcal{T}\} \times \mathbb{R}_{+}^{\mathcal{T}}\right)
$$

Hence, by conditioning with respect to $\sigma$ in Theorem 3.3.1, we get

$$
\begin{equation*}
\frac{1}{p!} N\left(\sigma^{p} e^{-\lambda \sigma}\right) \Theta_{(1)}\left(\{\mathcal{T}\} \times \mathbb{R}_{+}^{\mathcal{T}}\right)=\psi^{\prime}\left(\psi^{-1}(\lambda)\right)^{-|\mathcal{T}|} \prod_{v \in \mathcal{N}_{\mathcal{T}}} \frac{\left|\psi^{\left(k_{v}\right)}\left(\psi^{-1}(\lambda)\right)\right|}{k_{v}!} \tag{3.15}
\end{equation*}
$$

From this, we can compute $\Theta_{(1)}\left(\{\mathcal{T}\} \times \mathbb{R}_{+}^{\mathcal{T}}\right)$ by observing that, for every $k \geq 1$,

$$
\psi^{(k)}\left(\psi^{-1}(\lambda)\right)=\gamma(\gamma-1) \cdots(\gamma-k+1) \lambda^{1-\frac{k}{\gamma}}
$$

and

$$
N\left(\sigma^{p} e^{-\lambda \sigma}\right)=\left|\frac{d^{p}}{d \lambda^{p}} N\left(1-e^{-\lambda \sigma}\right)\right|=\left|\frac{d^{p}}{d \lambda^{p}} \psi^{-1}(\lambda)\right|=\left|\frac{1}{\gamma}\left(\frac{1}{\gamma}-1\right) \cdots\left(\frac{1}{\gamma}-p+1\right)\right| \lambda^{\frac{1}{\gamma}-p} .
$$

If we substitute these expressions in (3.15), the terms in $\lambda$ cancel and we get part (i) of the theorem.

To prove (ii), fix $\mathcal{T} \in \mathbf{T}_{p}^{*}$, and let $D$ be a bounded Borel subset of $\mathbb{R}_{+}^{\mathcal{T}}$. Write $p_{\sigma}(d u)$ for the law of $\sigma$ under $N$. Then by applying Theorem 3.3 .1 with $\Phi=1_{\{\mathcal{T}\} \times D}$ and $\Phi=1_{\{\mathcal{T}\} \times \mathbb{R}_{+}^{\tau}}$, we get

$$
\begin{align*}
& \int p_{\sigma}(d u) e^{-\lambda u} u^{p} \Theta_{(u)}\left(\left\{h_{v}\right\}_{v \in \mathcal{T}} \in D \mid \mathcal{T}\right)  \tag{3.16}\\
& =\left(\int p_{\sigma}(d u) e^{-\lambda u} u^{p}\right) \int_{D} \prod_{v \in \mathcal{T}} d h_{v} \psi^{\prime}\left(\psi^{-1}(\lambda)\right)^{|\mathcal{T}|} \exp \left(-\psi^{\prime}\left(\psi^{-1}(\lambda)\right) \sum_{v \in \mathcal{T}} h_{v}\right) .
\end{align*}
$$

By scaling (or inverting $N\left(1-e^{-\lambda \sigma}\right)=\lambda^{1 / \gamma}$ ), we have $p_{\sigma}(d u)=c u^{-1-\frac{1}{\gamma}} d u$. It follows that

$$
\begin{aligned}
& \int_{0}^{\infty} d u e^{-\lambda u} u^{p-1-\frac{1}{\gamma}} \Theta_{(u)}\left(\left\{h_{v}\right\}_{v \in \mathcal{T}} \in D \mid \mathcal{T}\right) \\
& \quad=\frac{\gamma^{|\mathcal{T}|} \Gamma\left(p-\frac{1}{\gamma}\right)}{\lambda^{\delta_{\mathcal{T}}}} \int_{D} \prod_{v \in \mathcal{T}} d h_{v} \exp \left(-\gamma \lambda^{1-\frac{1}{\gamma}} \sum_{v \in \mathcal{T}} h_{v}\right)
\end{aligned}
$$

where $\delta_{\mathcal{T}}=p-\left(1-\frac{1}{\gamma}\right)|\mathcal{T}|-\frac{1}{\gamma}$ as in the theorem. Suppose first that $p \geq 2$. To invert the Laplace transform, observe that the right side can be written as

$$
\begin{aligned}
& \frac{\gamma^{|\mathcal{T}|} \Gamma\left(p-\frac{1}{\gamma}\right)}{\Gamma\left(\delta_{\mathcal{T}}\right)} \int_{0}^{\infty} d u e^{-\lambda u} u^{\delta_{\mathcal{T}}-1} \times \int_{0}^{\infty} d u^{\prime} e^{-\lambda u^{\prime}}\left(\int_{D} \prod_{v \in \mathcal{T}} d h_{v} q\left(\gamma \sum_{v \in \mathcal{T}} h_{v}, u^{\prime}\right)\right) \\
& =\frac{\gamma^{|\mathcal{T}|} \Gamma\left(p-\frac{1}{\gamma}\right)}{\Gamma\left(\delta_{\mathcal{T}}\right)} \int_{0}^{\infty} d u e^{-\lambda u}\left(\int_{D} \prod_{v \in \mathcal{T}} d h_{v} \int_{0}^{u} d r r^{\delta_{\mathcal{T}}-1} q\left(\gamma \sum_{v \in \mathcal{T}} h_{v}, u-r\right)\right)
\end{aligned}
$$

The first formula of (ii) now follows. In the case $p=1$, we get

$$
\int_{0}^{\infty} d u e^{-\lambda u} u^{-\frac{1}{\gamma}} \Theta_{(u)}\left(h_{\emptyset} \in D\right)=\gamma \Gamma\left(1-\frac{1}{\gamma}\right) \int_{D} d h \exp \left(-\gamma \lambda^{1-\frac{1}{\gamma}} h\right)
$$

and the stated result follows by inverting the Laplace transform.
Remarks. (a) The previous proof also readily gives the analogue of Theorem 3.3.3 in the case $\psi(u)=u^{2}$, which corresponds to the finite-dimensional marginals of Aldous' continuum random tree (see Aldous [3], or Chapter 3 of [31]). In that case, the discrete skeleton of $\theta\left(H, t_{1}, \ldots, t_{p}\right)$ is with probability one a binary tree, meaning that $k_{v}=2$ for every $v \in \mathcal{T}$. The law of $\mathcal{T}\left(H, t_{1}, \ldots, t_{p}\right)$ is the uniform probability measure on the set of all binary trees in $\mathbf{T}_{p}^{*}$, so that the probability of each possible skeleton is

$$
\frac{p!}{2^{p-1}(1 \times 3 \times \cdots \times(2 p-3))}
$$

This formula can be deduced informally by letting $\gamma$ tend to 2 in Theorem 3.3.3 (i).
To obtain the analogue of (ii), note that there is an explicit formula for $q(s, u)$ when $\psi(u)=u^{2}$ :

$$
q(s, u)=\frac{s}{2 \sqrt{\pi} u^{3 / 2}} e^{-s^{2} /(4 u)}
$$

Observe that when the skeleton is binary, we have always $|\mathcal{T}|=2 p-1$. It follows that the powers of $\lambda$ cancel in the right side of (3.16), and after straightforward calculations, we obtain that the density of $\left(h_{v}\right)_{v \in \mathcal{T}}$ on $\mathbb{R}_{+}^{2 p-1}$ is

$$
2^{2 p-1} \Gamma\left(p-\frac{1}{2}\right) q\left(2 \sum h_{v}, 1\right)=2^{p}(1 \times 3 \times \cdots \times(2 p-3))\left(\sum h_{v}\right) \exp \left(-\left(\sum h_{v}\right)^{2}\right)
$$

Compare with Aldous [3] or Chapter 3 of [31], but note that constants are different because $\psi(u)=u^{2}$ corresponds to a Brownian motion with variance $2 t$ (also the CRT is coded by twice the normalized Brownian excursion in [3]).
(b) We could get rid of the factor

$$
\frac{p!}{\prod_{v \in \mathcal{N}_{\mathcal{T}}} k_{v}!}
$$

in Theorem 3.3.3 by considering rooted (unordered) trees with $p$ labelled leaves rather than rooted ordered trees : See the discussion at the end of Chapter 3 of [31].

## Chapter 4

## The Lévy snake

### 4.1 The construction of the Lévy snake

Our goal is now to combine the branching structure studied in the previous chapters with a spatial displacement prescribed by a Markov process $\xi$. Throughout this chapter, we assume that $H$ has continuous paths (the condition $\int^{\infty} d u / \psi(u)<\infty$ holds) although many of the results can presumably be extended to the general case.

### 4.1.1 Construction of the snake with a fixed lifetime process

We consider a Markov process $\xi$ with càdlàg paths and values in a Polish space $E$, whose topology is defined by a metric $\delta$. For simplicity, we will assume that $\xi$ is defined on the canonical space $\mathbb{D}\left(\mathbb{R}_{+}, E\right)$ of càdlàg functions from $\mathbb{R}_{+}$into $E$. For every $x \in E$, we denote by $\Pi_{x}$ the distribution of $\xi$ started at $x$. It is implicitly assumed in our definition of a Markov process that the mapping $x \rightarrow \Pi_{x}$ is measurable. We also assume that $\xi$ is continuous in probability under $\Pi_{x}$ (equivalently, $\xi$ has no fixed discontinuities, $\Pi_{x}\left[\xi_{s} \neq \xi_{s-}\right]=0$ for every $s>0$ ). On the other hand, we do not assume that $\xi$ is strong Markov.

For $x \in E$, we denote by $\mathcal{W}_{x}$ the space of all $E$-valued killed paths started at $x$. An element of $\mathcal{W}_{x}$ is a càdlàg mapping $\mathrm{w}:[0, \zeta) \longrightarrow E$ such that $\mathrm{w}(0)=x$. Here $\zeta \in(0, \infty)$ is called the lifetime of the path. When there is a risk of confusion we write $\zeta=\zeta_{\mathrm{w}}$. Note that we do not require the existence of the left limit $\mathrm{w}(\zeta-)$. By convention, the point $x$ is also considered as a killed path with lifetime 0 . We set $\mathcal{W}=\cup_{x \in E} \mathcal{W}_{x}$ and equip $\mathcal{W}$ with the distance

$$
d\left(\mathrm{w}, \mathrm{w}^{\prime}\right)=\delta\left(\mathrm{w}(0), \mathrm{w}^{\prime}(0)\right)+\left|\zeta-\zeta^{\prime}\right|+\int_{0}^{\zeta \wedge \zeta^{\prime}} d t\left(d_{t}\left(\mathrm{w}_{\leq t}, \mathrm{w}_{\leq t}^{\prime}\right) \wedge 1\right)
$$

where $d_{t}$ is the Skorokhod metric on the space $\mathbb{D}([0, t], E)$, and $\mathrm{w}_{\leq t}$ denotes the restriction of w to the interval $[0, t]$. It is then elementary to check that the space $(\mathcal{W}, d)$ is a Polish space. The space $(E, \delta)$ is embedded isometrically in $\mathcal{W}$ thanks to the previous convention.

Let $x \in E$ and $\mathrm{w} \in \mathcal{W}_{x}$. If $a \in\left[0, \zeta_{\mathrm{w}}\right)$ and $b \in[a, \infty)$, we can define a probability measure $R_{a, b}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)$ on $\mathcal{W}_{x}$ by requiring that:
(i) $R_{a, b}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)$ a.s., $\mathrm{w}^{\prime}(t)=\mathrm{w}(t), \forall t \in[0, a)$;
(ii) $R_{a, b}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)$ a.s., $\zeta_{\mathrm{w}^{\prime}}=b$;
(iii) the law of ( $\left.\mathrm{w}^{\prime}(a+t), 0 \leq t<b-a\right)$ under $R_{a, b}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)$ is the law of $\left(\xi_{t}, 0 \leq t<\right.$ $b-a)$ under $\Pi_{\mathrm{w}(a-)}$.

In (iii), $\mathrm{w}(0-)=x$ by convention. In particular, $R_{0, b}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)$ is the law of $\left(\xi_{t}, 0 \leq\right.$ $t<b)$ under $\Pi_{x}$, and $R_{0,0}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right)=\delta_{x}\left(d \mathrm{w}^{\prime}\right)$.

When $\mathrm{w}\left(\zeta_{\mathrm{w}}-\right)$ exists, we may and will extend the previous definition to the case $a=\zeta_{\mathrm{w}}$.

We denote by $\left(W_{s}, s \geq 0\right)$ the canonical process on the product space $(\mathcal{W})^{\mathbb{R}_{+}}$. We will abuse notation and also write $\left(W_{s}, s \geq 0\right)$ for the canonical process on the set $C\left(\mathbb{R}_{+}, \mathcal{W}\right)$ of all continuous mappings from $\mathbb{R}_{+}$into $\mathcal{W}$. Let us fix $x \in E$ and $\mathrm{w}_{0} \in \mathcal{W}_{x}$, and let $h \in C\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$be such that $h(0)=\zeta_{\mathrm{w}_{0}}$. For $0 \leq s \leq s^{\prime}$, we set

$$
m_{h}\left(s, s^{\prime}\right)=\inf _{s \leq r \leq s^{\prime}} h(r)
$$

We assume that either $\mathrm{w}_{0}\left(\zeta_{\mathrm{w}_{0}}-\right)$ exists or $m_{h}(0, r)<h(0)$ for every $r>0$. Then, the Kolmogorov extension theorem can be used to construct the (unique) probability measure $Q_{\mathrm{w}_{0}}^{h}$ on $\left(\mathcal{W}_{x}\right)^{\mathbb{R}_{+}}$such that, for $0=s_{0}<s_{1}<\cdots<s_{n}$,

$$
\begin{aligned}
& Q_{\mathrm{w}_{0}}^{h}\left[W_{s_{0}} \in A_{0}, \ldots, W_{s_{n}} \in A_{n}\right] \\
& =1_{A_{0}}\left(\mathrm{w}_{0}\right) \int_{A_{1} \times \cdots \times A_{n}} R_{m_{h}\left(s_{0}, s_{1}\right), h\left(s_{1}\right)}\left(\mathrm{w}_{0}, d \mathrm{w}_{1}\right) \ldots R_{m_{h}\left(s_{n-1}, s_{n}\right), h\left(s_{n}\right)}\left(\mathrm{w}_{n-1}, d \mathrm{w}_{n}\right) .
\end{aligned}
$$

Notice that our assumption on the pair $\left(\mathrm{w}_{0}, h\right)$ is needed already for $n=1$ to make sense of the measure $R_{m_{h}\left(s_{0}, s_{1}\right), h\left(s_{1}\right)}\left(\mathrm{w}_{0}, d \mathrm{w}_{1}\right)$.

From the previous definition, it is clear that, for every $s<s^{\prime}, Q_{\mathrm{w}_{0}}^{h}$ a.s.,

$$
W_{s^{\prime}}(t)=W_{s}(t), \quad \forall t<m_{h}\left(s, s^{\prime}\right)
$$

and furthermore $\zeta_{W_{s}}=h(s), \zeta_{W_{s^{\prime}}}=h\left(s^{\prime}\right)$. Hence,
$d\left(W_{s}, W_{s^{\prime}}\right) \leq\left|h(s)-h\left(s^{\prime}\right)\right|+\left|\left(h(s) \wedge h\left(s^{\prime}\right)\right)-m_{h}\left(s, s^{\prime}\right)\right|=\left(h(s) \vee h\left(s^{\prime}\right)\right)-m_{h}\left(s, s^{\prime}\right)$.
From this bound, it follows that the mapping $s \longrightarrow W_{s}$ is $Q_{\mathrm{w}_{0}}^{h}$ a.s. uniformly continuous on the bounded subsets of $[0, \infty) \cap \mathbb{Q}$. Hence this mapping has $Q_{\mathrm{w}_{0}}^{h}$ a.s. a continuous extension to the positive real line. We abuse notation and still denote by $Q_{\mathrm{w}_{0}}^{h}$ the induced probability measure on $C\left(\mathbb{R}_{+}, \mathcal{W}_{x}\right)$. By an obvious continuity argument, we have $\zeta_{W_{s}}=h(s)$, for every $s \geq 0, Q_{\mathrm{w}_{0}}^{h}$ a.s., and

$$
W_{s^{\prime}}(t)=W_{s}(t), \quad \forall t<m_{h}\left(s, s^{\prime}\right), \quad \forall s<s^{\prime}, \quad Q_{\mathrm{w}_{0}}^{h} \text { a.s. }
$$

We will refer to this last property as the snake property. The process $\left(W_{s}, s \geq 0\right)$ is under $Q_{\mathrm{w}_{0}}^{h}$ a time-inhomogeneous continuous Markov process.

### 4.1.2 The definition of the Lévy snake

Following the remarks of the end of Chapter 1, we now consider the exploration process $\rho$ as a Markov process with values in the set

$$
M_{f}^{0}=\left\{\mu \in M_{f}\left(\mathbb{R}_{+}\right): H(\mu)<\infty \text { and } \operatorname{supp} \mu=[0, H(\mu)]\right\} \cup\{0\}
$$

We denote by $\mathbf{P}_{\mu}$ the law of $\left(\rho_{s}, s \geq 0\right)$ started at $\mu$. We will write indifferently $H\left(\rho_{s}\right)$ or $H_{s}$.

We then define $\Theta$ as the set of all pairs $(\mu, \mathrm{w}) \in M_{f}^{0} \times \mathcal{W}$ such that $\zeta_{\mathrm{w}}=H(\mu)$, and at least one of the following two properties hold:
(i) $\mu(\{H(\mu)\})=0$;
(ii) $\mathrm{w}\left(\zeta_{\mathrm{w}}-\right)$ exists.

We equip $\Theta$ with the product distance on $M_{f}^{0} \times \mathcal{W}$. For every $y \in E$, we also set $\Theta_{y}=\{(\mu, \mathrm{w}) \in \Theta: \mathrm{w}(0)=y\}$.

From now on until the end of this section, we fix a point $x \in E$.
Notice that when $H(\mu)>0$ and $\mu(\{H(\mu)\})=0$, we have $\inf _{[0, s]} H\left(\rho_{r}\right)<H(\mu)$ for every $s>0, \mathbf{P}_{\mu}$ a.s. To see this, note that $I_{s}<0$ for every $s>0, P$-a.s., by the regularity of 0 for $(-\infty, 0)$, for the underlying Lévy process. Then, $P$-a.s. for every $s>0$ we can pick $r \in(0, s)$ such that $X_{r}=I_{r}<0$. Formula (1.11) then shows that $\rho_{r}^{\mu}=k_{-I_{r}} \mu$ and our assumption $\mu(\{H(\mu)\})=0$ implies that $H\left(\rho_{r}^{\mu}\right)<H(\mu)$, which gives the claim.

Using the last observation and the previous subsection, we can for every $(\mu, \mathrm{w}) \in$ $\Theta_{x}$ define a probability measure $\mathbb{P}_{\mu, \mathrm{w}}$ on $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$ by the formula

$$
\mathbb{P}_{\mu, \mathrm{w}}(d \rho d W)=\mathbf{P}_{\mu}(d \rho) Q_{\mathrm{w}}^{H(\rho)}(d W),
$$

where in the right side $H(\rho)$ obviously stands for the function $\left(H\left(\rho_{s}\right), s \geq 0\right)$, which is continuous $\mathbf{P}_{\mu}$ a.s.

We will write $\mathbb{P}_{x}$ instead of $\mathbb{P}_{0, x}$ when $\mu=0$.
Proposition 4.1.1 The process $\left(\rho_{s}, W_{s}\right)$ is under $\mathbb{P}_{\mu, \mathrm{w}}$ a càdlàg Markov process in $\Theta_{x}$.
Proof. We first verify that $\mathbb{P}_{\mu, \mathrm{w}}$ a.s. the process $\left(\rho_{s}, W_{s}\right)$ does not visit $\Theta_{x}^{c}$. We must check that $W_{s}\left(H_{s}-\right)$ exists whenever $\rho_{s}\left(\left\{H_{s}\right\}\right)>0$. Suppose thus that $\rho_{s}\left(\left\{H_{s}\right\}\right)>0$. Then, we have also $\rho_{s^{\prime}}\left(\left\{H_{s}\right\}\right)>0$, and so $H_{s^{\prime}} \geq H_{s}$, for all $s^{\prime}>s$ sufficiently close to $s$. In particular, we can find a rational $s_{1}>s$ such that $H_{s_{1}} \geq H_{s}$ and $\inf _{\left[s, s_{1}\right]} H_{r}=H_{s}$, which by the snake property implies that $W_{s_{1}}(t)=W_{s}(t)$ for every $t \in\left[0, H_{s}\right)$. However, from the construction of the measures $Q_{\mathrm{w}}^{h}$, it is clear that a.s. for every rational $r>0$, the killed path $W_{r}$ must have a left limit at every $t \in\left(0, H_{r}\right]$. We conclude that $W_{s}\left(H_{s}-\right)=W_{s_{1}}\left(H_{s}-\right)$ exists.

The càdlàg property of paths is obvious by construction. To obtain the Markov property, we consider nonnegative functions $f_{1}, \ldots, f_{n}$ on $M_{f}^{0}$ and $g_{1}, \ldots, g_{n}$ on $\mathcal{W}_{x}$. Then, if $0<s_{1}<\cdots<s_{n}$,

$$
\begin{aligned}
& \mathbb{E}_{\mu, \mathrm{w}}\left[f_{1}\left(\rho_{s_{1}}\right) g_{1}\left(W_{s_{1}}\right) \ldots f_{n}\left(\rho_{s_{n}}\right) g_{n}\left(W_{s_{n}}\right)\right] \\
& =\mathbf{E}_{\mu}\left[f_{1}\left(\rho_{s_{1}}\right) \ldots f_{n}\left(\rho_{s_{n}}\right) Q_{\mathrm{w}}^{H(\rho)}\left[g_{1}\left(W_{s_{1}}\right) \ldots g_{n}\left(W_{s_{n}}\right)\right]\right] \\
& = \\
& =\mathbf{E}_{\mu}\left[f_{1}\left(\rho_{s_{1}}\right) \ldots f_{n}\left(\rho_{s_{n}}\right) \int R_{m_{H(\rho)}\left(0, s_{1}\right), H\left(\rho_{s_{1}}\right)}\left(\mathrm{w}, d \mathrm{w}_{1}\right)\right. \\
& \left.\quad \ldots R_{m_{H(\rho)}\left(s_{n-1}, s_{n}\right), H\left(\rho_{s_{n}}\right)}\left(\mathrm{w}_{n-1}, d \mathrm{w}_{n}\right) g_{1}\left(\mathrm{w}_{1}\right) \ldots g_{n}\left(\mathrm{w}_{n}\right)\right] \\
& =\mathbb{E}_{\mu, \mathrm{w}}\left[f_{1}\left(\rho_{s_{1}}\right) g_{1}\left(W_{s_{1}}\right) \ldots f_{n-1}\left(\rho_{s_{n-1}}\right) g_{n-1}\left(W_{s_{n-1}}\right)\right. \\
& \left.\quad \mathbf{E}_{\rho_{s_{n-1}}}\left[f_{n}\left(\rho_{s_{n}-s_{n-1}}\right) \int R_{m_{H(\rho)}\left(0, s_{n}-s_{n-1}\right), H\left(\rho_{s_{n}-s_{n-1}}\right)}\left(W_{s_{n-1}}, d \mathrm{w}\right) g_{n}(\mathrm{w})\right]\right],
\end{aligned}
$$

where in the last equality we used the Markov property for $\left(\rho_{s}, s \geq 0\right)$ at time $s_{n-1}$. We get the desired result with a transition kernel given by

$$
\begin{aligned}
\mathbf{Q}_{r} G(\mu, \mathrm{w}) & =\int \mathbf{P}_{\mu}(d \rho) \int R_{m_{H(\rho)}(0, r), H\left(\rho_{r}\right)}\left(\mathrm{w}, d \mathrm{w}^{\prime}\right) G\left(\rho_{r}, \mathrm{w}^{\prime}\right) \\
& =\int \mathbf{P}_{\mu}(d \rho) \int Q_{\mathrm{w}}^{H(\rho)}(d W) G\left(\rho_{r}, W_{r}\right)
\end{aligned}
$$

In what follows we will often use the convenient notation $\bar{W}_{s}=\left(\rho_{s}, W_{s}\right)$. By our construction, the conditional distribution under $\mathbb{P}_{\mu, \mathrm{w}}$ of $\left(W_{s}, s \geq 0\right)$ knowing $\left(\rho_{s}, s \geq 0\right)$ is $Q_{\mathrm{w}}^{H(\rho)}$. In particular, if we write $\zeta_{s}=\zeta_{W_{s}}$ for the lifetime of $W_{s}$, we have

$$
\zeta_{s}=H\left(\rho_{s}\right)=H_{s} \text { for every } s \geq 0, \mathbb{P}_{\mu, \mathrm{w}} \text { a.s. }
$$

### 4.1.3 The strong Markov property

We denote by $\left(\mathcal{F}_{s}\right)_{s \geq 0}$ the canonical filtration on $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$.
Theorem 4.1.2 The process $\left(\bar{W}_{s}, s \geq 0 ; \mathbb{P}_{\mu, \mathrm{w}},(\mu, \mathrm{w}) \in \Theta_{x}\right)$ is strong Markov with respect to the filtration $\left(\mathcal{F}_{s+}\right)$.

Proof. Let $(\mu, \mathrm{w}) \in \Theta_{x}$. It is enough to prove that, if $T$ is a bounded stopping time of the filtration $\left(\mathcal{F}_{s+}\right)$, then, for any bounded $\mathcal{F}_{T+-}$-measurable functional $F$, for any bounded Lipschitz continuous function $f$ on $\Theta_{x}$, and for every $t>0$,

$$
\mathbb{E}_{\mu, \mathrm{w}}\left[F f\left(\bar{W}_{T+t}\right)\right]=\mathbb{E}_{\mu, \mathrm{w}}\left[F \mathbb{E}_{\bar{W}_{T}}\left[f\left(\bar{W}_{t}\right)\right]\right] .
$$

First observe that

$$
\begin{aligned}
\mathbb{E}_{\mu, \mathrm{w}}\left[F f\left(\bar{W}_{T+t}\right)\right] & =\lim _{n \rightarrow \infty} \sum_{k=1}^{\infty} \mathbb{E}_{\mu, \mathrm{w}}\left[F 1_{\left\{\frac{k-1}{n} \leq T<\frac{k}{n}\right\}} f\left(\bar{W}_{\frac{k}{n}+t}\right)\right] \\
& =\lim _{n \rightarrow \infty} \sum_{k=1}^{\infty} \mathbb{E}_{\mu, \mathrm{w}}\left[F 1_{\left\{\frac{k-1}{n} \leq T<\frac{k}{n}\right\}} \mathbf{Q}_{t} f\left(\bar{W}_{\frac{k}{n}}\right)\right] .
\end{aligned}
$$

In the first equality, we used the right continuity of paths, and in the second one the ordinary Markov property. We see that the desired result follows from the next lemma.

Lemma 4.1.3 Let $t>0$, let $T$ be a bounded stopping time of the filtration $\left(\mathcal{F}_{s+}\right)$ and let $f$ be a bounded Lipschitz continuous function on $\Theta_{x}$. Then the mapping $s \longrightarrow \mathbf{Q}_{t} f\left(\bar{W}_{s}\right)$ is $\mathbb{P}_{\mu, \mathrm{w}}$ a.s. right-continuous at $s=T$.

Proof of Lemma 4.1.3. We use the notation $Y_{s}=\left\langle\rho_{s}, 1\right\rangle$. Recall that $Y$ is distributed under $\mathbb{P}_{\mu, \mathrm{w}}$ as the reflected Lévy process $X-I$ started at $\langle\mu, 1\rangle$. Let $\varepsilon>0$. By the right-continuity of the paths of $Y$, if $s>T$ is sufficiently close to $T$, we have

$$
\varepsilon_{1}(s)=Y_{T}-\inf _{u \in[T, s]} Y_{u}<\varepsilon, \quad \varepsilon_{2}(s)=Y_{s}-\inf _{u \in[T, s]} Y_{u}<\varepsilon
$$

On the other hand, we know from (1.13) that $\rho_{T+s}=\left[k_{\varepsilon_{1}} \rho_{T}, \rho_{s}^{(T)}\right]$, and it follows that $k_{\varepsilon_{1}} \rho_{T}=k_{\varepsilon_{2}} \rho_{s}$. Furthermore, $\inf _{[T, s]} H\left(\rho_{u}\right)=H\left(k_{\varepsilon_{1}} \rho_{T}\right)=H\left(k_{\varepsilon_{2}} \rho_{s}\right)$, and by the snake property,

$$
W_{s}(u)=W_{T}(u), \quad \forall u \in\left[0, H\left(k_{\varepsilon_{1}} \rho_{T}\right)\right)
$$

Let us fix $\overline{\mathrm{w}}=(\mu, \mathrm{w}) \in \Theta_{x}$, and set

$$
\begin{gathered}
\mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})=\left\{\overline{\mathrm{w}}^{\prime}=\left(\mu^{\prime}, \mathrm{w}^{\prime}\right) \in \Theta_{x} ; \exists \varepsilon_{1}, \varepsilon_{2} \in[0, \varepsilon), k_{\varepsilon_{1}} \mu=k_{\varepsilon_{2}} \mu^{\prime},\right. \\
\text { and } \left.\mathrm{w}^{\prime}(u)=\mathrm{w}(u), \forall u \in\left[0, H\left(k_{\varepsilon_{1}} \mu\right)\right)\right\} .
\end{gathered}
$$

In view of the preceding observations, the proof of Lemma 4.1.3 reduces to checking that

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left(\sup _{\overline{\mathrm{w}}^{\prime} \in \mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})}\left|\mathbf{Q}_{t} f\left(\overline{\mathrm{w}}^{\prime}\right)-\mathbf{Q}_{t} f(\overline{\mathrm{w}})\right|\right)=0 \tag{4.1}
\end{equation*}
$$

We will use a coupling argument to obtain (4.1). More precisely, if $\overline{\mathrm{w}}^{\prime} \in \mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})$, we will introduce two (random) variables $\overline{\mathrm{w}}_{(1)}$ and $\overline{\mathrm{w}}_{(2)}$ such that $\overline{\mathrm{w}}_{(1)}$, resp. $\overline{\mathrm{w}}_{(2)}$, is distributed according to $\mathbf{Q}_{t}(\overline{\mathrm{w}}, \cdot)$, resp. $\mathbf{Q}_{t}\left(\overline{\mathrm{w}}^{\prime}, \cdot\right)$, and $\overline{\mathrm{w}}_{(1)}$ and $\overline{\mathrm{w}}_{(2)}$ are close to each other. Let us fix $\overline{\mathrm{w}}^{\prime} \in \mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})$ and let $\varepsilon_{1}, \varepsilon_{2} \in[0, \varepsilon)$ be associated with $\overline{\mathrm{w}}^{\prime}$ as in the definition of $\mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})$. For definiteness we assume that $\varepsilon_{1} \leq \varepsilon_{2}$ (the other case is treated in a symmetric way). Let $X^{(1)}$ be a copy of the Lévy process $X$ started at 0 and let $I^{(1)}$ and $\rho^{(1)}$ be the analogues of $I$ and $\rho$ for $X^{(1)}$. We can then define $\overline{\mathrm{w}}_{(1)}=\left(\mu_{(1)}, \mathrm{w}_{(1)}\right)$ by

$$
\begin{aligned}
& \mu_{(1)}=\left[k_{-I_{t}^{(1)}} \mu, \rho_{t}^{(1)}\right] \\
& \mathrm{w}_{(1)}(r)= \begin{cases}\mathrm{w}(r) & \text { if } r<H\left(k_{-I_{t}^{(1)}} \mu\right) \\
\xi^{(1)}\left(r-H\left(k_{-I_{t}^{(1)}} \mu\right)\right) & \text { if } H\left(k_{-I_{t}^{(1)}} \mu\right) \leq r<H\left(\mu_{(1)}\right),\end{cases}
\end{aligned}
$$

where, conditionally on $X^{(1)}, \xi^{(1)}=\left(\xi^{(1)}(t), t \geq 0\right)$ is a copy of the spatial motion $\xi$ started at $\mathrm{w}\left(H\left(k_{-I_{t}^{(1)}} \mu\right)-\right)$. Clearly, $\overline{\mathrm{w}}_{(1)}$ is distributed according to $\mathbf{Q}_{t}(\overline{\mathrm{w}}, \cdot)$.

The definition of $\overline{\mathrm{w}}_{(2)}$ is analogous but we use another copy of the underlying Lévy process. Precisely, we let $Z$ be a copy of $X$ independent of the pair $\left(X^{(1)}, \xi^{(1)}\right)$, and if $T_{*}(Z):=\inf \left\{r \geq 0: Z_{r}=\varepsilon_{1}-\varepsilon_{2}\right\}$, we set

$$
X_{s}^{(2)}= \begin{cases}Z_{s} & \text { if } 0 \leq s \leq T_{*}(Z) \\ \varepsilon_{1}-\varepsilon_{2}+X_{s-T_{*}(Z)}^{(1)} & \text { if } s>T_{*}(Z)\end{cases}
$$

We then take, with an obvious notation,

$$
\mu_{(2)}=\left[k_{-I_{t}^{(2)}} \mu^{\prime}, \rho_{t}^{(2)}\right] .
$$

The definition of $\mathrm{w}_{(2)}$ is somewhat more intricate. Let $\tau^{(1)}$ be the (a.s. unique) time of the minimum of $X^{(1)}$ over $[0, t]$. Consider the event

$$
A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)=\left\{T_{*}(Z)+\tau^{(1)}<t, I_{t}^{(1)}<-\varepsilon\right\}
$$

Notice that $T_{*}(Z)$ is small in probability when $\varepsilon$ is small, and $I_{t}^{(1)}<0$ a.s. It follows that $P\left[A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)\right] \geq 1-\alpha(\varepsilon)$, where the function $\alpha(\varepsilon)$ satisfies $\alpha(\varepsilon) \longrightarrow 0$ as $\varepsilon \rightarrow 0$. However, on the event $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$, we have $I_{t}^{(2)}=\varepsilon_{1}-\varepsilon_{2}+I_{t}^{(1)}$, and so

$$
k_{-I_{t}^{(2)}} \mu^{\prime}=k_{-I_{t}^{(1)}-\varepsilon_{1}} k_{\varepsilon_{2}} \mu^{\prime}=k_{-I_{t}^{(1)}-\varepsilon_{1}} k_{\varepsilon_{1}} \mu=k_{-I_{t}^{(1)}} \mu .
$$

Also recall that from the definition of $\mathcal{V}_{\varepsilon}(\overline{\mathrm{w}})$, we have $\mathrm{w}^{\prime}(r)=\mathrm{w}(r)$ for every $r<$ $H\left(k_{\varepsilon_{1}} \mu\right)$, hence for every $r<H\left(k_{-I_{t}^{(1)}} \mu\right)$ when $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$ holds.

We construct $\mathrm{w}_{(2)}$ by imposing that, on the set $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$,

$$
\mathrm{w}_{(2)}(r)= \begin{cases}\mathrm{w}^{\prime}(r)=\mathrm{w}(r) & \text { if } r<H\left(k_{-I_{t}^{(2)}} \mu^{\prime}\right)=H\left(k_{-I_{t}^{(1)}} \mu\right), \\ \xi^{(1)}\left(r-H\left(k_{-I_{t}^{(1)}} \mu\right)\right) & \text { if } H\left(k_{-I_{t}^{(1)}} \mu\right) \leq r<H\left(\mu_{(2)}\right),\end{cases}
$$

whereas on $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)^{c}$, we take

$$
\mathrm{w}_{(2)}(r)= \begin{cases}\mathrm{w}^{\prime}(r) & \text { if } r<H\left(k_{-I_{t}^{(2)}} \mu^{\prime}\right) \\ \xi^{(2)}\left(r-H\left(k_{-I_{t}^{(2)}} \mu^{\prime}\right)\right) & \text { if } H\left(k_{-I_{t}^{(2)}} \mu^{\prime}\right) \leq r<H\left(\mu_{(2)}\right)\end{cases}
$$

where, conditionally on $X^{(2)}, \xi^{(2)}$ is independent of $\xi^{(1)}$ and distributed according to the law of $\xi$ started at $\mathrm{w}^{\prime}\left(H\left(k_{-I_{t}^{(2)}} \mu^{\prime}\right)-\right)$. Note that, in the first case, we use the same process $\xi^{(1)}$ as in the definition of $\mathrm{w}_{(1)}$. It is again easy to verify that $\overline{\mathrm{w}}_{(2)}$ is distributed according to $\mathbf{Q}_{t}\left(\overline{\mathrm{w}}^{\prime}, \cdot\right)$.

To complete the proof, note that the distance in variation $d_{\mathrm{var}}\left(\mu_{(1)}, \mu_{(2)}\right)$ is equal to $d_{\mathrm{var}}\left(\rho_{t}^{(1)}, \rho_{t}^{(2)}\right)$ on the set $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$. Furthermore, from the construction of $X^{(2)}$, on the set $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$, we have also

$$
\rho_{t}^{(2)}=\rho_{t-T_{*}(Z)}^{(1)} .
$$

and thus $d_{\operatorname{var}}\left(\mu_{(1)}, \mu_{(2)}\right)=d_{\operatorname{var}}\left(\rho_{t}^{(1)}, \rho_{t-T_{*}(Z)}^{(1)}\right)$ is small in probability when $\varepsilon$ is small, because $t$ is a.s. not a discontinuity time of $\rho^{(1)}$. In addition, again on the set $A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)$, the paths $\mathrm{w}_{(1)}$ and $\mathrm{w}_{(2)}$ coincide on the interval $\left[0, H\left(\mu_{(1)}\right) \wedge H\left(\mu_{(2)}\right)\right)$, and so

$$
d\left(\mathrm{w}_{(1)}, \mathrm{w}_{(2)}\right) \leq\left|H\left(\mu_{(2)}\right)-H\left(\mu_{(1)}\right)\right|=\left|H\left(\rho_{t-T_{*}(Z)}^{(1)}\right)-H\left(\rho_{t}^{(1)}\right)\right|
$$

is small in probability when $\varepsilon$ goes to 0 . The limiting result (4.1) now follows from these observations and the fact that $P\left[A\left(\varepsilon, \varepsilon_{1}, \varepsilon_{2}\right)^{c}\right]$ tends to 0 as $\varepsilon$ goes to 0 .

### 4.1.4 Excursion measures

We know that $\mu=0$ is a regular recurrent point for the Markov process $\rho_{s}$, and the associated local time is the process $L_{s}^{0}$ of Section 1.3. It immediately follows that $(0, x)$ is also a regular recurrent point for the Lévy snake $\left(\rho_{s}, W_{s}\right)$, with associated local time $L_{s}^{0}$. We will denote by $\mathbb{N}_{x}$ the corresponding excursion measure. It is straightforward to verify that
(i) the law of ( $\rho_{s}, s \geq 0$ ) under $\mathbb{N}_{x}$ is the excursion measure $N(d \rho)$;
(ii) the conditional distribution of $\left(W_{s}, s \geq 0\right)$ under $\mathbb{N}_{x}$ knowing $\left(\rho_{s}, s \geq 0\right)$ is $Q_{x}^{H(\rho)}$.

From these properties and Proposition 1.2.5, we easily get for any nonnegative measurable function $F$ on $M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}$,

$$
\begin{equation*}
\mathbb{N}_{x}\left(\int_{0}^{\sigma} d s F\left(\rho_{s}, W_{s}\right)\right)=\int_{0}^{\infty} d a e^{-\alpha a} E^{0} \otimes \Pi_{x}\left[F\left(J_{a},\left(\xi_{r}, 0 \leq r \leq a\right)\right)\right] \tag{4.2}
\end{equation*}
$$

Here, as in Chapter $1, J_{a}(d r)$ stands for the measure $1_{[0, a]}(r) d U_{r}$, where $U$ is under the probability measure $P^{0}$ a subordinator with Laplace exponent $\tilde{\psi}(\lambda)-\alpha$, where $\tilde{\psi}(\lambda)=\psi(\lambda) / \lambda$. Note that the right side of (4.2) gives an invariant measure for the Lévy snake $\left(\rho_{s}, W_{s}\right)$.

The strong Markov property of the Lévy snake can be extended to the excursion measures in the following form. Let $T$ be a stopping time of the filtration $\left(\mathcal{F}_{s+}\right)$ such that $T>0, \mathbb{N}_{x}$ a.e., let $F$ be a nonnegative $\mathcal{F}_{T+\text {-measurable functional }}$ on $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$, and let $G$ be any nonnegative measurable functional on $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$. Then,

$$
\mathbb{N}_{x}\left[F G\left(\bar{W}_{T+s}, s \geq 0\right)\right]=\mathbb{N}_{x}\left[F \mathbb{E}_{\bar{W}_{T}}^{*}[G]\right]
$$

where $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ denotes the law under $\mathbb{P}_{\mu, \mathrm{w}}$ of the process $\left(\bar{W}_{s}, s \geq 0\right)$ stopped at $\inf \{s \geq$ $\left.0, \rho_{s}=0\right\}$. This statement follows from Theorem 4.1.2 by standard arguments.

### 4.2 The connection with superprocesses

### 4.2.1 Statement of the result

In this section, we state and prove the basic theorem relating the Lévy snake with the superprocess with spatial motion $\xi$ and branching mechanism $\psi$. This connection was already obtained in a less precise form in [33].

We start with a few simple observations. Let $\kappa(d s)$ be a random measure on $\mathbb{R}_{+}$, measurable with respect to the $\sigma$-field generated by $\left(\rho_{s}, s \geq 0\right)$. Then, from the form of the conditional distribution of ( $W_{s}, s \geq 0$ ) knowing ( $\rho_{s}, s \geq 0$ ), it is easy to see that, for any nonnegative measurable functional $F$ on $\mathcal{W}_{x}$,

$$
\mathbb{E}_{x}\left[\int \kappa(d s) F\left(W_{s}\right)\right]=\mathbb{E}_{x}\left[\int \kappa(d s) \Pi_{x}\left[F\left(\xi_{r}, 0 \leq r<H_{s}\right)\right]\right],
$$

and a similar formula holds under $\mathbb{N}_{x}$. This identity implies in particular that the left limit $W_{s}\left(H_{s}-\right)$ exists $\kappa(d s)$ a.e., $\mathbb{P}_{x}$ a.s. (or $\mathbb{N}_{x}$ a.e.). We will apply this simple observation to the random measure $d_{s} L_{s}^{a}$ associated with the local time of $H$ at level $a$ (cf Chapter 1). To simplify notation, we will write $\widehat{W}_{s}=W_{s}\left(H_{s}-\right)$ when the limit exists, and when the limit does not exist, we take $\widehat{W}_{s}=\Delta$, where $\Delta$ is a cemetery point added to $E$.

In order to state the main theorem of this section, we denote by $\mathcal{Z}_{a}=\mathcal{Z}_{a}(\rho, W)$ the random measure on $E$ defined by

$$
\left\langle\mathcal{Z}_{a}, f\right\rangle=\int_{0}^{\sigma} d_{s} L_{s}^{a} f\left(\widehat{W}_{s}\right) .
$$

This definition makes sense under the excursion measures $\mathbb{N}_{x}$.
Theorem 4.2.1 Let $\mu \in M_{f}(E)$ and let

$$
\sum_{i \in I} \delta_{\left(x_{i}, \rho^{i}, W^{i}\right)}
$$

be a Poisson point measure with intensity $\mu(d x) \mathbb{N}_{x}(d \rho d W)$. Set $Z_{0}=\mu$ and for every $a>0$

$$
Z_{a}=\sum_{i \in I} \mathcal{Z}_{a}\left(\rho^{i}, W^{i}\right) .
$$

The process $\left(Z_{a}, a \geq 0\right)$ is a superprocess with spatial motion $\xi$ and branching mechanism $\psi$, started at $\mu$.

This means that $\left(Z_{a}, a \geq 0\right)$ is a Markov process with values in $M_{f}(E)$, whose semigroup is characterized by the following Laplace functional. For every $0 \leq a \leq b$ and every function $f \in \mathcal{B}_{b+}(E)$,

$$
E\left[\exp -\left\langle Z_{b}, f\right\rangle \mid Z_{a}\right]=\exp -\left\langle Z_{a}, u_{b-a}\right\rangle
$$

where the function $\left(u_{t}(y), t \geq 0, y \in E\right)$ is the unique nonnegative solution of the integral equation

$$
\begin{equation*}
u_{t}(y)+\Pi_{y}\left(\int_{0}^{t} \psi\left(u_{t-r}\left(\xi_{r}\right)\right) d r\right)=\Pi_{y}\left(f\left(\xi_{t}\right)\right) . \tag{4.3}
\end{equation*}
$$

The proof of Theorem 4.2.1 is easily reduced to that of the following proposition.

Proposition 4.2.2 Let $0<a<b$ and let $f \in \mathcal{B}_{b+}(E)$. Then,

$$
\begin{equation*}
\mathbb{N}_{x}\left(\exp -\left\langle\mathcal{Z}_{b}, f\right\rangle \mid\left(\mathcal{Z}_{r}, 0 \leq r \leq a\right)\right)=\exp -\left\langle\mathcal{Z}_{a}, u_{b-a}\right\rangle \tag{4.4}
\end{equation*}
$$

where for every $t>0$ and $y \in E$,

$$
u_{t}(y)=\mathbb{N}_{y}\left(1-\exp -\left\langle\mathcal{Z}_{t}, f\right\rangle\right)
$$

Furthermore, if we set $u_{0}(y)=f(y)$, the function $\left(u_{t}(y), t \geq 0, y \in E\right)$ is the unique nonnegative solution of the integral equation (4.3).

Remark. Although $\mathbb{N}_{x}$ is an infinite measure, the conditioning in (4.4) makes sense because we can restrict our attention to the set $\left\{\mathcal{Z}_{a} \neq 0\right\}=\left\{L_{\sigma}^{a}>0\right\}$ which has finite $\mathbb{N}_{x}$-measure (cf Corollary 1.4.2). A similar remark applies in several places below, e.g. in the statement of Proposition 4.2.3.

Given Proposition 4.2.2, it is a straightforward exercise to verify that the process $\left(Z_{a}, a \geq 0\right)$ of Theorem 4.2 .1 has the finite-dimensional marginals of the superprocess with spatial motion $\xi$ and branching mechanism $\psi$, started at $\mu$. In fact the statement of Propostion 4.2 .2 means that the laws of $\left(\mathcal{Z}_{a}, a>0\right)$ under $\mathbb{N}_{y}, y \in E$ are the canonical measures of the superprocess with spatial motion $\xi$ and branching mechanism $\psi$, and given this fact, Theorem 4.2.1 is just the canonical representation of superprocesses.

The remaining part of this section is devoted to the proof of Proposition 4.2.2. We will proceed in two steps. In the first one, we introduce a $\sigma$-field $\mathcal{E}_{a}$ that contains $\sigma\left(\mathcal{Z}_{u}, 0 \leq u \leq a\right)$, and we compute $\mathbb{N}_{x}\left(\exp -\left\langle\mathcal{Z}_{b}, f\right\rangle \mid \mathcal{E}_{a}\right)$ in the form given by (4.4). In the second step, we establish the integral equation (4.3).

### 4.2.2 First step

Recall the notation of Section 1.3

$$
\tilde{\tau}_{s}^{a}=\inf \left\{r: \int_{0}^{r} d u 1_{\left\{H_{u} \leq a\right\}}>s\right\} .
$$

Note that $\tilde{\tau}_{s}^{a}<\infty$ for every $s \geq 0, \mathbb{N}_{x}$ a.e. For $a>0$, we let $\mathcal{E}_{a}$ be the $\sigma$-field generated by the right-continuous process $\left(\rho_{\tilde{\tau}_{s}^{a}}, W_{\tilde{\tau}_{s}^{a}} ; s \geq 0\right)$ and augmented with the class of all sets that are $\mathbb{N}_{x}$-negligible for every $x \in E$. From the second approximation of Proposition 1.3.3, it is easy to verify that $L_{\sigma}^{a}$ is measurable with respect to the $\sigma$ field generated by ( $\rho_{\tilde{\tau}_{s}^{a}}, s \geq 0$ ), and in particular with respect to $\mathcal{E}_{a}$ ( cf the beginning of the proof of Theorem 1.4.1).

We then claim that $\mathcal{Z}_{a}$ is $\mathcal{E}_{a}$-measurable. It is enough to check that, if $g$ is bounded and continuous on $\mathcal{W}_{x}$,

$$
\int_{0}^{\sigma} d L_{s}^{a} g\left(W_{s}\right)
$$

is $\mathcal{E}_{a}$-measurable. However, by Proposition 1.3.3, this integral is the limit in $\mathbb{N}_{x^{-}}$ measure as $\varepsilon \rightarrow 0$ of

$$
\frac{1}{\varepsilon} \int_{0}^{\sigma} d s 1_{\left\{a-\varepsilon<H_{s} \leq a\right\}} g\left(W_{s}\right) .
$$

For $\varepsilon<a$, this quantity coincides with

$$
\frac{1}{\varepsilon} \int_{0}^{\infty} d s 1\left\{a-\varepsilon<H_{\tilde{\tau}_{s}^{a}} \leq a\right\} g\left(W_{\tilde{\tau}_{s}^{a}}\right)
$$

and the claim follows from the definition of $\mathcal{E}_{a}$.
We then decompose the measure $\mathcal{Z}_{b}$ according to the contributions of the different excursions of the process $H$ above level $a$. Precisely, we let $\left(\alpha_{i}, \beta_{i}\right), i \in I$ be the excursion intervals of $H$ above $a$ over the time interval $[0, \sigma]$. We will use the following simple facts that hold $N$ a.e.: For every $i \in I$ and every $t>0$, we have

$$
\int_{0}^{\beta_{i}+t} 1_{\left\{H_{s} \leq a\right\}} d s>\int_{0}^{\beta_{i}} 1_{\left\{H_{s} \leq a\right\}} d s
$$

and

$$
L_{\beta_{i}+t}^{a}>L_{\beta_{i}}^{a} .
$$

The first assertion is an easy consequence of the strong Markov property of $\rho$, recalling that $\rho_{s}(\{a\})=0$ for every $s \geq 0, N$ a.e. To get the second one, we can use Proposition 1.3.1 and the definition of the local time $L^{a}$ to see that it is enough to prove that

$$
\int_{0}^{\beta_{i}+t} 1_{\left\{H_{s}>a\right\}} d s>\int_{0}^{\beta_{i}} 1_{\left\{H_{s}>a\right\}} d s
$$

for every $t>0$ and $i \in I$. Via a time-reversal argument (Corollary 3.1.6), it suffices to verify that, if $\sigma_{a}^{q}=\inf \left\{s>q: H_{s}>a\right\}$, we have

$$
\int_{0}^{\sigma_{a}^{q}+t} 1_{\left\{H_{s} \leq a\right\}} d s>\int_{0}^{\sigma_{a}^{q}} 1_{\left\{H_{s} \leq a\right\}} d s
$$

for every $t>0$ and every rational $q>0, N$ a.e. on the set $\left\{q<\sigma_{a}^{q}<\infty\right\}$. The latter fact is again a consequence of the strong Markov property of the process $\rho$.

As was observed in the proof of Proposition 1.3.1, for every $i \in I$, for every $s \in\left(\alpha_{i}, \beta_{i}\right)$, the restriction of $\rho_{s}$ to $[0, a]$ coincides with $\rho_{\alpha_{i}}=\rho_{\beta_{i}}$. Furthermore, the snake property implies that, for every $i \in I$, the paths $W_{s}, \alpha_{i}<s<\beta_{i}$ take the same value $x_{i}$ at time $a$, and this value must be the same as the left limit $\widehat{W}_{\alpha_{i}}=\widehat{W}_{\beta_{i}}$ (recall our assumption that $\xi$ has no fixed discontinuities). We can then define the pair $\left(\rho^{i}, W^{i}\right) \in \mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$ by setting

$$
\begin{array}{ll}
\left\langle\rho_{s}^{i}, g\right\rangle=\int_{(a, \infty)} \rho_{\alpha_{i}+s}(d r) g(r-a) & \text { if } 0<s<\beta_{i}-\alpha_{i} \\
\rho_{s}^{i}=0 & \text { if } s=0 \text { or } s \geq \beta_{i}-\alpha_{i}
\end{array}
$$

and

$$
\begin{array}{ll}
W_{s}^{i}(r)=W_{\alpha_{i}+s}(a+r), \zeta_{W_{s}^{i}}=H_{\alpha_{i}+s}-a & \text { if } 0<s<\beta_{i}-\alpha_{i} \\
W_{s}^{i}=x_{i} & \text { if } s=0 \text { or } s \geq \beta_{i}-\alpha_{i}
\end{array}
$$

Proposition 4.2.3 Under $\mathbb{N}_{x}$, conditionally on $\mathcal{E}_{a}$, the point measure

$$
\sum_{i \in I} \delta_{\left(\rho^{i}, W^{i}\right)}
$$

is a Poisson point measure with intensity

$$
\int \mathcal{Z}_{a}(d y) \mathbb{N}_{y}(\cdot)
$$

Proof. Let the process $\rho_{t}^{a}$ be defined as in Proposition 1.3.1. Note that under $N$ the definition of $\rho_{t}^{a}$ only makes sense for $t<\int_{0}^{\sigma} d s 1_{\left\{H_{s}>a\right\}}$. For convenience, we take $\rho_{t}^{a}=0$ if $t \geq \int_{0}^{\sigma} d s 1_{\left\{H_{s}>a\right\}}$. We also set

$$
\tilde{\rho}_{t}=\rho_{\tilde{\tau}_{t}^{a}}, \quad \tilde{W}_{t}=W_{\tilde{\tau}_{t}^{a}} .
$$

With these definitions, the processes $\rho^{i}, i \in I$ are exactly the excursions of the process $\rho^{a}$ away from 0 . For every $i \in I$, introduce the local time at the beginning (or the end) of excursion $\rho^{i}$ :

$$
\ell^{i}=L_{\alpha_{i}}^{a} .
$$

By Proposition 1.3.1 and standard excursion theory, we know that conditionally on the process $\tilde{\rho}_{t}$, the point measure

$$
\sum_{i \in I} \delta_{\left(\ell^{i}, \rho^{i}\right)}
$$

is Poisson with intensity $1_{\left[0, L_{\sigma}^{a}\right]}(\ell) d \ell N(d \rho)$ (recall that $L_{\sigma}^{a}$ is measurable with respect to the $\sigma$-field generated by $\tilde{\rho})$. Note that Proposition 1.3.1 is formulated under $\mathbb{P}_{x}$ : However, by considering the first excursion of $\rho$ away from 0 that hits the set $\left\{\sup H_{s}>a\right\}$, we can easily derive the previous assertion from Proposition 1.3.1.

Define $\tilde{L}_{s}^{a}=L_{\tilde{\tau}_{s}^{a}}^{a}$ (note that this is a continuous process), and let $\gamma^{a}(r)$ be the right-continuous inverse of $\tilde{L}^{a}$ :

$$
\gamma^{a}(r)=\inf \left\{s \geq 0: \tilde{L}_{s}^{a}>r\right\}
$$

Then, if $f$ is any nonnegative measurable function on $E$, we have $\mathbb{N}_{x}$ a.e.

$$
\begin{equation*}
\left\langle\mathcal{Z}_{a}, f\right\rangle=\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right)=\int_{0}^{\infty} d \tilde{L}_{s}^{a} f\left(\widehat{\tilde{W}}_{s}\right)=\int_{0}^{L_{\sigma}^{a}} d \ell f\left(\widehat{\tilde{W}}_{\gamma^{a}(\ell)}\right) \tag{4.5}
\end{equation*}
$$

Notice that both processes $\tilde{L}^{a}$ and $\gamma^{a}$ are measurable with respect to the $\sigma$-field generated by $\tilde{\rho}$ (for $\tilde{L}^{a}$, this follows again from Proposition 1.3.3).

Consider now the processes $\tilde{W}$ and $W^{i}, i \in I$. The following two properties are straightforward consequences of our construction:
(i) The law of $\tilde{W}$ under $Q_{x}^{H(\rho)}$ is $Q_{x}^{H(\tilde{\rho})}$.
(ii) Under $Q_{x}^{H(\rho)}$, conditionally on $\tilde{W}$, the "excursions" $W^{i}, i \in I$ are independent and the conditional distribution of $W^{i}$ is $Q_{x_{i}}^{H\left(\rho^{i}\right)}$, where $x_{i}=\widehat{W}_{\beta_{i}}=\widehat{W}_{\gamma^{a}\left(\ell^{i}\right)}$.

To verify the second expression for $x_{i}$, note that if $\tilde{A}_{s}^{a}=\int_{0}^{s} d r 1_{\left\{H_{r} \leq a\right\}}$, we have $W_{\beta_{i}}=\tilde{W}_{\tilde{A}_{\beta_{i}}^{a}}$ (because $\tilde{A}_{\beta_{i}+t}^{a}>\tilde{A}_{\beta_{i}}^{a}$ for every $t>0$ ) and $\tilde{A}_{\beta_{i}}^{a}=\gamma^{a}\left(\ell^{i}\right)$ (because $L_{\beta_{i}+t}^{a}>L_{\beta_{i}}^{a}=\ell_{i}$ for every $\left.t>0\right)$.

As a consequence of (i), the conditional distribution (under $\mathbb{N}_{x}$ ) of $\tilde{W}$ knowing $\rho$ depends only on $\tilde{\rho}$. Hence, $\tilde{W}$ and the point measure $\sum_{i \in I} \delta_{\left(\ell^{i}, \rho^{i}\right)}$ are conditionally independent given $\tilde{\rho}$ under $\mathbb{N}_{x}$.

Let $F$ be a nonnegative measurable function on $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$. We use the previous observations in the following calculation:

$$
\mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\sum_{i \in I} F\left(\rho^{i}, W^{i}\right)\right)\right)
$$

$$
\begin{aligned}
& =\int N(d \rho) Q_{x}^{H(\rho)}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\sum_{i \in I} F\left(\rho^{i}, W^{i}\right)\right)\right) \\
& =\int N(d \rho) Q_{x}^{H(\rho)}\left(G(\tilde{\rho}, \tilde{W}) \prod_{i \in I} Q_{\hat{W}_{\gamma^{a}\left(\ell^{i}\right)}^{H\left(\rho^{i}\right)}}^{H}\left(e^{-F\left(\rho^{i}, \cdot\right)}\right)\right) \\
& =\mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \prod_{i \in I} Q_{\tilde{\tilde{W}}_{\gamma^{a}\left(\ell^{i}\right)}^{H\left(\rho^{i}\right)}}\left(e^{-F\left(\rho^{i}, \cdot\right)}\right)\right) \\
& =\mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\int_{0}^{L_{\sigma}^{a}} d \ell \int N(d \rho) Q_{\hat{\tilde{W}}_{\gamma^{a}(\ell)}^{H(\rho)}}^{H}\left(1-e^{-F(\rho, \cdot)}\right)\right)\right) .
\end{aligned}
$$

The second equality follows from (ii) above. In the last one, we used the conditional independence of $\tilde{W}$ and of the point measure $\sum_{i \in I} \delta_{\left(\ell^{i}, \rho^{i}\right)}$, given $\tilde{\rho}$, and the fact that the conditional distribution of this point measure is Poisson with intensity $1_{\left[0, L_{\sigma}^{a}\right]}(\ell) d \ell N(d \rho)$. Using (4.5), we finally get

$$
\begin{aligned}
& \mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\sum_{i \in I} F\left(\rho^{i}, W^{i}\right)\right)\right) \\
& \quad=\mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\int_{0}^{L_{\sigma}^{a}} d \ell \mathbb{N}_{\hat{W}_{\gamma^{a}(\ell)}}\left(1-e^{-F(\rho, W)}\right)\right)\right) \\
& \quad=\mathbb{N}_{x}\left(G(\tilde{\rho}, \tilde{W}) \exp \left(-\int \mathcal{Z}_{a}(d y) \mathbb{N}_{y}\left(1-e^{-F}\right)\right)\right)
\end{aligned}
$$

This completes the proof.
Let $f$ be a nonnegative measurable function on $E$, and let $0 \leq a<b$. With the preceding notation, it is easy to verify that $\mathbb{N}_{x}$ a.s.

$$
\begin{aligned}
\left\langle\mathcal{Z}_{b}, f\right\rangle & =\int_{0}^{\sigma} d L_{s}^{b} f\left(\widehat{W}_{s}\right) \\
& =\sum_{i \in I} \int_{\alpha_{i}}^{\beta_{i}} d L_{s}^{b} f\left(\widehat{W}_{s}\right) \\
& =\sum_{i \in I} \int_{0}^{\infty} d L_{s}^{b-a}\left(\rho^{i}\right) f\left(\widehat{W}_{s}^{i}\right) \\
& =\sum_{i \in I}\left\langle\mathcal{Z}_{b-a}\left(\rho^{i}, W^{i}\right), f\right\rangle .
\end{aligned}
$$

As a consequence of Proposition 4.2.3, we have then

$$
\mathbb{N}_{x}\left[\exp -\left\langle\mathcal{Z}_{b}, f\right\rangle \mid \mathcal{E}_{a}\right]=\exp -\left\langle\mathcal{Z}_{a}, u_{b-a}\right\rangle
$$

where

$$
u_{r}(y)=\mathbb{N}_{y}\left[1-\exp -\left\langle\mathcal{Z}_{r}, f\right\rangle\right] .
$$

### 4.2.3 Second step

It remains to prove that the function $\left(u_{r}(y), r \geq 0, y \in E\right)$ introduced at the end of the first step solves the integral equation (4.3). By definition, we have for $a>0$,

$$
u_{a}(y)=\mathbb{N}_{y}\left(1-\exp -\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right)\right)
$$

$$
\begin{align*}
& =\mathbb{N}_{y}\left(\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right) \exp \left(-\int_{s}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)\right)\right) \\
& =\mathbb{N}_{y}\left(\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right) \mathbb{E}_{\rho_{s}, W_{s}}^{*}\left[\exp -\int_{0}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)\right]\right) \tag{4.6}
\end{align*}
$$

where we recall that $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ stands for the law of the Lévy snake started at $(\mu, \mathrm{w})$ and stopped when $\rho_{s}$ first hits 0 . In the last equality, we replaced $\exp -\int_{s}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)$ by its optional projection, using the strong Markov property of the Lévy snake to identify this projection.

We now need to compute for a fixed $(\mu, \mathrm{w}) \in \Theta_{x}$,

$$
\mathbb{E}_{\mu, \mathrm{w}}^{*}\left[\exp -\int_{0}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)\right]
$$

We will derive this calculation from a more general fact, that is also useful for forthcoming applications. First recall that $Y_{t}=\left\langle\rho_{t}, 1\right\rangle$ is distributed under $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ as the underlying Lévy process started at $\langle\mu, 1\rangle$ and stopped when it first hits 0 . We write $K_{t}=\inf _{r \leq t} Y_{r}$, and we denote by $\left(\alpha_{i}, \beta_{i}\right), i \in I$ the excursion intervals of $Y_{t}-K_{t}$ away from 0 . For every $i \in I$, we set $h_{i}=H_{\alpha_{i}}=H_{\beta_{i}}$. From the snake property, it is easy to verify that $W_{s}\left(h_{i}\right)=\mathrm{w}\left(h_{i}-\right)$ for every $s \in\left(\alpha_{i}, \beta_{i}\right), i \in I, \mathbb{P}_{\mu, \mathrm{w}}^{*}$ a.s. We then define the pair $\left(\rho^{i}, W^{i}\right)$ by the formulas

$$
\begin{array}{ll}
\left\langle\rho_{s}^{i}, g\right\rangle=\int_{\left(h_{i}, \infty\right)} \rho_{\alpha_{i}+s}(d r) g\left(r-h_{i}\right) & \text { if } 0 \leq s \leq \beta_{i}-\alpha_{i} \\
\rho_{s}^{i}=0 & \text { if } s>\beta_{i}-\alpha_{i}
\end{array}
$$

and

$$
\begin{array}{ll}
W_{s}^{i}(t)=W_{\alpha_{i}+s}\left(h_{i}+t\right), \quad \zeta_{s}^{i}=H_{\alpha_{i}+s}-h_{i} & \text { if } 0<s<\beta_{i}-\alpha_{i} \\
W_{s}^{i}=\mathrm{w}\left(h_{i}-\right) & \text { if } s=0 \text { or } s \geq \beta_{i}-\alpha_{i}
\end{array}
$$

Lemma 4.2.4 Let $(\mu, \mathrm{w}) \in \Theta_{x}$. The point measure

$$
\sum_{i \in I} \delta_{\left(h_{i}, \rho^{i}, W^{i}\right)}
$$

is under $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ a Poisson point measure with intensity

$$
\mu(d h) \mathbb{N}_{\mathrm{w}(h-)}(d \rho d W)
$$

Proof. Consider first the point measure

$$
\sum_{i \in I} \delta_{\left(h_{i}, \rho^{i}\right)}
$$

If $I_{s}=K_{s}-\langle\mu, 1\rangle$, we have $h_{i}=H\left(\rho_{\alpha_{i}}\right)=H\left(k_{-I_{\alpha_{i}}} \mu\right)$. Excursion theory for $Y_{s}-K_{s}$ ensures that

$$
\sum_{i \in I} \delta_{\left(-I_{\alpha_{i}}, \rho^{i}\right)}
$$

is under $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ a Poisson point measure with intensity $1_{[0,<\mu, 1>]}(u) d u N(d \rho)$. Since the image measure of $1_{[0,<\mu, 1>]}(u) d u$ under the mapping $u \longrightarrow H\left(k_{u} \mu\right)$ is precisely the measure $\mu$, it follows that

$$
\sum_{i \in I} \delta_{\left(h_{i}, \rho^{i}\right)}
$$

is a Poisson point measure with intensity $\mu(d h) N(d \rho)$. To complete the proof, it remains to obtain the conditional distribution of ( $W^{i}, i \in I$ ) knowing ( $\rho_{s}, s \geq 0$ ). However, the form of the conditional law $Q_{\mathrm{w}}^{H}$ easily implies that under $Q_{\mathrm{w}}^{H}$, the processes $W^{i}, i \in I$ are independent, and furthermore the conditional distribution of $W^{i}$ is $Q_{\mathrm{w}\left(h_{i}-\right)}^{H^{i}}$, where $H_{s}^{i}=H\left(\rho_{s}^{i}\right)$. It follows that

$$
\sum_{i \in I} \delta_{\left(h_{i}, \rho^{i}, W^{i}\right)}
$$

is a Poisson measure with intensity

$$
\mu(d h) N(d \rho) Q_{\mathrm{w}(h-)}^{H(\rho)}(d W)=\mu(d h) \mathbb{N}_{\mathrm{w}(h-)}(d \rho d W)
$$

This completes the proof.
We apply Lemma 4.2.4 to a pair $(\mu, \mathrm{w})$ such that $H(\mu) \leq a$ and $\mu(\{H(\mu)\})=0$. Then, it is easy to verify that $\mathbb{P}_{\mu, \mathrm{w}}^{*}$ a.s.

$$
\int_{0}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)=\sum_{i \in I} \int_{\alpha_{i}}^{\beta_{i}} d L_{r}^{a} f\left(\widehat{W}_{r}\right)=\sum_{i \in I} \int_{0}^{\infty} d L_{r}^{a-h_{i}}\left(\rho^{i}\right) f\left(\widehat{W}_{r}^{i}\right)
$$

and thus, by Lemma 4.2.4,

$$
\begin{equation*}
\mathbb{E}_{\mu, \mathrm{w}}^{*}\left[\exp -\int_{0}^{\infty} d L_{r}^{a} f\left(\widehat{W}_{r}\right)\right]=\exp \left(-\int \mu(d h) \mathbb{N}_{\mathrm{w}(h-)}\left[1-\exp -\left\langle\mathcal{Z}_{a-h}, f\right\rangle\right]\right) \tag{4.7}
\end{equation*}
$$

We now come back to formula (4.6). As a consequence of Proposition 1.2.5, we know that $\rho_{s}(\{a\})=0$ for every $s \geq 0, \mathbb{N}_{y}$ a.e. We can thus use (4.7) to get

$$
\begin{align*}
u_{a}(y) & =\mathbb{N}_{y}\left(\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right) \exp \left(-\int \rho_{s}(d h) \mathbb{N}_{W_{s}(h-)}\left[1-\exp -\left\langle\mathcal{Z}_{a-h}, f\right\rangle\right]\right)\right) \\
& =\mathbb{N}_{y}\left(\int_{0}^{\infty} d L_{s}^{a} f\left(\widehat{W}_{s}\right) \exp \left(-\int \rho_{s}(d h) u_{a-h}\left(W_{s}(h-)\right)\right)\right) \tag{4.8}
\end{align*}
$$

Let $J_{a}, P^{0}$ be as in (4.2).
Lemma 4.2.5 For any nonnegative measurable function $F$ on $\Theta_{y}$,

$$
\mathbb{N}_{y}\left(\int_{0}^{\infty} d L_{s}^{a} F\left(\rho_{s}, W_{s}\right)\right)=e^{-\alpha a} E^{0} \otimes \Pi_{y}\left[F\left(J_{a},\left(\xi_{r}, 0 \leq r<a\right)\right)\right]
$$

Proof. If $F\left(\rho_{s}, W_{s}\right)$ depends only on $\rho_{s}$, the result follows from Corollary 1.3.4. In the general case, we may take $F$ such that $F\left(\rho_{s}, W_{s}\right)=F_{1}\left(\rho_{s}\right) F_{2}\left(W_{s}\right)$, and we use the simple observation of the beginning of this section.

From (4.8) and Lemma 4.2.5, we get

$$
\begin{aligned}
u_{a}(y) & =e^{-\alpha a} E^{0} \otimes \Pi_{y}\left[f\left(\xi_{a}\right) \exp \left(-\int J_{a}(d h) u_{a-h}\left(\xi_{h-}\right)\right)\right] \\
& =\Pi_{y}\left[f\left(\xi_{a}\right) \exp \left(-\int_{0}^{a} \tilde{\psi}\left(u_{a-r}\left(\xi_{r}\right)\right) d r\right)\right]
\end{aligned}
$$

The proof of (4.3) is now completed by routine calculations. We have

$$
u_{a}(y)=\Pi_{y}\left[f\left(\xi_{a}\right)\right]-\Pi_{y}\left[f\left(\xi_{a}\right) \int_{0}^{a} d b \tilde{\psi}\left(u_{a-b}\left(\xi_{b}\right)\right) \exp \left(-\int_{b}^{a} \tilde{\psi}\left(u_{a-r}\left(\xi_{r}\right)\right) d r\right)\right]
$$

$$
\begin{aligned}
& =\Pi_{y}\left[f\left(\xi_{a}\right)\right]-\Pi_{y}\left[\int_{0}^{a} d b \tilde{\psi}\left(u_{a-b}\left(\xi_{b}\right)\right)\right. \\
& \left.\quad \Pi_{\xi_{b}}\left[f\left(\xi_{a-b}\right) \exp \left(-\int_{0}^{a-b} \tilde{\psi}\left(u_{a-b-r}\left(\xi_{r}\right)\right) d r\right)\right]\right] \\
& =\Pi_{y}\left[f\left(\xi_{a}\right)\right]-\Pi_{y}\left[\int_{0}^{a} d b \tilde{\psi}\left(u_{a-b}\left(\xi_{b}\right)\right) u_{a-b}\left(\xi_{b}\right)\right]
\end{aligned}
$$

which gives (4.3) and completes the proof of Proposition 4.2.2.
Remark. By combining the previous arguments, especially Lemma 4.2.4, with the duality properties of $\rho$ (see the derivation of (4.33) below), we could have obtained the following result. Let $f$ and $g$ be nonnegative measurable functions defined respectively on $E$ and on $M_{f}(E)$. Let $V$ be a subordinator with Laplace exponent $\psi^{\prime}(\lambda)-\alpha$ defined under the probability measure $P^{0}$. Then, for every $a>0$,

$$
\begin{align*}
& \mathbb{N}_{x}\left(\int \mathcal{Z}_{a}(d y) f(y) g\left(\mathcal{Z}_{a}\right)\right) \\
& \quad=e^{-\alpha a} E^{0} \otimes \Pi_{x}\left[f\left(\xi_{a}\right) E_{V, \xi}^{a}\left[g\left(\int \mathcal{N}(d s d \rho d W) \mathcal{Z}_{a-s}(\rho, W)\right)\right]\right] \tag{4.9}
\end{align*}
$$

where $\mathcal{N}(d s d \rho d W)$ is under $P_{V, \xi}^{a}$ a Poisson point measure on $\mathbb{R}_{+} \times \mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$ with intensity

$$
1_{[0, a]}(s) d V_{s} \mathbb{N}_{\xi_{s}}(d \rho d W) .
$$

Formula (4.9) identifies the Palm distributions associated with the random measure $\mathcal{Z}_{a}$ under $\mathbb{N}_{x}$. It should be compared with the results of Chapter 4 of [8], in particular Proposition 4.1.5. In the stable case considered in [8], $V$ is a stable subordinator, which arises analytically in the derivation of the Palm measure formula. Here $V$ can be interpreted probabilistically in terms of the measure-valued processes $\rho$ and $\eta$. As (4.9) will not be needed in the applications below, we will leave details of the proof to the reader.

### 4.3 Exit measures

Throughout this section, we consider an open set $D \subset E$, and we denote by $\tau$ the first exit time of $\xi$ from $D$ :

$$
\tau=\inf \left\{t \geq 0: \xi_{t} \notin D\right\}
$$

where $\inf \emptyset=\infty$ as usual. By abuse of notation, we will also denote by $\tau(\mathrm{w})$ the exit time from $D$ of a killed path $\mathrm{w} \in \mathcal{W}$,

$$
\tau(\mathrm{w})=\inf \left\{t \in\left[0, \zeta_{\mathrm{w}}\right): \mathrm{w}(t) \notin D\right\}
$$

Let $x \in D$. The next result is much analogous to Proposition 1.3.1.
Proposition 4.3.1 Assume that $\Pi_{x}(\tau<\infty)>0$. Then,

$$
\int_{0}^{\infty} d s 1_{\left\{\tau\left(W_{s}\right)<H_{s}\right\}}=\infty, \quad \mathbb{P}_{x} \text { a.s. }
$$

Furthermore, let

$$
\sigma_{s}^{D}=\inf \left\{t \geq 0: \int_{0}^{t} d r 1_{\left\{\tau\left(W_{r}\right)<H_{r}\right\}}>s\right\}
$$

and let $\rho_{s}^{D} \in M_{f}\left(\mathbb{R}_{+}\right)$be defined by

$$
\left\langle\rho_{s}^{D}, f\right\rangle=\int \rho_{\sigma_{s}^{D}}(d r) f\left(r-\tau\left(W_{\sigma_{s}^{D}}\right)\right) 1_{\left\{r>\tau\left(W_{\sigma_{s}^{D}}\right)\right\}}
$$

Then the process $\left(\rho_{s}^{D}, s \geq 0\right)$ has the same distribution under $\mathbb{P}_{x}$ as $\left(\rho_{s}, s \geq 0\right)$.
Remark. We could have considered the more general situation of a space-time open set $D$ (as a matter of fact, this is not really more general as we could replace $\xi_{t}$ by $\left.\left(t, \xi_{t}\right)\right)$. Taking $D=[0, a) \times E$, we would recover part of the statement of Proposition 1.3.1. This proposition contains an independence statement that could also be extended to the present setting.

Proof. To simplify notation, we set

$$
A_{s}^{D}=\int_{0}^{s} d r 1_{\left\{\tau\left(W_{r}\right)<H_{r}\right\}}
$$

By using (4.2), excursion theory and our assumption $\Pi_{x}(\tau<\infty)>0$, it is a simple exercise to verify that $A_{\infty}^{D}=\infty, \mathbb{P}_{x}$ a.s., and thus the definition of $\sigma_{s}^{D}$ makes sense for every $s \geq 0$, a.s. The arguments then are much similar to the proof of Proposition 1.3.1. For every $\varepsilon>0$, we introduce the stopping times $S_{\varepsilon}^{k}, T_{\varepsilon}^{k}, k \geq 1$, defined inductively by:

$$
\begin{aligned}
& S_{\varepsilon}^{1}=\inf \left\{s \geq 0: \tau\left(W_{s}\right)<\infty \text { and } \rho_{s}\left(\left(\tau\left(W_{s}\right), \infty\right)\right) \geq \varepsilon\right\} \\
& T_{\varepsilon}^{k}=\inf \left\{s \geq S_{\varepsilon}^{k}: \tau\left(W_{s}\right)=\infty\right\} \\
& S_{\varepsilon}^{k+1}=\inf \left\{s \geq T_{\varepsilon}^{k}: \tau\left(W_{s}\right)<\infty \text { and } \rho_{s}\left(\left(\tau\left(W_{s}\right), \infty\right)\right) \geq \varepsilon\right\}
\end{aligned}
$$

It is easy to see that these stopping times are a.s. finite, and $S_{\varepsilon}^{k} \uparrow \infty, T_{\varepsilon}^{k} \uparrow \infty$ as $k \uparrow \infty$.

From the key formula (1.13), we see that for

$$
S_{\varepsilon}^{k} \leq s<\inf \left\{r \geq S_{\varepsilon}^{k}:\left\langle\rho_{r}, 1\right\rangle \leq \rho_{S_{\varepsilon}^{k}}\left(\left[0, \tau\left(W_{S_{\varepsilon}^{k}}\right)\right]\right)\right\}
$$

we have $H_{s}>\tau\left(W_{S_{\varepsilon}^{k}}\right)$, and the paths $W_{s}$ and $W_{S_{\varepsilon}^{k}}$ coincide over $\left[0, \tau\left(W_{S_{\varepsilon}^{k}}\right)\right]$ (by the snake property), so that in particular $\tau\left(W_{s}\right)=\tau\left(W_{S_{\varepsilon}^{k}}\right)<\infty$. On the other hand, for

$$
s=\inf \left\{r \geq S_{\varepsilon}^{k}:\left\langle\rho_{r}, 1\right\rangle \leq \rho_{S_{\varepsilon}^{k}}\left(\left[0, \tau\left(W_{S_{\varepsilon}^{k}}\right)\right]\right)\right\}
$$

the path $W_{s}$ is the restriction of $W_{S_{\varepsilon}^{k}}$ to $\left[0, \tau\left(W_{S_{\varepsilon}^{k}}\right)\right)$ and thus $\tau\left(W_{s}\right)=\infty$. From these observations, we see that

$$
T_{\varepsilon}^{k}=\inf \left\{r \geq S_{\varepsilon}^{k}:\left\langle\rho_{r}, 1\right\rangle \leq \rho_{S_{\varepsilon}^{k}}\left(\left[0, \tau\left(W_{S_{\varepsilon}^{k}}\right)\right]\right)\right\}
$$

and that conditionally on the past up to time $S_{\varepsilon}^{k}$, the process

$$
Y_{s}^{k, \varepsilon}=\rho_{\left(S_{\varepsilon}^{k}+s\right) \wedge T_{\varepsilon}^{k}}\left(\left(\tau\left(W_{S_{\varepsilon}^{k}}\right), \infty\right)\right)
$$

is distributed as the underlying Lévy process started at $\rho_{S_{\varepsilon}^{k}}\left(\left(\tau\left(W_{S_{\varepsilon}^{k}}\right), \infty\right)\right)$ and stopped at its first hitting time of 0 .

The same argument as in the proof of (1.18) shows that, for every $t \geq 0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \sup _{\left\{k \geq 1, S_{\varepsilon}^{k} \leq t\right\}} \rho_{S_{\varepsilon}^{k}}\left(\left(\tau\left(W_{S_{\varepsilon}^{k}}\right), \infty\right)\right)=0, \quad \text { a.s. } \tag{4.10}
\end{equation*}
$$

The remaining part of the proof is very similar to the end of the proof of Proposition 1.3.1. Using (4.10) and the observations preceding (4.10), we get by a passage to the limit $\varepsilon \rightarrow 0$ that the total mass process $\left\langle\rho_{s}^{D}, 1\right\rangle=\rho_{\sigma_{s}^{D}}\left(\left(\tau\left(W_{\sigma_{s}^{D}}\right), \infty\right)\right)$ has the same distribution as the process $\left\langle\rho_{s}, 1\right\rangle$. Then the statement of the proposition follows by an argument similar to the second step of the proof of Proposition 1.3.1.

Let $\ell^{D}=\left(\ell^{D}(s), s \geq 0\right)$ be the local time at 0 of the process $\left\langle\rho^{D}, 1\right\rangle$. Recall the notation $A_{s}^{D}$ from the previous proof. We define the exit local time from $D$ by the formula

$$
L_{s}^{D}=\ell^{D}\left(A_{s}^{D}\right)=\ell^{D}\left(\int_{0}^{s} d r 1_{\left\{\tau\left(W_{r}\right)<H_{r}\right\}}\right) .
$$

Recall from (4.2) the notation $J_{a}, P^{0}$.
Proposition 4.3.2 For any nonnegative measurable function $\Phi$ on $M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}$,

$$
\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} \Phi\left(\rho_{s}, W_{s}\right)\right)=E^{0} \otimes \Pi_{x}\left[1_{\{\tau<\infty\}} e^{-\alpha \tau} \Phi\left(J_{\tau},\left(\xi_{r}, 0 \leq r<\tau\right)\right)\right]
$$

Proof. By applying Lemma 1.3.2 to the reflected Lévy process $\left\langle\rho^{D}, 1\right\rangle$, we get for every $s \geq 0$,

$$
\ell^{D}(s)=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \int_{0}^{s} d r 1_{\left\{0<H\left(\rho_{r}^{D}\right) \leq \varepsilon\right\}}
$$

in $L^{1}\left(\mathbb{P}_{x}\right)$. From a simple monotonicity argument, we have then for every $t \geq 0$

$$
\lim _{\varepsilon \rightarrow 0} \mathbb{E}_{x}\left[\sup _{s \leq t}\left|\ell^{D}(s)-\frac{1}{\varepsilon} \int_{0}^{s} d r 1_{\left\{0<H\left(\rho_{r}^{D}\right) \leq \varepsilon\right\}}\right|\right]=0
$$

Using the formulas $L_{s}^{D}=\ell^{D}\left(A_{s}^{D}\right)$ and $H\left(\rho_{\sigma_{r}^{D}}\right)=\tau\left(W_{\sigma_{r}^{D}}\right)+H\left(\rho_{r}^{D}\right)$ (the latter holding on the set $\left\{H\left(\rho_{\sigma_{r}^{D}}\right)>\tau\left(W_{\sigma_{r}^{D}}\right)\right\}$, by the definition of $\left.\rho^{D}\right)$, we obtain

$$
\lim _{\varepsilon \rightarrow 0} \mathbb{E}_{x}\left[\sup _{s \leq t}\left|L_{s}^{D}-\frac{1}{\varepsilon} \int_{0}^{s} d r 1_{\left\{\tau\left(W_{r}\right)<H_{r} \leq \tau\left(W_{r}\right)+\varepsilon\right\}}\right|\right]=0 .
$$

Arguing as in the derivation of (1.29), we get, for any measurable subset $V$ of $\mathbb{D}\left(\mathbb{R}_{+}, M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}\right)$ such that $\mathbb{N}_{x}(V)<\infty$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \mathbb{N}_{x}\left(1_{V} \sup _{s \leq t}\left|L_{s}^{D}-\frac{1}{\varepsilon} \int_{0}^{s} d r 1_{\left\{\tau\left(W_{r}\right)<H_{r} \leq \tau\left(W_{r}\right)+\varepsilon\right\}}\right|\right)=0 \tag{4.11}
\end{equation*}
$$

We then observe that for any bounded measurable function $F$ on $\mathbb{R}_{+} \times M_{f}\left(\mathbb{R}_{+}\right) \times$ $\mathcal{W}$, we have

$$
\begin{equation*}
\mathbb{N}_{x}\left(\int_{0}^{\sigma} d s F\left(s, \rho_{s}, W_{s}\right)\right)=E \otimes \Pi_{x}\left[\int_{0}^{L_{\infty}} d a F\left(L^{-1}(a), \Sigma_{a},\left(\xi_{r}, 0 \leq r<a\right)\right)\right] \tag{4.12}
\end{equation*}
$$

where the random measure $\Sigma_{a}$ is defined under $P$ by

$$
\left\langle\Sigma_{a}, g\right\rangle=\int_{0}^{L^{-1}(a)} d S_{s} g\left(a-L_{s}\right)
$$

Indeed, we observe that the special case where $F(s, \mu, \mathrm{w})$ does not depend on w ,

$$
N\left(\int_{0}^{\sigma} d s F\left(s, \rho_{s}\right)\right)=E\left[\int_{0}^{L_{\infty}} d a F\left(L^{-1}(a), \Sigma_{a}\right)\right]
$$

is a consequence of Proposition 1.1.4 (see the proof of Proposition 1.2.5), and it then suffices to use the conditional distribution of $W$ knowing ( $\rho_{s}, s \geq 0$ ).

After these preliminaries, we turn to the proof of the proposition. We let $F$ be a bounded continuous function on $\mathbb{R}_{+} \times M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}$, and assume in addition that there exist $\delta>0$ and $A>0$ such that $F(s, \mu, \mathrm{w})=0$ if $s \leq \delta$ or $s \geq A$. As a consequence of (4.11) and (4.12), we have then

$$
\begin{aligned}
& \mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} F\left(s, \rho_{s}, W_{s}\right)\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} \mathbb{N}_{x}\left(\frac{1}{\varepsilon} \int_{0}^{\sigma} d r F\left(r, \rho_{r}, W_{r}\right) 1_{\left\{\tau\left(W_{r}\right)<H_{r} \leq \tau\left(W_{r}\right)+\varepsilon\right\}}\right) \\
& \quad=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} E \otimes \Pi_{x}\left[\int_{0}^{L_{\infty}} d a F\left(L^{-1}(a), \Sigma_{a},\left(\xi_{r}, 0 \leq r<a\right)\right) 1_{\{\tau<a \leq \tau+\varepsilon\}}\right] \\
& \quad=E \otimes \Pi_{x}\left[1_{\left\{\tau<L_{\infty}\right\}} F\left(L^{-1}(\tau), \Sigma_{\tau},\left(\xi_{r}, 0 \leq r<\tau\right)\right)\right] .
\end{aligned}
$$

From this identity, we easily get

$$
\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} \Phi\left(\rho_{s}, W_{s}\right)\right)=E \otimes \Pi_{x}\left[1_{\left\{\tau<L_{\infty}\right\}} \Phi\left(\Sigma_{\tau},\left(\xi_{r}, 0 \leq r<\tau\right)\right)\right]
$$

Recall that $P\left[L_{\infty}>a\right]=e^{-\alpha a}$ and, that conditionally on $\left\{L_{\infty}>a\right\}, \Sigma_{a}$ has the same distribution as $J_{a}$. The last formula is thus equivalent to the statement of the proposition.

We now introduce an additional assumption. Namely we assume that for every $x \in D$, the process $\xi$ is $\Pi_{x}$ a.s. continuous at $t=\tau$, on the event $\{\tau<\infty\}$. Obviously this assumption holds if $\xi$ has continuous sample paths, but there are other cases of interest.

Under this assumption, Proposition 4.3.2 ensures that $\mathbb{N}_{x}$ a.e. the left limit $\widehat{W}_{s}$ exists $d L_{s}^{D}$ a.e. over $[0, \sigma]$ and belongs to $\partial D$. We define under $\mathbb{N}_{x}$ the exit measure $\mathcal{Z}^{D}$ from $D$ by the formula

$$
\left\langle\mathcal{Z}^{D}, g\right\rangle=\int_{0}^{\sigma} d L_{s}^{D} g\left(\widehat{W}_{s}\right)
$$

The previous considerations show that $\mathcal{Z}^{D}$ is a (finite) measure supported on $\partial D$. As a consequence of Proposition 4.3.2, we have for every nonnegative measurable function $g$ on $\partial D$,

$$
\mathbb{N}_{x}\left(\left\langle\mathcal{Z}^{D}, g\right\rangle\right)=\Pi_{x}\left(1_{\{\tau<\infty\}} e^{-\alpha \tau} g\left(\xi_{\tau}\right)\right)
$$

Theorem 4.3.3 Let $g$ be a bounded nonnegative measurable function on $\partial D$. For every $x \in D$ set

$$
u(x)=\mathbb{N}_{x}\left(1-\exp -\left\langle\mathcal{Z}^{D}, g\right\rangle\right)
$$

Then $u$ solves the integral equation

$$
u(x)+\Pi_{x}\left(\int_{0}^{\tau} d t \psi\left(u\left(\xi_{t}\right)\right)\right)=\Pi_{x}\left(1_{\{\tau<\infty\}} g\left(\xi_{\tau}\right)\right)
$$

Proof. Several arguments are analogous to the second step of the proof of Proposition 4.2.2 in Section 4, and so we will skip some details. By the definition of $\mathcal{Z}^{D}$, we have

$$
u(x)=\mathbb{N}_{x}\left(1-\exp -\int_{0}^{\sigma} d L_{s}^{D} g\left(\widehat{W}_{s}\right)\right)
$$

$$
\begin{aligned}
& =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} g\left(\widehat{W}_{s}\right) \exp \left(-\int_{s}^{\sigma} d L_{r}^{D} g\left(\widehat{W}_{r}\right)\right)\right) \\
& =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} g\left(\widehat{W}_{s}\right) \mathbb{E}_{\rho_{s}, W_{s}}^{*}\left[\exp -\int_{0}^{\infty} d L_{r}^{D} g\left(\widehat{W}_{r}\right)\right]\right)
\end{aligned}
$$

Note that the definition of the random measure $d L_{r}^{D}$ makes sense under $\mathbb{P}_{\mu, \mathrm{w}}^{*}$, provided that $\tau(\mathrm{w})=\infty$, thanks to Lemma 4.2.4 and the approximations used in the proof of Proposition 4.3.2. Using Lemma 4.2.4 as in subsection 4.2.3, we get if $(\mu, \mathrm{w}) \in \Theta_{x}$ is such that $\tau(\mathrm{w})=\infty$,

$$
\begin{aligned}
& \mathbb{E}_{\mu, \mathrm{w}}^{*}\left[\exp -\int_{0}^{\infty} d L_{r}^{D} g\left(\widehat{W}_{r}\right)\right] \\
& \quad=\exp \left(-\int \mu(d h) \mathbb{N}_{\mathrm{w}(h-)}\left(1-\exp -\int_{0}^{\sigma} d L_{r}^{D} g\left(\widehat{W}_{r}\right)\right)\right) \\
& \quad=\exp \left(-\int \mu(d h) \mathbb{N}_{\mathrm{w}(h-)}\left(1-e^{-\left\langle\mathcal{Z}^{D}, g\right\rangle}\right)\right) \\
& \quad=\exp \left(-\int \mu(d h) u(\mathrm{w}(h-))\right) .
\end{aligned}
$$

Hence, using also Proposition 4.3.2,

$$
\begin{aligned}
u(x) & =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{D} g\left(\widehat{W}_{s}\right) \exp \left(-\int \rho_{s}(d h) u\left(W_{s}(h-)\right)\right)\right. \\
& =E^{0} \otimes \Pi_{x}\left[1_{\{\tau<\infty\}} e^{-\alpha \tau} g\left(\xi_{\tau}\right) \exp \left(-\int J_{\tau}(d h) u\left(\xi_{h-}\right)\right)\right] \\
& =\Pi_{x}\left[1_{\{\tau<\infty\}} g\left(\xi_{\tau}\right) \exp \left(-\int_{0}^{\tau} d h \tilde{\psi}\left(u\left(\xi_{h}\right)\right)\right)\right] .
\end{aligned}
$$

The integral equation of the theorem now follows by the same routine arguments used in the end of the proof of Proposition 4.2.2.

### 4.4 Continuity properties of the Lévy snake

From now on until the end of this chapter we assume that the underlying spatial motion $\xi$ has continuous sample paths. The construction of Section 4.1 applies with the following minor simplification. Rather than considering càdlàg paths, we can define $\mathcal{W}_{x}$ as the set of all $E$-valued killed continuous paths started at $x$. An element of $\mathcal{W}$ is thus a continuous mapping $\mathrm{w}:[0, \zeta) \longrightarrow E$, and the distance between w and $\mathrm{w}^{\prime}$ is defined by

$$
\begin{equation*}
\left.d\left(\mathrm{w}, \mathrm{w}^{\prime}\right)=\delta\left(\mathrm{w}(0), \mathrm{w}^{\prime}(0)\right)+\left|\zeta-\zeta^{\prime}\right|+\int_{0}^{\zeta \wedge \zeta^{\prime}} d t \sup _{r \leq t} \delta\left(\mathrm{w}(r), \mathrm{w}^{\prime}(r)\right) \wedge 1\right) \tag{4.13}
\end{equation*}
$$

Without risk of confusion, we will keep the same notation as in Section 4.1. The construction developed there goes through without change with these new definitions.

Our goal is to provide conditions on $\psi$ and $\xi$ that will ensure that the process $W_{s}$ is continuous with respect to a distance finer than $d$, which we now introduce. We need to consider stopped paths rather than killed paths. A stopped (continuous) path is a continuous mapping $\mathrm{w}:[0, \zeta] \longrightarrow E$, where $\zeta \geq 0$. When $\zeta=0$, we identify w with $\mathrm{w}(0) \in E$. We denote by $\mathcal{W}^{*}$ the set of all stopped paths in $E$. The set $\mathcal{W}^{*}$
is equipped with the distance

$$
d^{*}\left(\mathrm{w}, \mathrm{w}^{\prime}\right)=\left|\zeta-\zeta^{\prime}\right|+\sup _{t \geq 0} \delta\left(\mathrm{w}(t \wedge \zeta), \mathrm{w}\left(t \wedge \zeta^{\prime}\right)\right)
$$

Note that $\left(\mathcal{W}^{*}, d^{*}\right)$ is a Polish space.
If $\mathrm{w} \in \mathcal{W}$ is a killed path such that $\zeta>0$ and the left limit $\widehat{\mathrm{w}}=\mathrm{w}(\zeta-)$ exists, we write $\mathrm{w}^{*}$ for the corresponding stopped path $\mathrm{w}^{*}(t)=\mathrm{w}(t)$ if $t<\zeta$, and $\mathrm{w}^{*}(\zeta)=\widehat{\mathrm{w}}$. When $\zeta=0$ we make the convention that $x^{*}=x$. Note that $W_{s}^{*}$ is well defined $\mathbb{P}_{x}$ a.s., for every fixed $s \geq 0$.

As in Chapter 1, we set

$$
\gamma=\sup \left\{a \geq 0: \lim _{\lambda \rightarrow \infty} \lambda^{-a} \psi(\lambda)=\infty\right\} \geq 1
$$

Proposition 4.4.1 Suppose that there exist three constants $p>0, q>0$ and $C<\infty$ such that for every $t>0$ and $x \in E$,

$$
\begin{equation*}
\Pi_{x}\left[\sup _{r \leq t} \delta\left(x, \xi_{r}\right)^{p}\right] \leq C t^{q} \tag{4.14}
\end{equation*}
$$

Suppose in addition that

$$
q\left(1-\frac{1}{\gamma}\right)>1
$$

Then the left limit $\widehat{W}_{s}=W_{s}\left(H_{s}-\right)$ exists for every $s \geq 0, \mathbb{P}_{x}$ a.s. or $\mathbb{N}_{x}$ a.e. Furthermore the process $\left(W_{s}^{*}, s \geq 0\right)$ has continuous sample paths with respect to the distance $d^{*}, \mathbb{P}_{x}$ a.s. or $\mathbb{N}_{x}$ a.e.

Remark. Only the small values of $t$ are relevant in our assumption (4.14) since we can always replace the distance $\delta$ by $\delta \wedge 1$. Uniformity in $x$ could also be relaxed, but we do not strive for the best conditions.
Proof. It is enough to argue under $\mathbb{P}_{x}$. Let us fix $t \geq 0$ and $s \in(0,1)$. Then,
$\mathbb{E}_{x}\left[d^{*}\left(W_{t}^{*}, W_{t+s}^{*}\right)^{p}\right] \leq 2^{p}\left(\mathbb{E}_{x}\left[\left|H_{t+s}-H_{t}\right|^{p}\right]+\mathbb{E}_{x}\left[\sup _{r \geq 0} \delta\left(W_{t}^{*}\left(r \wedge H_{t}\right), W_{t+s}^{*}\left(r \wedge H_{t+s}\right)\right)^{p}\right]\right)$.
To simplify notation, set $m=m_{H}(t, t+s)=\inf _{[t, t+s]} H_{r}$. From the conditional distribution of the process $W$ knowing $H$, we easily get

$$
\begin{aligned}
\mathbb{E}_{x} & {\left[\sup _{u \geq 0} \delta\left(W_{t}^{*}\left(u \wedge H_{s}\right), W_{t+s}^{*}\left(u \wedge H_{t+s}\right)\right)^{p} \mid H_{r}, r \geq 0\right] } \\
& \leq 2^{p}\left(\Pi_{x}\left[\Pi_{\xi_{m}}\left[\sup _{u \leq H_{t}-m} \delta\left(\xi_{0}, \xi_{u}\right)^{p}\right]\right]+\Pi_{x}\left[\Pi_{\xi_{m}}\left[\sup _{u \leq H_{t+s}-m} \delta\left(\xi_{0}, \xi_{u}\right)^{p}\right]\right]\right) \\
& \leq C 2^{p}\left(\left|H_{t}-m\right|^{q}+\left|H_{t+s}-m\right|^{q}\right)
\end{aligned}
$$

using our assumption (4.14) in the last bound. By combining the previous estimates with Lemma 1.4.6, we arrive at

$$
\mathbb{E}_{x}\left[d^{*}\left(W_{t}^{*}, W_{t+s}^{*}\right)^{p}\right] \leq 2^{2 p+1}\left(C_{p} \varphi(1 / s)^{-p}+C C_{q} \varphi(1 / s)^{-q}\right),
$$

where $\varphi(\lambda)=\lambda / \psi^{-1}(\lambda)$. Now choose $\alpha \in\left(0,1-\frac{1}{\gamma}\right)$ such that $q \alpha>1$. Notice that we may also assume $p \alpha>1$ since by replacing the distance $\delta$ by $\delta \wedge 1$, we can take $p$ as large as we wish. The condition $\alpha<1-\frac{1}{\gamma}$ and the definition of $\gamma$ imply that
$\varphi(\lambda) \geq c \lambda^{\alpha}$ for every $\lambda \geq 1$, for some constant $c>0$. Hence, there exists a constant $C^{\prime}$ independent of $t$ and $s$ such that

$$
\mathbb{E}_{x}\left[d^{*}\left(W_{t}^{*}, W_{t+s}^{*}\right)^{p}\right] \leq C^{\prime}\left(s^{p \alpha}+s^{q \alpha}\right)
$$

The Kolmogorov lemma then gives the existence of a continuous modification of the process $\left(W_{s}^{*}, s \geq 0\right)$ with respect to the distance $d^{*}$. The various assertions of the proposition follow easily, recalling that we already know that the process ( $W_{s}, s \geq 0$ ) has continuous paths for the distance $d$.

### 4.5 The Brownian motion case

In this section, we concentrate on the case when the underlying spatial motion $\xi$ is Brownian motion in $\mathbb{R}^{d}$. We will give a necessary and sufficient condition for the process $W^{*}$ to have a modification that is continuous with respect to the distance $d^{*}$.

To this end, we introduce the following condition on $\psi$ :

$$
\begin{equation*}
\int_{1}^{\infty}\left(\int_{0}^{t} \psi(u) d u\right)^{-1 / 2} d t<\infty \tag{4.15}
\end{equation*}
$$

Note that this condition is stronger than the condition $\int_{1}^{\infty} d u / \psi(u)<\infty$ for the path continuity of $H$. In fact, since $\psi$ is convex, there exists a positive constant $c$ such $\psi(t) \geq c t$ for every $t \geq 1$. Then, for $t \geq 1$,

$$
\int_{0}^{t} \psi(u) d u \leq t \psi(t) \leq c^{-1} \psi(t)^{2}
$$

and thus

$$
\int_{1}^{\infty} \frac{d u}{\psi(u)} \leq c^{-1 / 2} \int_{1}^{\infty}\left(\int_{0}^{t} \psi(u) d u\right)^{-1 / 2} d t
$$

Also note that (4.15) holds if $\gamma>1$. On the other hand, it is easy to produce examples where (4.15) does not hold although $H$ has continuous sample paths.

Condition (4.15) was introduced in connection with solutions of $\Delta u=\psi(u)$ in domains of $\mathbb{R}^{d}$. We briefly review the results that will be relevant to our needs (see [25],[41] and also Lemma 2.3 in [45]). We denote by $B_{r}$ the open ball of radius $r$ centered at the origin in $\mathbb{R}^{d}$.
A. If (4.15) holds, then, for every $r>0$, there exists a positive solution of the problem

$$
\left\{\begin{array}{l}
\frac{1}{2} \Delta u=\psi(u) \quad \text { in } B_{r}  \tag{4.16}\\
u_{\mid \partial B_{r}}=\infty .
\end{array}\right.
$$

Here the condition $u_{\mid \partial B_{r}}=\infty$ means that $u(x)$ tends to $+\infty$ as $x \rightarrow y, x \in B_{r}$, for every $y \in \partial B_{r}$.
B. If (4.15) does not hold, then for every $c>0$, there exists a positive solution of the problem

$$
\left\{\begin{array}{l}
\frac{1}{2} \Delta u=\psi(u) \quad \text { in } \mathbb{R}^{d}  \tag{4.17}\\
u(0)=c .
\end{array}\right.
$$

Connections between the Lévy snake and the partial differential equation $\Delta u=$ $\psi(u)$ follow from Theorem 4.3.3. Note that this is just a reformulation of the wellknown connections involving superprocesses. We use the notation of Section 4.3. A domain $D$ in $\mathbb{R}^{d}$ is regular if every point $y$ of $\partial D$ is regular for $D^{c}$, meaning that: $\inf \left\{t>0: \xi_{t} \notin D\right\}=0, \Pi_{y}$ a.s.

Proposition 4.5.1 Assume that $\xi$ is Brownian motion in $\mathbb{R}^{d}$. Let $D$ be a bounded regular domain in $\mathbb{R}^{d}$, and let $g$ be a nonnegative continuous function on $\partial D$. Then the function

$$
u(x)=\mathbb{N}_{x}\left(1-\exp -\left\langle\mathcal{Z}^{D}, g\right\rangle\right)
$$

is twice continuously differentiable in $D$ and is the unique nonnegative solution of the problem

$$
\left\{\begin{array}{l}
\frac{1}{2} \Delta u=\psi(u) \quad \text { in } D  \tag{4.18}\\
u_{\mid \partial D}=g .
\end{array}\right.
$$

Proof. This follows from Theorem 4.3 .3 by standard arguments. In the context of superprocesses, the result is due to Dynkin [12]. See e.g. Chapter 5 in [31] for a proof in the case $\psi(u)=u^{2}$, which is readily extended.

We can now state our main result.
Theorem 4.5.2 Assume that $\xi$ is Brownian motion in $\mathbb{R}^{d}$. The following conditions are equivalent.
(i) $\mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)<\infty$ for some $r>0$.
(ii) $\mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)<\infty$ for every $r>0$.
(iii) The left limit $\widehat{W}_{s}=W_{s}\left(\zeta_{s}-\right)$ exists for every $s \geq 0, \mathbb{P}_{0}$ a.s., and the mapping $s \rightarrow \widehat{W}_{s}$ is continuous, $\mathbb{P}_{0}$ a.s.
(iv) The left limit $\widehat{W}_{s}=W_{s}\left(\zeta_{s}-\right)$ exists for every $s \geq 0, \mathbb{P}_{0}$ a.s., and the mapping $s \rightarrow W_{s}^{*}$ is continuous for the metric $d^{*}, \mathbb{P}_{0}$ a.s.
(v) Condition (4.15) holds.

Remark. The conditions of Theorem 4.5.2 are also equivalent to the a.s. compactness of the range of the superprocess with spatial motion $\xi$ and branching mechanism $\psi$, started at a nonzero initial measure $\mu$ with compact support. This fact, that follows from Theorem 5.1 in [45], can be deduced from the representation of Theorem 4.2.1.

Proof. The equivalence between (i),(ii) and (v) is easy given facts A. and B. recalled above. We essentially reproduce arguments of [45]. By fact B., if (v) does not hold, then we can for every $c>0$ find a nonnegative function $v_{c}$ such that $v_{c}(0)=c$ and $\frac{1}{2} \Delta v_{c}=\psi\left(v_{c}\right)$ in $\mathbb{R}^{d}$. Let $r>0$ and $\lambda>0$. By Proposition 4.5.1, the nonnegative function

$$
u_{\lambda, r}(x)=\mathbb{N}_{x}\left(1-\exp -\lambda\left\langle\mathcal{Z}^{B_{r}}, 1\right\rangle\right), \quad x \in B_{r}
$$

solves $\frac{1}{2} \Delta u_{\lambda, r}=\psi\left(u_{\lambda, r}\right)$ in $B_{r}$ with boundary condition $\lambda$ on $\partial B_{r}$. By choosing $\lambda$ sufficiently large so that $\sup \left\{v_{c}(y), y \in \partial B_{r}\right\}<\lambda$, and using the comparison principle for nonnegative solutions of $\frac{1}{2} \Delta u=\psi(u)$ (see Lemma V. 7 in [31]), we see that $v_{c} \leq$ $u_{\lambda, r}$ in $B_{r}$. In particular,

$$
c=v_{c}(0) \leq u_{\lambda, r}(0) \leq \mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)
$$

Since $c$ was arbitrary, we get $\mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)=\infty$ and we have proved that $(\mathrm{i}) \Rightarrow(\mathrm{v})$. Trivially (ii) $\Rightarrow$ (i).

Suppose now that (v) holds. Let $r>0$. By fact A., we can find a function $u_{(r)}$ such that $\frac{1}{2} \Delta u_{(r)}=\psi\left(u_{(r)}\right)$ in $B_{r}$ with boundary condition $+\infty$ on $\partial B_{r}$. The maximum principle then implies that, for every $\lambda>0, u_{\lambda, r} \leq u_{(r)}$. Hence

$$
\mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)=\lim _{\lambda \uparrow \infty} \uparrow u_{\lambda, r}(0) \leq u_{(r)}(0)<\infty
$$

and (ii) holds. We have thus proved the equivalence of (i),(ii) and (v).
Let us prove that (iii) $\Rightarrow$ (ii). We assume that (iii) holds. Let $r>0$. Then on the event $\left\{\mathcal{Z}^{B_{r}} \neq 0\right\}$, there exists $s \in(0, \sigma)$ such that $\widehat{W}_{s} \in \partial B_{r}$. It follows that

$$
\mathbb{N}_{0}\left(\sup _{s \in(0, \sigma)}\left|\widehat{W}_{s}\right| \geq r\right) \geq \mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)
$$

Let $T_{1}=\inf \left\{s \geq 0: L_{s}^{0}>1\right\}$ (in agreement with the notation of Chapter 1). The path continuity of $\widehat{W}_{s}$ ensures that $\mathbb{P}_{0}$ a.s. there are only finitely many excursions intervals $\left(\alpha_{i}, \beta_{i}\right)$ of $H_{s}$ away from 0 , before time $T_{1}$, such that

$$
\sup _{s \in\left(\alpha_{i}, \beta_{i}\right)}\left|\widehat{W}_{s}\right| \geq r
$$

On the other hand, excursion theory implies that the number of such intervals is Poisson with parameter

$$
\mathbb{N}_{0}\left(\sup _{s \in(0, \sigma)}\left|\widehat{W}_{s}\right| \geq r\right)
$$

We conclude that the latter quantity is finite, and so $\mathbb{N}_{0}\left(\mathcal{Z}^{B_{r}} \neq 0\right)<\infty$.
Note that (iv) $\Rightarrow$ (iii). Thus, to complete the proof of Theorem 4.5.2, it remains to verify that (ii) $\Rightarrow$ (iv). From now on until the end of the proof, we assume that (ii) holds.

We use the following simple lemma.
Lemma 4.5.3 Let $D$ be a domain in $\mathbb{R}^{d}$ containing 0 , and let

$$
S=\inf \left\{s \geq 0: W_{s}(t) \notin D \text { for some } t \in\left[0, H_{s}\right)\right\}
$$

Then $\mathbb{N}_{0}\left(\mathcal{Z}^{D} \neq 0\right) \geq \mathbb{N}_{0}(S<\infty)$.
Proof. By excursion theory, we have

$$
\mathbb{P}_{0}\left[S \geq T_{1}\right]=\exp \left(-\mathbb{N}_{0}(S<\infty)\right)
$$

Then, let $\tau$ be as previously the exit time from $D$. If there exists $s<T_{1}$ such that $\tau\left(W_{s}\right)<H_{s}$, then the same property holds for every $s^{\prime}>s$ such that $s^{\prime}-s$ is sufficiently small, by the continuity of $H$ and the snake property. Hence,

$$
\left\{S<T_{1}\right\} \subset\left\{\int_{0}^{T_{1}} d s 1_{\left\{\tau\left(W_{s}\right)<H_{s}\right\}}>0\right\}, \quad \mathbb{P}_{0} \text { a.e. }
$$

It follows that

$$
\mathbb{P}_{0}\left[S \geq T_{1}\right] \geq \mathbb{P}_{0}\left[\int_{0}^{T_{1}} d s 1_{\left\{\tau\left(W_{s}\right)<H_{s}\right\}}=0\right]=\mathbb{P}_{0}\left[L_{T_{1}}^{D}=0\right]
$$

where the second equality is a consequence of the formula

$$
L_{T_{1}}^{D}=\ell^{D}\left(\int_{0}^{T_{1}} d s 1_{\left\{\tau\left(W_{s}\right)<H_{s}\right\}}\right)
$$

together with the fact that $\ell^{D}(s)>0$ for every $s>0$, a.s.
Using again excursion theory and the construction of the exit measure under $\mathbb{N}_{0}$, we get

$$
\mathbb{P}_{0}\left[S \geq T_{1}\right] \geq \mathbb{P}_{0}\left[L_{T_{1}}^{D}=0\right]=\exp \left(-\mathbb{N}_{0}\left(\mathcal{Z}^{D} \neq 0\right)\right)
$$

By comparing with the first formula of the proof, we get the desired inequality.
Let $\varepsilon>0$. We specialize the previous lemma to the case $D=B_{\varepsilon}$ and write $S=S_{1}^{\varepsilon}$. Then, for every $r>0$,

$$
\mathbb{P}_{0}\left[S_{1}^{\varepsilon} \geq T_{r}\right]=\exp \left(-r \mathbb{N}_{0}\left(S_{1}^{\varepsilon}<\infty\right)\right) \geq \exp \left(-r \mathbb{N}_{0}\left(\mathcal{Z}^{B_{\varepsilon}} \neq 0\right)\right)
$$

From (ii), it follows that $S_{1}^{\varepsilon}>0, \mathbb{P}_{0}$ a.e. Also note that $S_{1}^{\varepsilon}$ is a stopping time of the filtration $\left(\mathcal{F}_{s+}\right)$ and that $\left\{W_{S_{1}^{\varepsilon}}(t): 0 \leq t<H_{S_{1}^{\varepsilon}}\right\} \subset \bar{B}_{\varepsilon}$ (if this inclusion were not true, the snake property would contradict the definition of $\left.S_{1}^{\varepsilon}\right)$.

Recall the notation $m_{H}\left(s, s^{\prime}\right)=\inf _{\left[s, s^{\prime}\right]} H_{r}$ for $s \leq s^{\prime}$. We define inductively a sequence $\left(S_{n}^{\varepsilon}\right)_{n \geq 1}$ of stopping times (for the filtration $\left(\mathcal{F}_{s+}\right)$ ) by setting

$$
S_{n+1}^{\varepsilon}=\inf \left\{s>S_{n}^{\varepsilon}:\left|W_{s}(t)-W_{s}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s\right)\right)\right|>\varepsilon \text { for some } t \in\left[0, H_{s}\right)\right\}
$$

At this point we need another lemma.
Lemma 4.5.4 Let $T$ be a stopping time of the filtration $\left(\mathcal{F}_{s_{+}}\right)$, such that $T<\infty, \mathbb{P}_{0}$ a.s. For every $s \geq 0$, define a killed path $\tilde{W}_{s}$ with lifetime $\tilde{H}_{s}$ by setting
$\tilde{W}_{s}(t)=W_{s}\left(m_{H}(T, T+s)+t\right)-W_{s}\left(m_{H}(T, T+s)\right), 0 \leq t<\tilde{H}_{s}:=H_{T+s}-m_{H}(T, T+s)$
with the convention that $\tilde{W}_{s}=0$ if $\tilde{H}_{s}=0$. Then the process $\left(\tilde{W}_{s}, s \geq 0\right)$ is independent of $\mathcal{F}_{T+}$ and has the same distribution as $\left(W_{s}, s \geq 0\right)$ under $\mathbb{P}_{0}$.

This lemma follows from the strong Markov property of the Lévy snake, together with Lemma 4.2.4. The translation invariance of the spatial motion is of course crucial here.

As a consequence of the preceding lemma, we obtain that the random variables $S_{1}^{\varepsilon}, S_{2}^{\varepsilon}-S_{1}^{\varepsilon}, \ldots, S_{n+1}^{\varepsilon}-S_{n}^{\varepsilon}, \ldots$ are independent and identically distributed. Recall that these variables are positive a.s. Also observe that

$$
\left\{W_{S_{n+1}^{\varepsilon}}(t)-W_{S_{n+1}^{\varepsilon}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, S_{n+1}^{\varepsilon}\right)\right): t \in\left[0, H_{S_{n+1}^{\varepsilon}}\right)\right\} \subset \bar{B}_{\varepsilon},
$$

by the same argument as used previously for $\left\{W_{S_{1}^{\varepsilon}}(t): t \in\left[0, H_{S_{1}^{\varepsilon}}\right)\right\}$.
Let $a>0$. We claim that $\mathbb{N}_{0}$ a.s. we can choose $\delta_{1}>0$ small enough so that, for every $s, s^{\prime} \in\left[0, T_{a}\right]$ such that $s \leq s^{\prime} \leq s+\delta_{1}$,

$$
\begin{equation*}
\left|W_{s^{\prime}}(t)-W_{s^{\prime}}\left(m_{H}\left(s, s^{\prime}\right) \wedge t\right)\right| \leq 3 \varepsilon, \quad \text { for every } t \in\left[0, H_{s^{\prime}}\right) \tag{4.19}
\end{equation*}
$$

Let us verify that the claim holds if we take

$$
\delta_{1}=\inf \left\{S_{n+1}^{\varepsilon}-S_{n}^{\varepsilon} ; n \geq 1, S_{n}^{\varepsilon} \leq T_{a}\right\}>0 .
$$

Consider $s, s^{\prime} \in\left[0, T_{a}\right]$ with $s \leq s^{\prime} \leq s+\delta_{1}$. Then two cases may occur.
Either $s, s^{\prime}$ belong to the same interval $\left[S_{n}^{\varepsilon}, S_{n+1}^{\varepsilon}\right]$. Then, from the definition of $S_{n+1}^{\varepsilon}$ we know that

$$
\begin{equation*}
\left|W_{s^{\prime}}(t)-W_{s^{\prime}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)\right| \leq \varepsilon \text { for every } t \in\left[0, H_{s^{\prime}}\right) \tag{4.20}
\end{equation*}
$$

Since $m_{H}\left(s, s^{\prime}\right) \geq m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)$ we can replace $t$ by $t \wedge m\left(s, s^{\prime}\right)$ to get

$$
\left|W_{s^{\prime}}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)-W_{s^{\prime}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)\right| \leq \varepsilon \text { for every } t \in\left[0, H_{s^{\prime}}\right)
$$

and our claim (4.19) follows by combining this bound with the previous one.

Then we need to consider the case where $s \in\left[S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right]$ and $s^{\prime} \in\left(S_{n}^{\varepsilon}, S_{n+1}^{\varepsilon}\right]$ for some $n \geq 1$ (by convention $S_{0}^{\varepsilon}=0$ ). If $m_{H}\left(s, s^{\prime}\right)=m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)$, then the same argument as in the first case goes through. Therefore we can assume that $m_{H}\left(s, s^{\prime}\right)<$ $m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)$, which implies $m_{H}\left(S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right)<m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)$. Note that the bound (4.20) still holds. We also know that

$$
\begin{equation*}
\left|W_{S_{n}^{\varepsilon}}(t)-W_{S_{n}^{\varepsilon}}\left(t \wedge m_{H}\left(S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right)\right)\right| \leq \varepsilon \text { for every } t \in\left[0, H_{S_{n}^{\varepsilon}}\right) \tag{4.21}
\end{equation*}
$$

We replace $t$ by $t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)$ in this bound, and note that $W_{S_{n}^{\varepsilon}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)=$ $W_{s^{\prime}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)$ for every $t \in\left[0, H_{S_{n}^{\varepsilon}} \wedge H_{s^{\prime}}\right)$, by the snake property. It follows that

$$
\begin{equation*}
\left|W_{s^{\prime}}\left(t \wedge m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)-W_{S_{n}^{\varepsilon}}\left(t \wedge m_{H}\left(S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right)\right)\right| \leq \varepsilon \text { for every } t \in\left[0, H_{s^{\prime}}\right) \tag{4.22}
\end{equation*}
$$

Similarly, we can replace $t$ by $t \wedge m_{H}\left(s, s^{\prime}\right)$ in (4.21), using again the snake property to write $W_{S_{n}^{\varepsilon}}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)=W_{s^{\prime}}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)\left(\right.$ note that $m_{H}\left(S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right) \leq m_{H}\left(s, s^{\prime}\right)<$ $\left.m_{H}\left(S_{n}^{\varepsilon}, s^{\prime}\right)\right)$. It follows that

$$
\begin{equation*}
\left|W_{s^{\prime}}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)-W_{S_{n}^{\varepsilon}}\left(t \wedge m_{H}\left(S_{n-1}^{\varepsilon}, S_{n}^{\varepsilon}\right)\right)\right| \leq \varepsilon \text { for every } t \in\left[0, H_{s^{\prime}}\right) . \tag{4.23}
\end{equation*}
$$

Our claim (4.19) is now a consequence of (4.20), (4.22) and (4.23).
We can already derive from (4.19) the fact that the left limit $\widehat{W}_{s}$ exists for every $s \in\left[0, T_{a}\right], \mathbb{P}_{0}$ a.s. We know that this left limit exists for every rational $s \in\left[0, T_{a}\right]$, $\mathbb{P}_{0}$ a.s. Let $s \in\left(0, T_{a}\right]$, and let $s_{n}$ be a sequence of rationals increasing to $s$. Then the sequence $m_{H}\left(s_{n}, s\right)$ also increases to $H_{s}$. If $m_{H}\left(s_{n}, s\right)=H_{s}$ for some $n$, then the snake property shows that $W_{s}(t)=W_{s_{n}}(t)$ for every $t \in\left[0, H_{s}\right)$ and the existence of $\widehat{W}_{s}$ is an immediate consequence. Otherwise, (4.19) shows that for $n$ large enough,

$$
\sup _{t \in\left[0, H_{s}\right)}\left|W_{s}(t)-W_{s}\left(t \wedge m_{H}\left(s_{n}, s\right)\right)\right| \leq 3 \varepsilon
$$

and by applying this to a sequence of values of $\varepsilon$ tending to 0 we also get the existence of $\widehat{W}_{s}$.

We finally use a time-reversal argument. From Corollary 3.1.6, we know that the processes $\left(H_{t \wedge T_{a}}, t \geq 0\right)$ and $\left(H_{\left(T_{a}-t\right)^{+}}, t \geq 0\right)$ have the same distribution. By considering the conditional law of $W$ knowing $H$, we immediately obtain that the processes $\left(W_{t \wedge T_{a}}, t \geq 0\right)$ and $\left(W_{\left(T_{a}-t\right)^{+}}, t \geq 0\right)$ also have the same distribution. Thanks to this observation and the preceding claim, we get that $\mathbb{P}_{0}$ a.s. there exists $\delta_{2}>0$ such that for every $s, s^{\prime} \in\left[0, T_{a}\right]$ with $s \leq s^{\prime} \leq s+\delta_{2}$,

$$
\begin{equation*}
\left|W_{s}(t)-W_{s}\left(m_{H}\left(s, s^{\prime}\right) \wedge t\right)\right| \leq 3 \varepsilon, \quad \text { for every } t \in\left[0, H_{s}\right) \tag{4.24}
\end{equation*}
$$

To complete the proof, note that the snake property implies that

$$
W_{s}^{*}\left(m_{H}\left(s, s^{\prime}\right)\right)=W_{s^{\prime}}^{*}\left(m_{H}\left(s, s^{\prime}\right)\right)
$$

using a continuity argument in the case $m_{H}\left(s, s^{\prime}\right)=H_{s} \wedge H_{s^{\prime}}$. Thus, if $s \leq s^{\prime} \leq T_{a}$ and $s^{\prime}-s \leq \delta_{1} \wedge \delta_{2}$,

$$
\begin{aligned}
& \sup _{t \geq 0}\left|W_{s}^{*}\left(t \wedge H_{s}\right)-W_{s^{\prime}}^{*}\left(t \wedge H_{s^{\prime}}\right)\right| \\
& \leq \sup _{t \in\left[0, H_{s}\right]}\left|W_{s}^{*}(t)-W_{s}^{*}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)\right|+\sup _{t \in\left[0, H_{s^{\prime}}\right]}\left|W_{s^{\prime}}^{*}(t)-W_{s^{\prime}}^{*}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)\right| \\
& =\sup _{t \in\left[0, H_{s}\right)}\left|W_{s}(t)-W_{s}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)\right|+\sup _{t \in\left[0, H_{s^{\prime}}\right)}\left|W_{s^{\prime}}(t)-W_{s^{\prime}}\left(t \wedge m_{H}\left(s, s^{\prime}\right)\right)\right| \\
& \leq 6 \varepsilon .
\end{aligned}
$$

This gives the continuity of the mapping $s \rightarrow W_{s}^{*}$ with respect to the distance $d^{*}$, and completes the proof of (iv).

### 4.6 The law of the Lévy snake at a first exit time

Our goal in this section is to give explicit formulas for the law of the Lévy snake at its first exit time from a domain. We keep assuming that $H$ has continuous sample paths and in addition we suppose that the process $W^{*}$ has continuous sample paths with respect to the metric $d^{*}$. Note that the previous two sections give sufficient conditions for this property to hold.

Let $D$ be an open set in $E$ and $x \in D$. We slightly abuse notation by writing $\tau(\mathrm{w})=\inf \left\{t \in\left[0, \zeta_{\mathrm{w}}\right]: \mathrm{w}(t) \notin D\right\}$ for any stopped path $\mathrm{w} \in \mathcal{W}^{*}$. We also set

$$
T_{D}=\inf \left\{s>0: \tau\left(W_{s}^{*}\right)<\infty\right\} .
$$

The continuity of $W^{*}$ with respect to the metric $d^{*}$ immediately implies that $T_{D}>0$, $\mathbb{N}_{x}$ a.e. or $\mathbb{P}_{x}$ a.e. Furthermore, on the event $\left\{T_{D}<\infty\right\}$ the path $W_{T_{D}}^{*}$ hits the boundary of $D$ exactly at its lifetime. The main result of this section determines the law of the pair ( $\rho_{T_{D}}, W_{T_{D}}$ ) under $\mathbb{N}_{x}\left(\cdot \cap\left\{T_{D}<\infty\right\}\right)$.

Before stating this result, we need some notation and a preliminary lemma. For every $y \in D$, we set

$$
u(y)=\mathbb{N}_{y}\left(T_{D}<\infty\right)<\infty
$$

Recall that, for every $a, b \geq 0$, we have defined

$$
\gamma_{\psi}(a, b)= \begin{cases}(\psi(a)-\psi(b)) /(a-b) & \text { if } a \neq b \\ \psi^{\prime}(a) & \text { if } a=b\end{cases}
$$

Note that $\gamma_{\psi}(a, 0)=\tilde{\psi}(a)$ (by convention $\tilde{\psi}(0)=\psi^{\prime}(0)=\alpha$ ). The following formulas will be useful: For every $a, b \geq 0$,

$$
\begin{align*}
& \int \pi(d r) \int_{0}^{r} d \ell\left(1-e^{-a \ell-b(r-\ell)}\right)=\gamma_{\psi}(a, b)-\alpha-\beta(a+b)  \tag{4.25}\\
& \int \pi(d r) \int_{0}^{r} d \ell e^{-a \ell}\left(1-e^{-b(r-\ell)}\right)=\gamma_{\psi}(a, b)-\tilde{\psi}(a)-\beta b \tag{4.26}
\end{align*}
$$

The first formula is easily obtained by observing that, if $a \neq b$,

$$
\int_{0}^{r} d \ell\left(1-e^{-\ell a-(r-\ell) b}\right)=\frac{1}{a-b}\left(r(a-b)+\left(e^{-r a}-e^{-r b}\right)\right) .
$$

The second one is a consequence of the first one and the identity

$$
\tilde{\psi}(a)=\alpha+\beta a+\int \pi(d r) \int_{0}^{r} d \ell\left(1-e^{-a \ell}\right)
$$

Recall from Section 3.1 the definition of the probability measures $\mathbb{M}_{a}$ on $M_{f}\left(\mathbb{R}_{+}\right)^{2}$.
Lemma 4.6.1 (i) Let $a>0$ and let $F$ be a nonnegative measurable function on $M_{f}\left(\mathbb{R}_{+}\right) \times M_{f}\left(\mathbb{R}_{+}\right) \times \mathcal{W}$. Then,

$$
\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} F\left(\rho_{s}, \eta_{s}, W_{s}\right)\right)=e^{-\alpha a} \int \mathbb{M}_{a}(d \mu d \nu) \Pi_{x}\left[F\left(\mu, \nu,\left(\xi_{r}, 0 \leq r<a\right)\right)\right] .
$$

(ii) Let $f, g$ be two nonnegative measurable functions on $\mathbb{R}_{+}$. Then,

$$
N\left(\int_{0}^{\sigma} d L_{s}^{a} \exp \left(-\left\langle\rho_{s}, f\right\rangle-\left\langle\eta_{s}, g\right\rangle\right)\right)=\exp \left(-\int_{0}^{a} \gamma_{\psi}(f(t), g(t)) d t\right)
$$

Proof. (i) As in the proof of Lemma 4.2.5, we may restrict our attention to a function $F\left(\rho_{s}, \eta_{s}, W_{s}\right)=F\left(\rho_{s}, \eta_{s}\right)$. Then the desired result follows from Corollary 1.3.4 in the same way as Proposition 3.1.3 was deduced from Proposition 1.1.4.
(ii) By part (i) we have

$$
N\left(\int_{0}^{\sigma} d L_{s}^{a} \exp \left(-\left\langle\rho_{s}, f\right\rangle-\left\langle\eta_{s}, g\right\rangle\right)\right)=e^{-\alpha a} \int \mathbb{M}_{a}(d \mu d \nu) \exp (-\langle\mu, f\rangle-\langle\nu, g\rangle) .
$$

From the definition of $\mathbb{M}_{a}$ this is equal to

$$
\exp \left(-\alpha a-\beta \int_{0}^{a}(f(t)+g(t)) d t-\int_{0}^{a} d t \int \pi(d r) \int_{0}^{r} d \ell\left(1-e^{-\ell f(t)-(r-\ell) g(t)}\right)\right) .
$$

The stated result now follows from (4.25).
Theorem 4.6.2 Assume that $u(x)>0$. Let $a>0$, let $F$ be a nonnegative measurable function on $\mathcal{W}_{x}^{*}$ and let $g$ be a nonnegative measurable function on $\mathbb{R}_{+}$with support contained in $[0, a]$. Then

$$
\begin{align*}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} 1_{\left\{a<H_{T_{D}}\right\}} F\left(W_{T_{D}}(t), 0 \leq t \leq a\right) \exp \left(-\left\langle\rho_{T_{D}}, g\right\rangle\right)\right) \\
& \quad=\Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) \exp \left(-\int_{0}^{a} \gamma_{\psi}\left(u\left(\xi_{r}\right), g(r)\right) d r\right)\right] . \tag{4.27}
\end{align*}
$$

Alternatively, the law of $W_{T_{D}}$ under $\mathbb{N}_{x}\left(\cdot \cap\left\{T_{D}<\infty\right\}\right)$ is characterized by:

$$
\begin{align*}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} 1_{\left\{a<H_{\left.T_{D}\right\}}\right.} F\left(W_{T_{D}}(t), 0 \leq t \leq a\right)\right) \\
& \quad=\Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) \exp \left(-\int_{0}^{a} \tilde{\psi}\left(u\left(\xi_{r}\right)\right) d r\right)\right], \tag{4.28}
\end{align*}
$$

and the conditional law of $\rho_{T_{D}}$ knowing $W_{T_{D}}$ is the law of

$$
\beta 1_{\left[0, H_{T_{D}}\right]}(r) d r+\sum_{i \in I}\left(v_{i}-\ell_{i}\right) \delta_{r_{i}}
$$

where $\sum \delta_{\left(r_{i}, v_{i}, \ell_{i}\right)}$ is a Poisson point measure on $\mathbb{R}_{+}^{3}$ with intensity

$$
\begin{equation*}
1_{\left[0, H_{T_{D}}\right]}(r) 1_{[0, v]}(\ell) e^{-\ell u\left(W_{T_{D}}(r)\right)} d r \pi(d v) d \ell . \tag{4.29}
\end{equation*}
$$

Proof. We will rely on results obtained in Section 4.2 above. As in subsection 4.2.2, we denote by $\left(\rho^{i}, W^{i}\right), i \in I$ the "excursions" of the Lévy snake above height $a$. We let $\left(\alpha_{i}, \beta_{i}\right)$ be the time interval corresponding to the excursion $\left(\rho^{i}, W^{i}\right)$ and $\ell^{i}=L_{\alpha_{i}}^{a}$. We also use the obvious notation

$$
T_{D}\left(W^{i}\right)=\inf \left\{s \geq 0: \tau\left(W_{s}^{i *}\right)<\infty\right\}
$$

For every $s \geq 0$, set

$$
G_{s}=1_{\left\{s<T_{D}\right\}} F\left(W_{s}^{*}\right) \exp \left(-\left\langle\rho_{s}, g\right\rangle\right) .
$$

Then it is easy to verify that

$$
\begin{equation*}
\sum_{i \in I} G_{\alpha_{i}} 1_{\left\{T_{D}\left(W^{i}\right)<\infty\right\}}=1_{\left\{T_{D}<\infty\right\}} 1_{\left\{a<H_{T_{D}}\right\}} F\left(W_{T_{D}}(t), 0 \leq t \leq a\right) \exp \left(-\left\langle\rho_{T_{D}}, g\right\rangle\right) . \tag{4.30}
\end{equation*}
$$

In fact, the sum in the left side contains at most one nonzero term, and exactly one iff $T_{D}<\infty$ and $a<H_{T_{D}}$. On this event, $T_{D}$ belongs to one excursion interval above height $a$, say $\left(\alpha_{j}, \beta_{j}\right)$, and then the restriction of $\rho_{T_{D}}$ to $[0, a]$ coincides with $\rho_{\alpha_{j}}$ (see the second step of the proof of Proposition 1.3.1), whereas the snake property ensures that the paths $W_{T_{D}}$ and $W_{\alpha_{i}}$ are the same over $[0, a)$. Our claim (4.30) follows.

Recall the notation $\tilde{W}, \tilde{\rho}, \gamma^{a}(\ell)$ introduced in the proof of Proposition 4.2.3. The proof of this proposition shows that conditionally on the $\sigma$-field $\mathcal{E}_{a}$, the point measure

$$
\sum_{i \in I} \delta_{\left(\ell^{i}, \rho^{i}, W^{i}\right)}
$$

is Poisson with intensity

$$
1_{\left[0, L_{\sigma}^{a}\right]}(\ell) d \ell \mathbb{N}_{\hat{\tilde{W}}_{\gamma^{a}(\ell)}}(d \rho d W) .
$$

Note that the statement of Proposition 4.2 .3 is slightly weaker than this, but the preceding assertion follows readily from the proof.

We now claim that we can find a deterministic function $\Delta$ and an $\mathcal{E}_{a}$-measurable random variable $Z$ such that, for every $j \in I$, we have

$$
\begin{equation*}
G_{\alpha_{j}}=\Delta\left(Z, \ell^{j},\left(\ell^{i}, W^{i}\right)_{i \in I}\right) . \tag{4.31}
\end{equation*}
$$

Precisely, this relation holds if we take for every $\ell \geq 0$,

$$
\begin{gathered}
\Delta\left(Z, \ell,\left(\ell^{i}, W^{i}\right)_{i \in I}\right)=\left(\prod_{i: \ell^{i}<\ell} 1_{\left\{T_{D}\left(W^{i}\right)=\infty\right\}}\right) 1_{\left\{\tilde{W}_{r}(t) \in D, \forall r \in\left[0, \gamma^{a}(\ell)\right], t \in[0, a]\right\}} \\
\times F\left(\tilde{W}_{\gamma^{a}(\ell)}^{*}\right) \exp \left(-\left\langle\tilde{\rho}_{\gamma^{a}(\ell)}, g\right\rangle\right) .
\end{gathered}
$$

Note that the right side of the last formula depends on $\ell$, on $\left(\ell^{i}, W^{i}\right)_{i \in I}$ and on the triple $\left(\tilde{W}, \tilde{\rho}, \gamma^{a}\right)$ which is $\mathcal{E}_{a}$-measurable, and thus can be written in the form of the left side. Then, to justify (4.31), note that

$$
1_{\left\{\alpha_{j}<T_{D}\right\}}=\left(\prod_{i: \ell^{i}<\ell^{j}} 1_{\left\{T_{D}\left(W^{i}\right)=\infty\right\}}\right) 1_{\left\{\tilde{W}_{r}(t) \in D, \forall r \in\left[0, \gamma^{a}\left(\ell^{j}\right)\right], t \in[0, a]\right\}},
$$

since $\gamma^{a}\left(\ell^{j}\right)=\int_{0}^{\alpha_{j}} d r 1_{\left\{H_{r} \leq a\right\}}$ as observed in the proof of Proposition 4.2.3. The latter proof also yields the identities

$$
W_{\alpha_{j}}=W_{\beta_{j}}=\tilde{W}_{\gamma^{a}\left(\ell^{j}\right)}, \quad \rho_{\alpha_{j}}=\rho_{\beta_{j}}=\tilde{\rho}_{\gamma^{a}\left(\ell^{j}\right)}
$$

from which (4.31) follows.
Then, by an application of Lemma 3.1.5 to the point measure $\sum_{i \in I} \delta_{\left(\ell^{i}, \rho^{i}, W^{i}\right)}$, which is Poisson conditional on $\mathcal{E}_{a}$, we have

$$
\begin{aligned}
& \mathbb{N}_{x}\left(\sum_{j \in I} G_{\alpha_{j}} 1_{\left\{T_{D}\left(W^{j}\right)<\infty\right\}} \mid \mathcal{E}_{a}\right) \\
& \quad=\mathbb{N}_{x}\left(\sum_{j \in I} \Delta\left(Z, \ell^{j},\left(\ell^{i}, W^{i}\right)_{i \in I}\right) 1_{\left\{T_{D}\left(W^{j}\right)<\infty\right\}} \mid \mathcal{E}_{a}\right) \\
& \quad=\mathbb{N}_{x}\left(\int_{0}^{L_{\sigma}^{a}} d \ell \int \mathbb{N}_{\hat{\tilde{W}}_{\gamma^{a}(\ell)}}\left(d \rho^{\prime} d W^{\prime}\right) 1_{\left\{T_{D}\left(W^{\prime}\right)<\infty\right\}} \Delta\left(Z, \ell,\left(\ell^{i}, W^{i}\right)_{i \in I}\right) \mid \mathcal{E}_{a}\right) \\
& \quad=\mathbb{N}_{x}\left(\int_{0}^{L_{\sigma}^{a}} d \ell \Delta\left(Z, \ell,\left(\ell^{i}, W^{i}\right)_{i \in I}\right) \mathbb{N}_{\hat{\tilde{W}}_{\gamma^{a}(\ell)}}\left(T_{D}<\infty\right) \mid \mathcal{E}_{a}\right) .
\end{aligned}
$$

Now use the definition of $\Delta\left(Z, \ell,\left(\ell^{i}, W^{i}\right)_{i \in I}\right)$ to get

$$
\begin{align*}
\mathbb{N}_{x}( & \left.\sum_{j \in I} G_{\alpha_{j}} 1_{\left\{T_{D}\left(W^{j}\right)<\infty\right\}}\right) \\
= & \mathbb{N}_{x}\left(\int_{0}^{L_{\sigma}^{a}} d \ell u\left(\widehat{\tilde{W}}_{\gamma^{a}(\ell)}\right)\left(\prod_{i: \ell^{i}<\ell} 1_{\left\{T_{D}\left(W^{i}\right)=\infty\right\}}\right) 1_{\left\{T_{D}(\tilde{W})>\gamma^{a}(\ell)\right\}}\right. \\
& \left.\times F\left(\tilde{W}_{\gamma^{a}(\ell)}^{*}\right) \exp \left(-\left\langle\tilde{\rho}_{\gamma^{a}(\ell)}, g\right\rangle\right)\right) \\
= & \mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) \exp \left(-\left\langle\rho_{s}, g\right\rangle\right) 1_{\left\{s<T_{D}\right\}}\right) . \tag{4.32}
\end{align*}
$$

The last equality is justified by the change of variables $\ell=L_{s}^{a}$ and the fact that $d L_{s}^{a}$ a.e.,

$$
\tilde{W}_{\gamma^{a}\left(L_{s}^{a}\right)}=\tilde{W}_{\tilde{A}_{s}^{a}}=W_{s}, \quad \tilde{\rho}_{\gamma^{a}\left(L_{s}^{a}\right)}=\tilde{\rho}_{\tilde{A}_{s}^{a}}=\rho_{s},
$$

(where $\tilde{A}_{s}^{a}=\int_{0}^{s} d r 1_{\left\{H_{r} \leq a\right\}}$ as previously) and similarly, $d L_{s}^{a}$ a.e.,

$$
\begin{aligned}
\left(\prod_{i: \ell^{i}<L_{s}^{a}} 1_{\left\{T_{D}\left(W^{i}\right)=\infty\right\}}\right) 1_{\left\{T_{D}(\tilde{W})>\gamma^{a}\left(L_{s}^{a}\right)\right\}} & =1_{\left\{W_{r}^{*}(t) \in D, \forall r \leq s, t \in\left[a, H_{r}\right]\right\}} 1_{\left\{T_{D}(\tilde{W})>\tilde{A}_{s}^{a}\right\}} \\
& =1_{\left\{s<T_{D}\right\}} .
\end{aligned}
$$

To evaluate the right side of (4.32), we use a duality argument. It follows from Corollary 3.1.6 and the construction of the Lévy snake that the triples

$$
\left(\rho_{s}, L_{s}^{a}, W_{s} ; 0 \leq s \leq \sigma\right)
$$

and

$$
\left(\eta_{(\sigma-s)-}, L_{\sigma}^{a}-L_{\sigma-s}^{a}, W_{\sigma-s} ; 0 \leq s \leq \sigma\right)
$$

have the same distribution under $\mathbb{N}_{x}$. From this we get

$$
\begin{align*}
& \mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) \exp \left(-\left\langle\rho_{s}, g\right\rangle\right) 1_{\left\{s<T_{D}\right\}}\right) \\
& =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) \exp \left(-\left\langle\eta_{s}, g\right\rangle\right) 1_{\left\{\tau\left(W_{r}^{*}\right)=\infty, \forall r \geq s\right\}}\right) . \tag{4.33}
\end{align*}
$$

Now we can use the strong Markov property of the Lévy snake (as in the second step of the proof of Proposition 4.2.2), and then Lemma 4.2.4, to get

$$
\begin{align*}
& \mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) \exp \left(-\left\langle\eta_{s}, g\right\rangle\right) 1_{\left\{\tau\left(W_{r}^{*}\right)=\infty, \forall r \geq s\right\}}\right) \\
& =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) \exp \left(-\left\langle\eta_{s}, g\right\rangle\right) 1_{\left\{\tau\left(W_{s}^{*}\right)=\infty\right\}} \mathbb{P}_{\rho_{s}, W_{s}}^{*}\left[T_{D}=\infty\right]\right) \\
& =\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right)\right. \\
& \left.\quad \quad \times \exp \left(-\left\langle\eta_{s}, g\right\rangle\right) 1_{\left\{\tau\left(W_{s}^{*}\right)=\infty\right\}} \exp \left(-\int \rho_{s}(d t) u\left(W_{s}(t)\right)\right)\right) \tag{4.34}
\end{align*}
$$

Finally, we use Lemma 4.6 .1 to write
$\mathbb{N}_{x}\left(\int_{0}^{\sigma} d L_{s}^{a} u\left(\widehat{W}_{s}\right) F\left(W_{s}^{*}\right) 1_{\left\{\tau\left(W_{s}^{*}\right)=\infty\right\}} \exp \left(-\left\langle\eta_{s}, g\right\rangle\right) \exp \left(-\int \rho_{s}(d t) u\left(W_{s}(t)\right)\right)\right)$
$=e^{-\alpha a} \int \mathbb{M}_{a}(d \mu d \nu) \Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) e^{-<\nu, g>} \exp \left(-\int \mu(d r) u\left(\xi_{r}\right)\right)\right]$
$=e^{-\alpha a} \Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) \int \mathbb{M}_{a}(d \mu d \nu) e^{-<\nu, g>} \exp \left(-\int \mu(d r) u\left(\xi_{r}\right)\right)\right]$
$=\Pi_{x}\left[1_{\{a<\tau\}} u\left(\xi_{a}\right) F\left(\xi_{r}, 0 \leq r \leq a\right) \exp \left(-\int_{0}^{a} \gamma_{\psi}\left(u\left(\xi_{r}\right), g(r)\right) d r\right)\right]$.
Formula (4.27) follows by combining this equality with (4.30), (4.32), (4.33) and (4.34).

Formula (4.28) is the special case $g=0$ in (4.27). To prove the last assertion, let $\zeta\left(W_{T_{D}}, d \mu\right)$ be the law of the random measure

$$
\beta 1_{\left[0, H_{T_{D}}\right]}(r) d r+\sum_{i \in I}\left(v_{i}-\ell_{i}\right) \delta_{r_{i}}
$$

where $\sum \delta_{\left(r_{i}, v_{i}, \ell_{i}\right)}$ is a Poisson point measure on $\mathbb{R}_{+}^{3}$ with intensity given by formula (4.29). Then, for every $a>0$, we can use (4.28) to compute

$$
\begin{aligned}
& \mathbb{N}_{x}\left(F\left(W_{T_{D}}(r), 0 \leq r \leq a\right) 1_{\left\{T_{D}<\infty\right\}} 1_{\left\{a<H_{T_{D}}\right\}} \int \zeta\left(W_{T_{D}}, d \mu\right) e^{-\langle\mu, g\rangle}\right) \\
& =\Pi_{x}\left[1_{\{a<\tau\}} F\left(\xi_{r}, 0 \leq r \leq a\right) u\left(\xi_{a}\right) \exp \left(-\int_{0}^{a} \tilde{\psi}\left(u\left(\xi_{r}\right)\right) d r\right)\right. \\
& \left.\times \exp \left(-\beta \int_{0}^{a} d r g(r)-\int_{0}^{a} d r \int \pi(d v) \int_{0}^{v} d \ell e^{-\ell u\left(\xi_{r}\right)}\left(1-e^{-(v-\ell) g(r)}\right)\right)\right] \\
& =\Pi_{x}\left[1_{\{a<\tau\}} F\left(\xi_{r}, 0 \leq r \leq a\right) u\left(\xi_{a}\right) \exp \left(-\int_{0}^{a} \gamma_{\psi}\left(u\left(\xi_{r}\right), g(r)\right) d r\right)\right]
\end{aligned}
$$

using (4.26) in the last equality.
Set $\mathbb{N}_{x}^{D}=\mathbb{N}_{x}\left(\cdot \mid T_{D}<\infty\right)$ to simplify notation. By comparing with (4.27), we see that for any nonnegative measurable function $g$ with support in $[0, a]$, we have

$$
\mathbb{N}_{x}^{D}\left[e^{-\left\langle\rho_{T_{D}}, g\right\rangle} \mid W_{T_{D}}\right]=\int \zeta\left(W_{T_{D}}, d \mu\right) e^{-\langle\mu, g\rangle}
$$

a.s. on the set $\left\{H_{T_{D}}>a\right\}$. This is enough to conclude that $\zeta\left(W_{T_{D}}, d \mu\right)$ is the conditional distribution of $\rho_{T_{D}}$ knowing $W_{T_{D}}$, provided that we already know that $\rho_{T_{D}}\left(\left\{H_{T_{D}}\right\}\right)=0$ a.s. The latter fact however is a simple consequence of (4.2). This completes the proof of the theorem.

The case of Brownian motion. Suppose that the spatial motion $\xi$ is $d$ dimensional Brownian motion and that $D$ is a domain in $\mathbb{R}^{d}$. Then, it is easy to see that the function $u(x)=\mathbb{N}_{x}\left(T_{D}<\infty\right), x \in D$ is of class $C^{2}$ and solves $\frac{1}{2} \Delta u=\psi(u)$. In the context of superprocesses, this was observed by Dynkin [12]. We may argue as follows. First note that the set of nonnegative solutions of $\frac{1}{2} \Delta u=\psi(u)$ in a domain is closed under pointwise convergence (for a probabilistic proof, reproduce the arguments of the proof of Proposition 9 (iii) in [31]). Then let $\left(D_{n}\right)$ be a sequence of bounded regular subdomains of $D$, such that $D_{n} \subset D_{n+1}$ and $D=\lim \uparrow D_{n}$. For every $n \geq 0$, set

$$
v_{n}(x)=\mathbb{N}_{x}\left(\mathcal{Z}^{D_{n}} \neq 0\right), \quad u_{n}(x)=\mathbb{N}_{x}\left(T_{D_{n}}<\infty\right), \quad x \in D_{n}
$$

From the properties of the exit measure, it is immediate to see that $v_{n} \leq u_{n}$. On the other hand, by writing

$$
v_{n}(x)=\lim _{\lambda \uparrow \infty} \uparrow \mathbb{N}_{x}\left(1-\exp -\lambda\left\langle\mathcal{Z}^{D_{n}}, 1\right\rangle\right),
$$

we deduce from Proposition 4.5.1 and the stability of the set of nonnegative solutions under pointwise convergence that $v_{n}$ is of class $C^{2}$ and solves $\frac{1}{2} \Delta v_{n}=\psi\left(v_{n}\right)$ in $D$. Since the function $x \longrightarrow \mathbb{N}_{x}\left(1-\exp -\lambda\left\langle\mathcal{Z}^{D_{n}}, 1\right\rangle\right)$ has boundary value $\lambda$ on $\partial D_{n}$ (Proposition 4.5.1), we also see that $v_{n}$ has boundary value $+\infty$ on $\partial D_{n}$.

Then, it follows from Lemma 4.5.3 and our assumption $\bar{D}_{n} \subset D_{n+1}$ that $v_{n}(x) \geq$ $u_{n+1}(x)$ for $x \in D_{n}$. Since it is easy to see that $u_{n}(x)$ decreases to $u(x)$ as $n \rightarrow \infty$, for every $x \in D$, we conclude from the inequalities $u_{n+1}(x) \leq v_{n}(x) \leq u_{n}(x)$ that $v_{n}(x)$ also converges to $u(x)$ pointwise as $n \rightarrow \infty$. Hence $u$ is a nonnegative solution of $\frac{1}{2} \Delta u=\psi(u)$ in $D$. The preceding argument gives more. Let $v$ be any nonnegative solution of $\Delta v=\frac{1}{2} \psi(v)$ in $D$. Since $v_{n \mid \partial D_{n}}=+\infty$, the comparison principle (Lemma V. 7 in [31]) implies that $v \leq v_{n}$ in $D_{n}$. By passing to the limit $n \rightarrow \infty$, we conclude that $v \leq u$. Hence $u$ is the maximal nonnegative solution of $\frac{1}{2} \Delta u=\psi(u)$ in $D$.

Suppose that $u(x)>0$ for some $x \in D$. It is easy to see that this implies $u(y)>0$ for every $y \in D$ (use a suitable Harnack principle or a probabilistic argument relying on the fact that $u\left(\xi_{t}\right) \exp \left(-\int_{0}^{t} \tilde{\psi}\left(u\left(\xi_{r}\right)\right) d r\right)$ is a martingale). By applying Itô's formula to $\log u\left(\xi_{t}\right)$, we see that $\Pi_{x}$ a.s. on $\{t<\tau\}$,

$$
\begin{aligned}
\log u\left(\xi_{t}\right) & =\log u(x)+\int_{0}^{t} \frac{\nabla u}{u}\left(\xi_{r}\right) \cdot d \xi_{r}+\frac{1}{2} \int_{0}^{t} \Delta(\log u)\left(\xi_{r}\right) d r \\
& =\log u(x)+\int_{0}^{t} \frac{\nabla u}{u}\left(\xi_{r}\right) \cdot d \xi_{r}+\int_{0}^{t}\left(\tilde{\psi}\left(u\left(\xi_{r}\right)\right)-\frac{1}{2}\left|\frac{\nabla u}{u}\right|^{2}\left(\xi_{r}\right)\right) d r
\end{aligned}
$$

We can then rewrite (4.28) in the form

$$
\begin{aligned}
& \mathbb{N}_{x}^{D}\left(1_{\left\{T_{D}<\infty\right\}} 1_{\left\{t<H_{T_{D}}\right\}} F\left(W_{T_{D}}(r), 0 \leq r \leq t\right)\right) \\
& \quad=\Pi_{x}\left[1_{\{t<\tau\}} \exp \left(\int_{0}^{t} \frac{\nabla u}{u}\left(\xi_{r}\right) \cdot d \xi_{r}-\frac{1}{2} \int_{0}^{t}\left|\frac{\nabla u}{u}\right|^{2}\left(\xi_{r}\right) d r\right) F\left(\xi_{r}, 0 \leq r \leq t\right)\right]
\end{aligned}
$$

An application of Girsanov's theorem then shows that $W_{T_{D}}$ is distributed as the solution of the stochastic differential equation

$$
\begin{aligned}
& d x_{t}=d B_{t}+\frac{\nabla u}{u}\left(x_{t}\right) d t \\
& x_{0}=x
\end{aligned}
$$

(where $B$ is a standard $d$-dimensional Brownian motion) which can be defined up to its first hitting time of $\partial D$. See [29] for a discussion and another interpretation of this distribution on paths in the case $\psi(u)=u^{2}$.

### 4.7 The reduced tree in an open set

We keep the notation and assumptions of the previous section. In particular, we assume that $W^{*}$ has continuous sample paths with respect to the distance $d^{*}, D$ is an open set in $E, x \in D, T_{D}=\inf \left\{s>0: \tau\left(W_{s}^{*}\right)<\infty\right\}$ and $u(x)=\mathbb{N}_{x}\left(T_{D}<\infty\right)<\infty$. To avoid trivialities, we assume that $u(x)>0$, and we recall the notation $\mathbb{N}_{x}^{D}=\mathbb{N}_{x}(\cdot \mid$ $\left.T_{D}<\infty\right)$. We will assume in addition that

$$
\begin{equation*}
\sup _{y \in K} u(y)<\infty \tag{4.35}
\end{equation*}
$$

for every compact subset $K$ of $D$. This assumption holds in particular when $\xi$ is Brownian motion in $\mathbb{R}^{d}$, under the condition (4.15) (use translation invariance and the fact that $u(0)<\infty$ when $D$ is an open ball centered at the origin).

We also set:

$$
L_{D}=\sup \left\{s \geq 0: \tau\left(W_{s}^{*}\right)<\infty\right\}
$$

and

$$
m_{D}=\inf _{T_{D} \leq s \leq L_{D}} H_{s}
$$

As a consequence of the first lemma below, we will see that $m_{D}<H_{T_{D}}, \mathbb{N}_{x}^{D}$ a.s.
Our goal is to describe the genealogical structure of the paths $W_{s}$ that exit $D$, up to their first exit time from $D$, under the probability measure $\mathbb{N}_{x}^{D}$. To be more precise, all paths $W_{s}$ such that $\tau\left(W_{s}^{*}\right)<\infty$ must coincide up to level $m_{D}$. At level $m_{D}$ there is a branching point with finitely many branches, each corresponding to an excursion of $H$ above level $m_{D}$ during which $W$ hits $D^{c}$. In each such excursion, the paths $W_{s}$ that hit $D^{c}$ will be the same up to a level (strictly greater than $m_{D}$ ) at which there is another branching point, and so on.

We will describe this genealogical structure in a recursive way. We will first derive the law of the common part to the paths $W_{s}$ that do exit $D$. This common part is represented by a stopped path $W_{0}^{D}$ in $\mathcal{W}_{x}$ with lifetime $\zeta_{W_{0}^{D}}=m_{D}$. Then we will obtain the distribution of the "number of branches" at level $m_{D}$, that is the number of excursions of $W$ above height $m_{D}$ that hit $D^{c}$. Finally, we will see that conditionally on $W_{0}^{D}$, these excursions are independent and distributed according to $\mathbb{N}_{\hat{W}_{0}^{D}}\left(\cdot \mid T_{D}<\infty\right)$. This completes our recursive description since we can apply to each of these excursions the results obtained under $\mathbb{N}_{x}^{D}$.

Before coming to the main result of this section, we state an important lemma.
Lemma 4.7.1 The point $T_{D}$ is not isolated in $\left\{s \geq 0: \tau\left(W_{s}^{*}\right)<\infty\right\}$, $\mathbb{N}_{x}$ a.e. on $\left\{T_{D}<\infty\right\}$.

Proof. We start with some preliminary observations. Let $(\mu, \mathrm{w}) \in \Theta_{x}$ be such that $\mu(\{H(\mu)\})=0$ and $\mathrm{w}(t) \in D$ for every $t \in[0, H(\mu))$. As an application of Lemma 4.2.4, we have

$$
\begin{aligned}
\mathbb{P}_{\mu, \mathrm{w}}^{*}\left[T_{D}<\infty\right] & =1-\exp -\int_{[0, H(\mu))} \mathbb{N}_{\mathrm{w}(t)}\left(T_{D}<\infty\right) \mu(d t) \\
& =1-\exp -\int_{[0, H(\mu))} u(\mathrm{w}(t)) \mu(d t)
\end{aligned}
$$

By the previous formula, the equality $\mathbb{P}_{\mu, \mathrm{w}}^{*}\left[T_{D}=0\right]=1$ can only hold if

$$
\begin{equation*}
\int_{[0, H(\mu))} u(\mathrm{w}(t)) \mu(d t)=\infty . \tag{4.36}
\end{equation*}
$$

Conversely, condition (4.36) also implies that $\mathbb{P}_{\mu, \mathrm{w}}^{*}\left[T_{D}=0\right]=1$. To see this, first note that our assumption (4.35) guarantees that for every $\varepsilon>0$,

$$
\int_{[0, H(\mu)-\varepsilon]} u(\mathrm{w}(t)) \mu(d t)<\infty
$$

and thus we have also under (4.36)

$$
\int_{(H(\mu)-\varepsilon, H(\mu))} u(\mathrm{w}(t)) \mu(d t)=\infty .
$$

Then write $\mu_{\varepsilon}$ for the restriction of $\mu$ to $[0, H(\mu)-\varepsilon]$, and set

$$
S_{\varepsilon}=\inf \left\{s \geq 0:\left\langle\rho_{s}, 1\right\rangle=\left\langle\mu_{\varepsilon}, 1\right\rangle\right\} .
$$

Lemma 4.2.4 again implies that

$$
\mathbb{P}_{\mu, \mathrm{w}}^{*}\left[T_{D} \leq S_{\varepsilon}\right]=1-\exp -\int_{(H(\mu)-\varepsilon, H(\mu))} u(\mathrm{w}(t)) \mu(d t)=1 .
$$

Since $S_{\varepsilon} \downarrow 0$ as $\varepsilon \downarrow 0, \mathbb{P}_{\mu, \mathrm{w}}^{*}$ a.s., we get that $\mathbb{P}_{\mu, \mathrm{w}}^{*}\left[T_{D}=0\right]=1$, which was the desired result.

Let us prove the statement of the lemma. Thanks to the strong Markov property, it is enough to prove that $\mathbb{P}_{\rho_{T_{D}}, W_{T_{D}}}^{*}\left[T_{D}=0\right]=1, \mathbb{N}_{x}$ a.e. on $\left\{T_{D}<\infty\right\}$. Note that we have $\rho_{T_{D}}\left(\left\{H_{T_{D}}\right\}\right)=0$ and $W_{T_{D}}(t) \in D$ for every $t<H_{T_{D}}, \mathbb{N}_{x}$ a.e. on $\left\{T_{D}<\infty\right\}$. By the preceding observations, it is enough to prove that

$$
\begin{equation*}
\int_{\left[0, H_{T_{D}}\right)} u\left(W_{T_{D}}(t)\right) \rho_{T_{D}}(d t)=\infty, \quad \text { a.e. on }\left\{T_{D}<\infty\right\} . \tag{4.37}
\end{equation*}
$$

To this end, set for every $s>0$,

$$
M_{s}=\mathbb{N}_{x}\left(T_{D}<\infty \mid \mathcal{F}_{s}\right)
$$

The Markov property at time $s$ shows that we have for every $s>0, \mathbb{N}_{x}$ a.e.,

$$
\begin{aligned}
M_{s} & =1_{\left\{T_{D} \leq s\right\}}+1_{\left\{s<T_{D}\right\}} \mathbb{P}_{\rho_{s}, W_{s}}^{*}\left[T_{D}<\infty\right] \\
& =1_{\left\{T_{D} \leq s\right\}}+1_{\left\{s<T_{D}\right\}}\left(1-\exp -\int u\left(W_{s}(t)\right) \rho_{s}(d t)\right)
\end{aligned}
$$

Since the process $\left(\rho_{s}\right)$ is right-continuous for the variation distance on measures, it is easy to verify that the process $1_{\left\{s<T_{D}\right\}}\left(1-\exp -\int u\left(W_{s}(t)\right) \rho_{s}(d t)\right)$ is rightcontinuous. Because $\left(M_{s}, s>0\right)$ is a martingale with respect to the filtration $\left(\mathcal{F}_{s}\right)$, a standard result implies that this process also has left limits at every $s>0, \mathbb{N}_{x}$ a.e. In particular the left limit at $T_{D}$

$$
\lim _{s \uparrow T_{D}, s<T_{D}} M_{s}=\lim _{s \uparrow T_{D}, s<T_{D}}\left(1-\exp -\int u\left(W_{s}(t)\right) \rho_{s}(d t)\right)
$$

exists $\mathbb{N}_{x}$ a.e. on $\left\{T_{D}<\infty\right\}$. It is not hard to verify that this limit is equal to 1: If $D_{n}=\left\{y \in D: \operatorname{dist}\left(y, D^{c}\right)>n^{-1}\right\}$ and $T_{n}=T_{D_{n}}$, we have $T_{n}<T_{D}$ and $T_{n} \uparrow T_{D}$ on $\left\{T_{D}<\infty\right\}$, and $M_{T_{n}}=\mathbb{N}_{x}\left(T_{D}<\infty \mid \mathcal{F}_{T_{n}}\right)$ converges to 1 as $n \rightarrow \infty$ on the set $\left\{T_{D}<\infty\right\}$ because $T_{D}$ is measurable with respect to the $\sigma$-field $\bigvee \mathcal{F}_{T_{n}}$.

Summarizing, we have proved that

$$
\begin{equation*}
\lim _{s \uparrow T_{D}, s<T_{D}} \int u\left(W_{s}(t)\right) \rho_{s}(d t)=+\infty \tag{4.38}
\end{equation*}
$$

$\mathbb{N}_{x}$ a.e. on $\left\{T_{D}<\infty\right\}$. Then, for every rational $a>0$, consider on the event $\left\{T_{D}<\infty\right\} \cap\left\{H_{T_{D}}>a\right\}$, the number $\alpha_{(a)}$ defined as the left end of the excursion interval of $H$ above $a$ that straddles $T_{D}$. As a consequence of the considerations in subsection 4.2.2, the following two facts hold on $\left\{T_{D}<\infty\right\} \cap\left\{H_{T_{D}}>a\right\}$ :

$$
\begin{aligned}
& \rho_{\alpha_{(a)}} \text { is the restriction of } \rho_{T_{D}} \text { to }[0, a) \\
& W_{\alpha_{(a)}}(t)=W_{T_{D}}(t), \quad \text { for every } t \in[0, a) .
\end{aligned}
$$

Thus, we have also on the same event

$$
\int u\left(W_{\alpha_{(a)}}(t)\right) \rho_{\alpha_{(a)}}(d t)=\int_{[0, a)} u\left(W_{T_{D}}(t)\right) \rho_{T_{D}}(d t) .
$$

Now on the event $\left\{T_{D}<\infty\right\}$ we can pick a sequence $\left(a_{n}\right)$ of rationals strictly increasing to $H_{T_{D}}$. We observe that $\alpha_{\left(a_{n}\right)}$ also converges to $T_{D}$ (if $S$ is the increasing limit of $\alpha_{\left(a_{n}\right)}$, the snake property implies that $\widehat{W}_{S}=\widehat{W}_{T_{D}} \in D^{c}$ and so we have $S \geq T_{D}$, whereas the other inequality is trivial). Therefore, using (4.38),

$$
\infty=\lim _{n \rightarrow \infty} \int u\left(W_{\alpha_{\left(a_{n}\right)}}(t)\right) \rho_{\alpha_{\left(a_{n}\right)}}(d t)=\lim _{n \rightarrow \infty} \int_{\left[0, a_{n}\right)} u\left(W_{T_{D}}(t)\right) \rho_{T_{D}}(d t),
$$

which yields (4.37).
Lemma 4.7.1 implies that $T_{D}<L_{D}, \mathbb{N}_{x}^{D}$ a.s. Since we know that $\rho_{T_{D}}\left(\left\{H_{T_{D}}\right\}\right)=0$, $\mathbb{N}_{x}^{D}$ a.s., an application of the strong Markov property at time $T_{D}$ shows that $m_{D}<$ $H_{T_{D}}, \mathbb{N}_{x}^{D}$ a.s. We define $W_{0}^{D}$ as the stopped path which is the restriction of $W_{T_{D}}$ to $\left[0, m_{D}\right]$. Then we define the excursions of $W$ above level $m_{D}$ in a way analogous to subsection 4.2.2. If

$$
R_{D}=\sup \left\{s \leq T_{D}: H_{s}=m_{D}\right\}, \quad S_{D}=\inf \left\{s \geq L_{D}: H_{s}=m_{D}\right\}
$$

we let $\left(a_{j}, b_{j}\right), j \in J$ be the connected components of the open set $\left(R_{D}, S_{D}\right) \cap\{s \geq$ $\left.0: H_{s}>m_{D}\right\}$. For each $j \in J$, we can then define the process $W^{(j)} \in C\left(\mathbb{R}_{+}, \mathcal{W}\right)$ by setting

$$
\begin{array}{ll}
W_{s}^{(j)}(r)=W_{a_{j}+s}\left(m_{D}+r\right), \zeta_{W_{s}^{(j)}}=H_{a_{j}+s}-m_{D} & \text { if } 0<s<b_{j}-a_{j} \\
W_{s}^{(j)}=\widehat{W}_{0}^{D} & \text { if } s=0 \text { or } s \geq b_{j}-a_{j}
\end{array}
$$

By a simple continuity argument, the set $\left\{j \in J: T_{D}\left(W^{(j)}\right)<\infty\right\}$ is finite a.s., and we set

$$
N_{D}=\operatorname{Card}\left\{j \in J: T_{D}\left(W^{(j)}\right)<\infty\right\} .
$$

We write $W^{D, 1}, W^{D, 2}, \ldots, W^{D, N_{D}}$ for the excursions $W^{(j)}$ such that $T_{D}\left(W^{(j)}\right)<\infty$, listed in chronological order.

We are now ready to state our main result.
Theorem 4.7.2 For every $r \geq 0$, set $\theta(r)=\psi^{\prime}(r)-\tilde{\psi}(r)$. Then the law of $W_{0}^{D}$ is characterized by the following formula, valid for any nonnegative measurable function $F$ on $\mathcal{W}^{*}$ :

$$
\begin{align*}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} F\left(W_{0}^{D}\right)\right) \\
& =\int_{0}^{\infty} d b \Pi_{x}\left[1_{\{b<\tau\}} u\left(\xi_{b}\right) \theta\left(u\left(\xi_{b}\right)\right) \exp \left(-\int_{0}^{b} \psi^{\prime}\left(u\left(\xi_{r}\right)\right) d r\right) F\left(\xi_{r}, 0 \leq r \leq b\right)\right] \tag{4.39}
\end{align*}
$$

The conditional distribution of $N_{D}$ knowing $W_{0}^{D}$ is given by:

$$
\begin{equation*}
\mathbb{N}_{x}^{D}\left[r^{N_{D}} \mid W_{0}^{D}\right]=r \frac{\psi^{\prime}(U)-\gamma_{\psi}(U,(1-r) U)}{\psi^{\prime}(U)-\gamma_{\psi}(U, 0)}, \quad 0 \leq r \leq 1 \tag{4.40}
\end{equation*}
$$

where $U=u\left(\widehat{W}_{0}^{D}\right)$. Finally, conditionally on the pair $\left(W_{0}^{D}, N_{D}\right)$, the processes $W^{D, 1}, W^{D, 2}, \ldots, W^{D, N_{D}}$ are independent and distributed according to $\mathbb{N}_{\hat{W}_{0}^{D}}^{D}$.

Proof. Our first objective is to compute the conditional distribution of $m_{D}$ knowing $W_{T_{D}}$. To this end, we will apply the strong Markov property of the Lévy snake at time $T_{D}$. We have for every $b>0$

$$
\mathbb{N}_{x}^{D}\left[m_{D}>b \mid \rho_{T_{D}}, W_{T_{D}}\right]=\mathbb{P}_{\rho_{T_{D}}, W_{T_{D}}}^{*}\left[\inf _{0 \leq s \leq L_{D}} H_{s}>b\right]
$$

By Lemma 4.2.4, the latter expression is equal to the probability that in a Poisson point measure with intensity

$$
\rho_{T_{D}}(d h) \mathbb{N}_{W_{T_{D}}}(h)(d \rho d W)
$$

there is no atom $\left(h_{i}, \rho^{i}, W^{i}\right)$ such that $h_{i} \leq b$ and $T_{D}\left(W^{i}\right)<\infty$. We conclude that

$$
\begin{align*}
\mathbb{N}_{x}^{D}\left[m_{D}>b \mid \rho_{T_{D}}, W_{T_{D}}\right] & =\exp -\int_{[0, b]} \rho_{T_{D}}(d h) \mathbb{N}_{W_{T_{D}}(h)}\left(T_{D}<\infty\right) \\
& =\exp -\int_{[0, b]} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right) \tag{4.41}
\end{align*}
$$

Recall that the conditional law of $\rho_{T_{D}}$ knowing $W_{T_{D}}$ is given in Theorem 4.6.2. Using this conditional distribution we see that
$\mathbb{N}_{x}^{D}\left[m_{D}>b \mid W_{T_{D}}\right]=\exp \left(-\beta \int_{0}^{b} d a u\left(W_{T_{D}}(a)\right)\right) E\left[\exp -\sum_{r_{i} \leq b}\left(v_{i}-\ell_{i}\right) u\left(W_{T_{D}}\left(r_{i}\right)\right)\right]$, where $\sum \delta_{\left(r_{i}, v_{i}, \ell_{i}\right)}$ is a Poisson point measure with intensity (depending on $W_{T_{D}}$ ) given by (4.29). By exponential formulas for Poisson measures, we have

$$
\begin{aligned}
& E\left[\exp -\sum_{r_{i} \leq b}\left(v_{i}-\ell_{i}\right) u\left(W_{T_{D}}\left(r_{i}\right)\right)\right] \\
& \quad=\exp -\int_{0}^{b} d r \int \pi(d v) \int d \ell e^{-\ell u\left(W_{T_{D}}(r)\right)}\left(1-e^{-(v-\ell) u\left(W_{T_{D}}(r)\right)}\right)
\end{aligned}
$$

By substituting this in the previous displayed formula, and using (4.26), we get

$$
\begin{equation*}
\mathbb{N}_{x}^{D}\left[m_{D}>b \mid W_{T_{D}}\right]=\exp \left(-\int_{0}^{b} d r\left(\psi^{\prime}\left(u\left(W_{T_{D}}(r)\right)\right)-\tilde{\psi}\left(u\left(W_{T_{D}}(r)\right)\right)\right)\right) \tag{4.42}
\end{equation*}
$$

Hence, if $\theta(r)=\psi^{\prime}(r)-\tilde{\psi}(r)$ as in the statement of the theorem, the conditional law of $m_{D}$ knowing $W_{T_{D}}$ has density

$$
1_{\left[0, H_{T_{D}}\right)}(b) \theta\left(u\left(W_{T_{D}}(r)\right)\right) \exp \left(-\int_{0}^{b} \theta\left(u\left(W_{T_{D}}(r)\right)\right) d r\right)
$$

It follows that

$$
\begin{aligned}
& \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} F\left(W_{0}^{D}\right)\right) \\
& =\mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} F\left(W_{T_{D}}(t), 0 \leq t \leq m_{D}\right)\right) \\
& =\mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} \int_{0}^{H_{T_{D}}} d b \theta\left(u\left(W_{T_{D}}(b)\right)\right) \exp \left(-\int_{0}^{b} \theta\left(u\left(W_{T_{D}}(r)\right)\right) d r\right)\right. \\
& \left.\quad \times F\left(W_{T_{D}}(t), 0 \leq t \leq b\right)\right)
\end{aligned}
$$

$$
\begin{gathered}
=\int_{0}^{\infty} d b \mathbb{N}_{x}\left(1_{\left\{T_{D}<\infty\right\}} 1_{\left\{b<H_{T_{D}}\right\}} \theta\left(u\left(W_{T_{D}}(b)\right)\right) \exp \left(-\int_{0}^{b} \theta\left(u\left(W_{T_{D}}(r)\right)\right) d r\right)\right. \\
\left.\times F\left(W_{T_{D}}(t), 0 \leq t \leq b\right)\right) \\
=\int_{0}^{\infty} d b \Pi_{x}\left[1_{\{b<\tau\}} u\left(\xi_{b}\right) \theta\left(u\left(\xi_{b}\right)\right) \exp \left(-\int_{0}^{b} \psi^{\prime}\left(u\left(\xi_{r}\right)\right) d r\right) F\left(\xi_{r}, 0 \leq r \leq b\right)\right],
\end{gathered}
$$

using (4.28) in the last equality. This gives the first assertion of the theorem.
We now turn to the distribution of $N_{D}$. We use again the strong Markov property at time $T_{D}$ and Lemma 4.2.4 to analyse the conditional distribution of the pair $\left(m_{D}, N_{D}\right)$ knowing $\left(\rho_{T_{D}}, W_{T_{D}}\right)$. Conditional on $\left(\rho_{T_{D}}, W_{T_{D}}\right)$, let $\sum \delta_{\left(h_{i}, \rho^{i}, W^{i}\right)}$ be a Poisson point measure with intensity

$$
\rho_{T_{D}}(d h) \mathbb{N}_{W_{T_{D}}}(h)(d \rho d W)
$$

Set

$$
\begin{aligned}
& m=\inf \left\{h_{i}: T_{D}\left(W^{i}\right)<\infty\right\} \\
& M=\operatorname{Card}\left\{i: h_{i}=m \text { and } T_{D}\left(W^{i}\right)<\infty\right\}
\end{aligned}
$$

Then Lemma 4.2.4 and the strong Markov property show that the pairs $(m, 1+M)$ and $\left(m_{D}, N_{D}\right)$ have the same distribution conditional on $\left(\rho_{T_{D}}, W_{T_{D}}\right)$. Recall that the conditional distribution of $m_{D}$ (or of $m$ ) is given by (4.41).

Now note that:

- If $\rho_{T_{D}}(\{m\})=0$, then $M=1$ because the Poisson measure $\sum \delta_{\left(h_{i}, \rho^{i}, W^{i}\right)}$ cannot have two atoms at a level $h$ such that $\rho_{T_{D}}(\{h\})=0$.
- Let $b \geq 0$ be such that $\rho_{T_{D}}(\{b\})>0$. The event $\{m=b\}$ occurs with probability

$$
\exp \left(-\int_{[0, b)} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right)\right)\left(1-e^{-\rho_{T_{D}}(\{b\}) u\left(W_{T_{D}}(b)\right)}\right)
$$

Conditionally on this event, $M$ is distributed as a Poisson variable with parameter $c=\rho_{T_{D}}(\{b\}) u\left(W_{T_{D}}(b)\right)$ and conditioned to be (strictly) positive, whose generating function is

$$
\frac{e^{-c(1-r)}-e^{-c}}{1-e^{-c}}
$$

Since the continuous part of the law of $m$ has density

$$
\beta u\left(W_{T_{D}}(b)\right) \exp \left(-\int_{[0, b)} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right)\right)
$$

we get by combining the previous two cases that

$$
\begin{align*}
& \mathbb{N}_{x}^{D}\left[f\left(m_{D}\right) r^{N_{D}} \mid \rho_{T_{D}}, W_{T_{D}}\right] \\
& \quad=\beta r^{2} \int_{0}^{H_{T_{D}}} d b f(b) u\left(W_{T_{D}}(b)\right) \exp \left(-\int_{[0, b)} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right)\right) \\
& \quad+r \sum_{\rho_{T_{D}}(\{b\})>0} f(b) \exp \left(-\int_{[0, b)} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right)\right) \\
& \quad \times\left(e^{-\rho_{T_{D}}(\{b\}) u\left(W_{T_{D}}(b)\right)(1-r)}-e^{-\rho_{T_{D}}(\{b\}) u\left(W_{T_{D}}(b)\right)}\right) . \tag{4.43}
\end{align*}
$$

We now need to integrate the right side of (4.43) with respect to the conditional law of $\rho_{T_{D}}$ knowing $W_{T_{D}}$. We get

$$
\mathbb{N}_{x}^{D}\left[f\left(m_{D}\right) r^{N_{D}} \mid W_{T_{D}}\right]=A_{1}+A_{2}
$$

where

$$
\begin{align*}
A_{1} & =\beta r^{2} \mathbb{N}_{x}^{D}\left[\int_{0}^{H_{T_{D}}} d b f(b) u\left(W_{T_{D}}(b)\right) \exp \left(-\int_{[0, b)} \rho_{T_{D}}(d h) u\left(W_{T_{D}}(h)\right)\right) \mid W_{T_{D}}\right] \\
& =\beta r^{2} \int_{0}^{H_{T_{D}}} d b f(b) u\left(W_{T_{D}}(b)\right) \exp \left(-\int_{[0, b)} \theta\left(u\left(W_{T_{D}}(h)\right)\right) d h\right) \tag{4.44}
\end{align*}
$$

by the calculation used in the proof of (4.42). We then compute $A_{2}$. To this end, let $\mathcal{N}(d b d v d \ell)$ be (conditionally on $W_{T_{D}}$ ) a Poisson point measure in $\mathbb{R}_{+}^{3}$ with intensity

$$
1_{\left[0, H_{T_{D}}\right]}(b) 1_{[0, v]}(\ell) e^{-\ell u\left(W_{T_{D}}(b)\right)} d b \pi(d v) d \ell
$$

From Theorem 4.6.2, we get

$$
\begin{aligned}
A_{2}=r \mathbb{N}_{x}^{D}\left[\int \mathcal{N}( \right. & d b d v d \ell) f(b) e^{-\beta \int_{0}^{b} d a u\left(W_{T_{D}}(a)\right)-\int_{\{a<b\}} \mathcal{N}\left(d a d v^{\prime} d \ell^{\prime}\right)\left(v^{\prime}-\ell^{\prime}\right) u\left(W_{T_{D}}(a)\right)} \\
& \left.\times\left(e^{-(v-\ell) u\left(W_{T_{D}}(b)\right)(1-r)}-e^{-(v-\ell) u\left(W_{T_{D}}(b)\right)}\right) \mid W_{T_{D}}\right]
\end{aligned}
$$

From Lemma 3.1.5 and (once again) the calculation used in proving (4.42), we arrive at

$$
\begin{aligned}
A_{2}= & r \int_{0}^{H_{T_{D}}} d b f(b) \exp \left(-\int_{[0, b)} \theta\left(u\left(W_{T_{D}}(a)\right)\right) d a\right) \\
& \times \int \pi(d v) \int_{0}^{v} d \ell e^{-\ell u\left(W_{T_{D}}(b)\right)}\left(e^{-(v-\ell)(1-r) u\left(W_{T_{D}}(b)\right)}-e^{-(v-\ell) u\left(W_{T_{D}}(b)\right)}\right) .
\end{aligned}
$$

From (4.26), we have

$$
\begin{aligned}
& \int \pi(d v) \int_{0}^{v} d \ell e^{-\ell u\left(W_{T_{D}}(b)\right)}\left(e^{-(v-\ell)(1-r) u\left(W_{T_{D}}(b)\right)}-e^{-(v-\ell) u\left(W_{T_{D}}(b)\right)}\right) \\
& \left.\quad=\psi^{\prime}\left(u\left(W_{T_{D}}(b)\right)\right)-\gamma_{\psi}\left(u\left(W_{T_{D}}(b)\right),(1-r) u\left(W_{T_{D}}(b)\right)\right)-\beta r u\left(W_{T_{D}}(b)\right)\right) .
\end{aligned}
$$

By substituting this identity in the previous formula for $A_{2}$, and then adding the formula for $A_{1}$, we arrive at:

$$
\begin{aligned}
& \mathbb{N}_{x}^{D}\left[f\left(m_{D}\right) r^{N_{D}} \mid W_{T_{D}}\right] \\
& \begin{aligned}
&=r \int_{0}^{H_{T_{D}}} d b f(b) \exp \left(-\int_{0}^{b} d a \theta\left(u\left(W_{T_{D}}(a)\right)\right)\right) \\
& \times\left(\psi^{\prime}\left(u\left(W_{T_{D}}(b)\right)\right)-\gamma_{\psi}\left(u\left(W_{T_{D}}(b)\right),(1-r) u\left(W_{T_{D}}(b)\right)\right)\right) \\
&= \mathbb{N}_{x}^{D}\left[\left.f\left(m_{D}\right) r \frac{\psi^{\prime}\left(u\left(W_{T_{D}}\left(m_{D}\right)\right)\right)-\gamma_{\psi}\left(u\left(W_{T_{D}}\left(m_{D}\right)\right),(1-r) u\left(W_{T_{D}}\left(m_{D}\right)\right)\right)}{\psi^{\prime}\left(u\left(W_{T_{D}}\left(m_{D}\right)\right)\right)-\gamma_{\psi}\left(u\left(W_{T_{D}}\left(m_{D}\right)\right), 0\right)} \right\rvert\, W_{T_{D}}\right] .
\end{aligned}
\end{aligned}
$$

In the last equality we used the conditional distribution of $m_{D}$ knowing $W_{T_{D}}$, and the fact that $\theta(u)=\psi^{\prime}(u)-\tilde{\psi}(u)=\psi^{\prime}(u)-\gamma_{\psi}(u, 0)$.

Finally, if $U=u\left(W_{T_{D}}\left(m_{D}\right)\right)=u\left(\widehat{W}_{0}^{D}\right)$, we have obtained

$$
\mathbb{N}_{x}^{D}\left[r^{N_{D}} \mid W_{0}^{D}\right]=r \frac{\psi^{\prime}(U)-\gamma_{\psi}(U,(1-r) U)}{\psi^{\prime}(U)-\gamma_{\psi}(U, 0)}
$$

which is formula (4.40) of the theorem.
It remains to obtain the last assertion of the theorem. Here again, we will rely on Lemma 4.2.4 and the strong Markov property at time $T_{D}$. We need to restate the result of Lemma 4.2 .4 in a slightly different form. Let $(\mu, \mathrm{w}) \in \Theta_{x}$ with $\mu(\{H(\mu)\}=0$ and $\mathrm{w}(t) \in D$ for every $t<H(\mu)$. Under $\mathbb{P}_{\mu, \mathrm{w}}^{*}$, we write $Y_{t}=\left\langle\rho_{t}, 1\right\rangle, K_{t}=\inf _{r \leq t} Y_{r}$ and $I_{t}=K_{t}-\langle\mu, 1\rangle$. If $\left(\alpha_{i}, \beta_{i}\right), i \in I$ are the excursion intervals of $Y-K$ away from 0 , we introduce the "excursions" $\left(\rho_{i}, W^{i}\right), i \in I$ as defined before the statement of Lemma 4.2.4. The starting height of excursion $\left(\rho_{i}, W^{i}\right)$ is $h_{i}=H_{\alpha_{i}}=H\left(k_{-I_{\alpha_{i}}} \mu\right)$. The proof of Lemma 4.2 .4 shows that the point measure

$$
\sum_{i \in I} \delta_{\left(-I_{\alpha_{i}}, \rho^{i}, W^{i}\right)}
$$

is Poisson with intensity $1_{[0,<\mu, 1>]}(u) d u \mathbb{N}_{\mathrm{w}\left(H\left(k_{u} \mu\right)\right)}(d \rho d W)$ (this is slightly more precise than the statement of Lemma 4.2.4).

We then write $i_{1}, i_{2}, \ldots$ for the indices $i \in I$ such that $T_{D}\left(W^{i}\right)<\infty$, ranked in such a way that $I_{\alpha_{i_{1}}}<I_{\alpha_{i_{2}}}<\cdots$. Our assumption (4.35) guarantees that this ordering is possible, and we have clearly $h_{i_{1}} \leq h_{i_{2}} \leq \cdots$. By well-known properties of Poisson measures, the processes $W^{i_{1}}, W^{i_{2}}, \ldots$ are independent conditionally on the sequence $h_{i_{1}}, h_{i_{2}}, \ldots$, and the conditional distribution of $W^{i_{\ell}}$ is $\mathbb{N}_{\mathrm{w}\left(h_{i_{\ell}}\right)}^{D}$.

If we apply the previous considerations to the shifted process $\left(\rho_{T_{D}+s}, W_{T_{D}+s} ; s \geq\right.$ $0)$, taking $\mu=\rho_{T_{D}}$ and $\mathrm{w}=W_{T_{D}}$ and relying on the strong Markov property at $T_{D}$ we can easily identify

$$
\begin{aligned}
& m_{D}=h_{i_{1}} \\
& N_{D}=1+\sup \left\{k \geq 1: h_{i_{k}}=h_{i_{1}}\right\} \\
& W^{D, N_{D}}=W^{i_{1}}, W^{D, N_{D}-1}=W^{i_{2}}, \ldots, W^{D, 2}=W^{i_{N_{D}-1}} .
\end{aligned}
$$

By a preceding observation, we know that conditionally on $\left(m_{D}, N_{D}\right)$, the processes $W^{i_{1}}, \ldots, W^{i_{N_{D}-1}}$ are independent and distributed according to $\mathbb{N}_{\mathrm{w}\left(m_{D}\right)}^{D}$.

Combining this with the strong Markov property at time $T_{D}$, we see that, conditionally on ( $N_{D}, W_{0}^{D}$ ), the processes $W^{D, 2}, \ldots, W^{D, N_{D}}$ are independent and distributed according to $\mathbb{N}_{\hat{W}_{0}^{D}}^{D}$ (recall that $\left.\widehat{W}_{0}^{D}=W_{T_{D}}\left(m_{D}\right)\right)$. An argument similar to the end of the proof of Theorem 3.2.1 (relying on independence properties of Poisson measures) also shows that, conditionally on $\left(N_{D}, W_{0_{\nu}}^{D}\right)$, the vector $\left(W^{D, 2}, \ldots, W^{D, N_{D}}\right)$ is independent of $W^{D, 1}$. Furthermore, denote by $\tilde{W}^{D, \ell}$ the timereversed processes

$$
\check{W}_{s}^{D, \ell}=W_{\left(\sigma\left(W^{D, \ell}\right)-s\right)^{+}}^{D, \ell} .
$$

The time-reversal property already used in the proof of Theorem 4.6.2 implies that the vectors $\left(W^{D, 1}, \ldots, W^{D, N_{D}}\right)$ and ( $\check{W}^{D, N_{D}}, \ldots, \check{W}^{D, 1}$ ) have the same conditional distribution given $\left(N_{D}, W_{0}^{D}\right)$. Hence, the conditional distribution of $\breve{W}^{D, 1}$, or equivalently that of $W^{D, 1}$, is also equal to $\mathbb{N}_{\hat{W}_{0}^{D}}$. This completes the proof of Theorem 4.7.2.

Remarks. (i) By considering the special case where the spatial motion $\xi_{t}$ is deterministic, $\xi_{t}=t$, and $E=\mathbb{R}_{+}, x=0$ and $D=[0, T)$ for some fixed $T>0$, we obtain an alternative proof of formulas derived in Theorem 2.7.1. In particular, formula (2.44) is a special case of (4.40). Similarly, (2.43) can be seen as a special case of (4.42).
(ii) In the stable case $\psi(u)=c u^{\alpha}$, the variable $N_{D}$ is independent of $W_{0}^{D}$, and its law is given by

$$
\mathbb{N}_{x}^{D}\left[r^{N_{D}}\right]=\frac{(1-r)^{\alpha}-1+\alpha r}{\alpha-1}
$$

Of course when $\alpha=2$, we have $N_{D}=2$.

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## Notation Index

$X$ Lévy process (Sect. 1.1.1)
$P$ probability measure under which $X$ starts from 0
$\psi$ Laplace functional of $X$ (Sect. 1.1.1)

$$
\psi(u)=\alpha u+\beta u^{2}+\int_{(0, \infty)} \pi(d r)\left(e^{-r u}-1+r u\right)
$$

$$
\tilde{\psi}(\lambda)=\frac{\psi(\lambda)}{\lambda}
$$

$$
\varphi(\lambda)=\frac{\lambda}{\psi^{-1}(\lambda)}
$$

$$
S_{t}=\sup _{s \leq t} X_{s}
$$

$I_{t}=\inf _{s \leq t} X_{s}$
$I_{t}^{s}=\inf _{s \leq r \leq t} X_{r}$
$T_{x}=\inf \left\{t \geq 0: X_{t}=-x\right\}$
$N$ excursion measure of $X-I$ (Sect. 1.1.2)
$N^{*}$ excursion measure of $X-S$ (Sect. 1.1.2)
$\sigma$ duration of the excursion under $N$ or under $N^{*}$
$L_{t}$ local time of $X-S$ at level 0
$L^{-1}(t)=\inf \left\{s \geq 0: L_{s}>t\right\}$
$X_{s}^{(T)}=X_{T+s}-X_{T}$
$\widehat{X}_{s}^{(t)}$ process time-reversed at $t: \widehat{X}_{s}^{(t)}=X_{t}-X_{(t-s)-}$
$H_{t}$ height process (Sect. 1.2)
$H(\mu)=\sup (\operatorname{supp}(\mu))$
$\rho_{t}$ exploration process (Sect. 1.2)
$M_{f}\left(\mathbb{R}_{+}\right)$set of all finite measures on $\mathbb{R}_{+}$
$k_{a} \mu$ "killing operator" on measures (Sect. 1.2, Sect. 3.1)
[ $\mu, \nu$ ] concatenation of the measures $\mu$ and $\nu$ (Sect. 1.2)
$\rho^{\mu}$ exploration process started at $\mu \in M_{f}\left(\mathbb{R}_{+}\right)$
$\mathbf{P}^{\mu}$ law of $\rho^{\mu}$
$\mathbf{M}$ invariant measure of $\rho$ (Sect. 1.2)
$L_{t}^{a}$ local times of the height process

$$
\begin{aligned}
u_{t}(\lambda) & =N\left(1-\exp \left(-\lambda L_{\sigma}^{t}\right)\right) \\
v(t) & =N\left(\sup _{s \geq 0} H_{s}>t\right) \\
\gamma_{\psi}(a, b) & =\frac{\psi(a)-\psi(b)}{a-b}
\end{aligned}
$$

$\eta_{t}$ dual of the exploration process (Sect. 3.1)
$\mathbb{M}$ invariant measure of $(\rho, \eta)$ (Sect. 3.1)
T set of all finite rooted ordered trees (Sect. 0.1)
$k_{v}(\mathcal{T})$ number of children of the vertex $v$ in the tree $\mathcal{T}$
$h_{v}$ mark (lifetime) of $v$
$\theta\left(e, t_{1}, \ldots, t_{p}\right)$ tree associated with the function $e$ and the times $t_{1}, \ldots, t_{p}$ (Sect. 3.2)
$\xi$ Markov process in $E$ (Sect. 4.1)
$\Pi_{x}$ law of $\xi$ started at $x$
$\mathcal{W}$ set of all $E$-valued killed paths
$\zeta=\zeta_{\mathrm{w}}$ lifetime of the killed path w
$\widehat{\mathrm{w}}=\mathrm{w}\left(\zeta_{\mathrm{w}}-\right)$
$m_{h}\left(s, s^{\prime}\right)=\inf _{s \leq r \leq s^{\prime}} h(r)$
$\bar{W}_{s}=\left(\rho_{s}, W_{s}\right)$ Lévy snake (Sect. 4.1)
$\mathbb{P}_{\mu, \mathrm{w}}$ law of the Lévy snake started at $(\mu, \mathrm{w})$
$\mathbf{Q}_{\mathbf{s}}$ transition kernels of the Lévy snake
$\mathbb{N}_{x}$ excursion measure for the Lévy snake
$\mathbb{P}_{\mu, \mathrm{w}}^{*}$ law of the Lévy snake started at $(\mu, \mathrm{w})$ and stopped at $\inf \left\{s \geq 0: \rho_{s}=0\right\}$
$\left\langle\mathcal{Z}_{a}, f\right\rangle=\int_{0}^{\sigma} d L_{s}^{a} f\left(\widehat{W}_{s}\right)$
$\tau(\mathrm{w})=\inf \{t \geq 0: \mathrm{w}(t) \notin D\}$
$L_{s}^{D}$ exit local time from $D$ (Sect. 4.3)
$\mathcal{Z}^{D}$ exit measure from $D$ (Sect. 4.3)
$\mathcal{W}^{*}$ set of all $E$-valued stopped paths
$T_{D}=\inf \left\{s \geq 0: \tau\left(W_{s}^{*}\right)<\infty\right\}$

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