Control and stabilization from geodesic domains

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Control of waves

Consider the wave equation on a Riemanian manifold M_g , $a \in L^{\infty}(M), a \ge 0, \ T > 0$

$$(\partial_t^2 - \Delta)u = f \times 1_{(0,T)} \times a(x), \quad (u \mid_{t=0}, \partial_t u \mid_{t=0}) = (u_0, u_1)$$

Given $(u_0, u_1) \in \mathcal{H}^1 = H^1(M) \times L^2(M)$ initial data and $(v_0, v_1) \in \mathcal{H}^1$ target data in energy space, can we choose f in suitable space such that

$$(u \mid_{t=T}, \partial_t u \mid_{t=T}) = (v_0, v_1)?$$

Natural space for f is $L^2((0, T) \times M)$. If answer yes: exact controlability

Stabilization for waves

$$(\partial_t^2 - \Delta + a(x)\partial_t)u = 0,$$

$$(u \mid_{t=0}, \partial_t u \mid_{t=0}) = (u_0, u_1) \in H^1 \times L^2 = \mathcal{H}^2$$

The natural energy is decaying $(a \ge 0)$

$$E(u)(t) = \int_{M} |\nabla_{x}u|^{2} + |\partial_{t}u|^{2} dx, \frac{d}{dt}E(t) = \int_{M} -a(x)|\partial_{t}u|^{2} dx$$

Question: speed of decay of E(u)(t)?

• The energy of all solutions tend to 0 iff there exists no non trivial stationary equilibrium, i.e.

$$-\Delta e = \lambda^2 e, a \times e = 0 \Rightarrow e = 0.$$

• Semi-group property: If there exists a uniform rate f(t),

$$\forall (u_0, u_1) \in \mathcal{H}^1, E(u)(t) \leq f(t)E(u)(0), \lim_{t \to +\infty} f(t) = 0,$$

then can choose $f(t) = Ce^{-ct}$ (uniform) stabilization.

Observation and HUM duality imply equivalence

• There exists a rate f(t) such that $\lim_{t\to+\infty} f(t) = 0$ and $\forall (u_0, u_1) \in H^1(M) \times L^2(M), E(u)(t) \leq f(t)E(u)(0).$

(and then can choose $f(t) = Ce^{-ct}$)

• $\exists T > 0, c > 0; \forall (u_0, u_1) \in H^1(M) \times L^2(M)$, if u is the solution to the damped wave equation, then

$$E(u)(0) \leq C \int_0^T \int_M 2a(x) |\partial_t u|^2 dx dt.$$

• $\exists T > 0, c > 0; \forall (u_0, u_1) \in H^1(M) \times L^2(M)$, if u is the solution to the undamped wave equation then

$$E(u)(0) \leq C \int_0^T \int_M 2a(x) |\partial_t u|^2 dx dt.$$

 There exists T > 0 such that The wave equation is exactly controlable in time T (and we can take the time given by observation) The geometric control assumption for waves

 $(a \in C^0(M), T)$ controls geometrically (M, g) if every geodesic starting from any point $x_0 \in M$ in any direction ξ_0 , $\gamma_{(x_0,\xi_0)}(s)$, encounters $\{a > 0\}$ in time smaller that T

Theorem (Rauch-Taylor, Bardos-Lebeau-Rauch 88', N.B- P.G.) $a \in C^0(M)$ geometric control is equivalent to observability (and hence control and stabilization) for wave equations. $a \in L^{\infty}(M)$ Strong Geometric Control is sufficient for observability which implies Weak Geometric Control.

$$\exists T, c > 0; \forall \rho_0 \in S^*M, \exists s \in (0, T), \exists \delta > 0; \\ a \ge c \text{ a.e. on } B(\gamma_{\rho_0}(s), \delta).$$
 (SGCC)

 $\exists T > 0; \forall \rho_0 \in S^*M, \exists s \in (0, T); \gamma_{\rho_0}(s) \in \text{supp}(a) \quad (\mathsf{WGCC})$

supp(a) is the support (in the distributional sense) of a,

The geometric control assumption



Some examples on tori



Figure: Checkerboards: the damping *a* is equal to 1 in the blue region, 0 elsewhere. The geodesics are (periodized) straight lines. The first example satisfies (SGCC) while all others satisfy (WGCC) but not (SGCC)

Stabilization for wave equations: the result Theorem (Does Stabilization holds? a = 1 in blue region 0 otherwise)









NO

(NB-PG 16)

Another geometric condition

When the manifold is a two dimensional torus and the damping *a* is a linear combination of characteristic functions of rectangles, i.e. there exists *N* rectangles (or polygons), $R_j, j = 1, ..., N$ (disjoint and non necessarily vertical), and $0 < a_j, j = 1, ..., N$ such that

$$a(x) = \sum_{j=1}^{N} a_j 1_{x \in R_j},$$
 (1)

(piecewise smooth domains, no infinite contact with geodesics= much easier)

Theorem (NB–P. Gérard 15-17)

Stabilization holds for the waves on \mathbb{T}^2 iff there exists T > 0 such that all geodesics (straight lines) of length T either encounters the interior of one of the rectangles or follows for some time one of the sides of a rectangle R_{j_1} on the left and for some time one of the sides of another (possibly the same) rectangle R_{j_2} on the right.

Stabilization for wave equations: the result



A geometric control condition for control When the manifold is a two dimensional surface and a is a linear combination of characteristic functions of geodesic polygons, i.e. there exists N polygons, $R_j, j = 1, ..., N$ (disjoint and non necessarily vertical), and $0 < a_j, j = 1, ..., N$ such that

$$a(x) = \sum_{j=1}^{N} a_j 1_{x \in R_j},$$
 (2)

Theorem (NB 17)

Let T > 0. Then exact controlability holds for the waves on M^2 if and only if a generalized geometric condition (defined in terms of an ODE on a sphere bundle over S^*M) is satisfied. Roughly speaking it says that all geodesics of length T either encounters the interior of one of the polygons or follows for some time one of the sides of a polygon R_{j_1} and there exists s > 0 such all neighbour geodesics spend an amount of time of order at least s in the interior of one of the polygons

The result on the sphere and the torus

Theorem (NB 17)

Let T > 0. Then exact controlability holds for the waves on M^2 if and only if There exists $\alpha > 0$ such that for almost every $(x_0, \xi_0) \in S^*M$,

$$\int_0^T a(x(s,x_0,\xi_0))ds \ge \alpha > 0.$$

here $(x(s, x_0, \xi_0), \xi(s, x_0, \xi_0)$ is the bicharacteristic starting from (x_0, ξ_0) at s = 0).

Control for wave equations: the result on the sphere

On spheres, geodesics are great circles and the generalized geometric condition reduces to checking

-The geodesic enters the *interior* of the control region $\omega = \{a(x) > 0\}$ or

- The geodesic follows (some) sides of (some) polygons and we check the following algorithm. Write the whole oriented geodesic circle

 $\Gamma = \Gamma^u \cup \Gamma^d \cup \Gamma^0,$

(parts of Γ which encounter the side of a polygon on the upper (lower) hemisphere–or not). Then any choice of oriented diameter D separates any piece of geodesic $\gamma(0, T)$ into three (possibly empty) pieces

$$\gamma' \cup \gamma^r \cup \gamma^c,$$

corresponding to the part on the left (right) of the diameter or on the diameter. Then we assume that we never have

$$(\gamma^{\prime} \subset \Gamma^{u} \text{ and } \gamma^{r} \subset \Gamma^{d}).$$

Contradiction argument

Want to prove observation estimates for half wave solutions with spectrally localized initial data (*h* small enough)

$$(\partial_t^2 - \Delta)u = 0, u_0 = 1_{a < -h^2 \Delta < b} u_0, u_1 = 1_{a < -h^2 \Delta < b} u_1 \qquad a < 1 < b$$
$$\|u_0\|_{L^2}^2 \le C \int_0^T \int_\omega |u|^2 (x, t) dx dt \ \omega = \{a > 0\}$$

Assume false then there exists sequences

$$a_n, b_n \rightarrow 1, u_n \in L^2, 1_{a_n < -h_n^2 \Delta < b_n} u_n = u_n,$$

such that

$$||u_n||_{L^2} = 1, \int_0^T \int_\omega |u_n|^2(x, t) dx dt = o(1)$$

First microlocalization Scales: $t, X \sim 1, \tau, \Xi \sim h^{-1}$. Consider operators $a(t, X, hD_t, hD_X), a \in C_0^{\infty}(T^*M)$.

If $a \ge 0$ then (Gårding) $a(t, X, hD_t, hD_X) \ge -Ch$.

Proposition

There exists a subsequence (now we drop all sub-indexes) and a positive measure μ (on continuous functions on T^*M) such that

$$\lim_{n\to+\infty} (a(t,X,h_nD_t,h_nD_X)u_n,u_n)_{L^2_{t,X}} = \langle \mu,a\rangle.$$

$$supp (\mu) \subset \{(t,\tau,X,\Xi); 1 = \tau^2 = \|\Xi\|_{g(x)}^2 = p(X,\Xi)$$
$$\iota(T^*(0,Y) \times M) = T, \qquad \partial_t \mu = H_p \mu, \qquad \mu \mid_{(0,T) \times \omega} = 0.$$

As a consequence, μ is supported on bicharacteristics which do not encounter ω but hence graze $\partial \omega$ on left or right

Second microlocalization

Understand at finer scales how the mass can concentrate on the geodesic from left or right. Work in a geodesic coordinate system (x, y) where

$$-\Delta = -\partial_x^2 - \partial_y^2 (1 + x^2 \kappa(y) + O(x^3)),$$

where the geodesic is given by $\{x = 0\}$ (and the bicharacteristic by $\{(x = 0, \xi = 0)\}$) and $\kappa(y)$ is the gauss curvature of the surface at point y.

Scales: 2 different regimes

• Transversal HF

$$t, y \sim 1, \tau = \eta = 1 + o(1), h^{1/2} \ll ||x, \xi|| = o(1)$$

Transversal LF

$$t, y \sim 1, \tau = \eta = 1 + o(1), \quad ||(x, \xi)|| \le Ch^{1/2}$$

Describe concentration at these scales and conclude contradiction.

2-pseudodifferential operators

– Symbols: functions $a(t, y, z, \tau, \eta, \zeta) \in S$ the class of smooth compactly supported in the (y, η) variables and polyhomogeneous of degree 0 near infinity in the (z, ζ) variables

$$|\partial^lpha_{t,y, au,\eta}\partial^\gamma_z\partial^\delta_\zeta a|\leq C(1+|z|+|\zeta|)^{-(\gamma+\delta)}.$$

– Operators : $\chi\in \mathit{C}^\infty_0(\mathbb{R}^2)$ equal to 1 near 0,

$$\begin{aligned} \mathsf{Op}_h(a) &= a(t, y, hD_t, hD_y, h^{-1/2}x, h^{1/2}D_x) \\ \mathsf{Op}_{h,\epsilon}(a) &= \mathsf{Op}(a \times \chi(\epsilon z, \epsilon \zeta)), \\ \mathsf{Op}_h^{\epsilon}(a) &= \mathsf{Op}(a \times (1 - \chi)(\epsilon z, \epsilon \zeta)), \end{aligned}$$

– Bad pseudodifferential calculus for Op(*a*) and Op_{*e*}(*a*) *L*² boundedness but no symbolic calculus , no Gårding – Good pseudodifferential calculus Op^{*e*}(*a*) (gain *e*²) symbolic calculus, and Gårding Approach inspired from works by Fermanian, Nier and Anantharaman-Macia for Schrödinger on tori. Here $S_{\frac{1}{2},\frac{1}{2}}$ calculus, Nier *S*_{1,1}, A–M, *S*_{0,0} calculus 2-microlocal measure: transversal HF (for tori $\epsilon
ightarrow h^{\epsilon}$)

$$t, y \sim 1, \tau = \eta = 1 + o(1), h^{1/2} \ll ||x, \xi|| = o(1),$$

If
$$a \ge 0$$
 then (Gårding) $Op^{\epsilon}(a) \ge -C\epsilon$.

Proposition

There exists a subsequence and a positive measure ν^+ (on continuous functions on $T^*N \times \mathbb{R}^2$ homogeneous of degree 0 at infinity in (z,ζ)) such that

$$\lim_{\epsilon\to 0}\lim_{n\to +\infty} \left(Op_{h_n}^{\epsilon}(a)u_n, u_n \right)_{L^2_{t,y,x}} = \langle \nu^+, a \rangle.$$

$$\mathsf{supp}\ (
u^+) \subset \{(t, au, y, z, \eta, \zeta); 1 = au = \eta\}$$

The projection of ν on the (t, y, τ, η) variables is bounded by the previous (1)-microlocal measure and additional propagation holds

$$(\partial_t - \partial_y - \zeta \partial_z + z \kappa(y) \partial_\zeta) \nu^+ = 0$$

Proof of propagation

Key remark

$$-h^{2}\Delta = -h^{2}\partial_{y}^{2}(1 - x^{2}\kappa(y) - h^{2}\partial_{x}^{2} + O(x^{3}))$$

= Op($\eta^{2}(1 + hz^{2}\kappa(y) + h\zeta^{2} + O(hz^{2}x))$ (3)

Compute

$$\frac{i}{2h} \Big[(h^2 \partial_t^2 - h^2 \Delta, \operatorname{Op}_{h_n}^{\epsilon}(a) \Big] = Op_{h_n}^{\epsilon} (-\tau \partial_t(a) + \eta \partial_y(a) + \zeta \partial_z(a) - \eta^2 \kappa(y) z \partial_{\zeta}(a) + O(\epsilon^2) + O(x) \Big]$$

implies

$$0 = \lim_{\epsilon \to 0} \lim_{n \to +\infty} \frac{i}{2h} \Big(\Big[(h^2 \partial_t^2 - h^2 \Delta, \operatorname{Op}_{h_n}^{\epsilon}(a) \Big] u_n, u_n \Big)_{L^2_{t,y,x}} \\ = \langle (\partial_t - \partial_y - \zeta \partial_z + \kappa(y) z \partial_\zeta) \nu, a \rangle$$

Conclusion in the transversal HF regime

- The measure ν^+ is invariant by the flow defined by previous equation
- It is supported on geodesics grazing $\partial \omega$
- By contradiction assumption

$$\int_{(0,T)}\int_{\omega}|u|^{2}(t,x)dxdt=o(1),$$

We deduce that if near points (t, y, x = 0) such that $(y, x = 0) \in \partial \omega^r$ then ν_+ is supported in $\{z \le 0\}$

- The geometric hypothesis implies that $u_+\equiv 0$

2-microlocal measure: transversal LF

We are looking at $(\chi_{\epsilon} = \chi(\epsilon \cdot))$

$$\lim_{\epsilon \to 0} \lim_{n \to +\infty} \left((a \times \chi_{\epsilon})(t, y, h_n D_t, h_n D_y, h_n^{-1/2} x, h_n^{1/2} D_x) u_n, u_n \right)_{L^2}$$

change variables in x, $z = h^{-1/2}x$, $v_n(z) = h^{1/4}u_n(h^{1/2}z)$ We are now looking at

$$\lim_{\epsilon \to 0} \lim_{n \to +\infty} \left((a \times \chi_{\epsilon})(t, y, h_n D_t, h_n D_y, z, D_z) v_n, v_n \right)_{L^2}$$

Due to the presence of the cut off χ_ϵ for any fixed ϵ , the operators

$$(a \times \chi_{\epsilon})(t, y, h_n D_t, h_n D_y, z, D_z)$$

are semi-classical operators in the (t, y) variable with values compact operators on $L^2_x = H$ The sequence v_n is bounded in $L^2_{loc}(\mathbb{R}^2_{t,y}; H)$.

2-microlocal measure: transversal LF

Proposition (P. Gérard 90')

There exists a subsequence and a positive measure ν^- on continuous functions on $T^*\mathbb{R}^2$ with values trace class operators

$$\lim_{n \to +\infty} \left((a \times \chi_{\epsilon})(t, y, h_n D_t, h_n D_y, z, D_z) v_n, v_n \right)_{L^2} = Tr \langle \nu^-, (a \chi_{\epsilon})(t, y, \tau, \eta, z, D_z) \rangle.$$

Radon-Nikodym: $\nu^{-} = A(t, y, \tau, \eta) d\rho$, where A is trace class

Proposition (Saut Scheurer?)

Consider the following classical one dimensional harmonic oscillator

$$(i\partial_s - \frac{\partial_z^2}{2} + \frac{\kappa(-s)z^2}{2})u = 0$$

Assume that u vanishes on $(\alpha, \beta)_s \times \mathbb{R}^+_z$. Then $u \equiv 0$.

Theorem

As soon as Γ^+ or Γ^- is non empty, the measure ν^- is identically 0.

The harmonic oscillator

$$t, y \sim 1, \tau = \eta = 1 + o(1), |z| + |D_z| \le \epsilon^{-1},$$

 $(ih_n\partial_t + \sqrt{-h_n^2\Delta})u_n = 0,$
 $-h^2\Delta = -h^2\partial_y^2(1 + h\kappa(y)z^2) - h\partial_z^2 + O_\epsilon(h^{3/2})$

$$\sqrt{-h^2\Delta} = -ih\partial_y \sqrt{\left(1 + h\kappa(y)z^2\right) - h\frac{\partial_z^2}{-h^2\partial_y^2}} + O_\epsilon(h^{3/2})$$
$$= -ih\partial_y + h\left(-\frac{\partial_z^2}{2} + \frac{\kappa(y)z^2}{2}\right) + o_\epsilon(h) \quad (4)$$

$$(ih_n(\partial_t - \partial_y) + h_n(-\frac{\partial_z^2}{2} + \frac{\kappa(y)z^2}{2})v_n = o_\epsilon(h).$$

$$s = \frac{t-y}{2}, r = \frac{t+y}{2} \Rightarrow (i\partial_s + H_{r-s})w_n = o_\epsilon(h)$$

$$\Leftrightarrow w_n = S_r(s,0)(w_n \mid_{s=0}) + o(1). \quad (5)$$

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An elementary argument in infinite dimension Conjugate (micro-locally) with the inverse of the evolution to replace the equation $(i\partial_s + H_{r-s})w_n = 0$ by $i\partial_s \tilde{w}_n = 0$.

$$\tilde{w}_n = S_r(0,s)w_n = S_r(s,0)^*w_n,$$

Let $\tilde{\nu}^-$ be the measure of the new sequence \tilde{w}_n . Then

$$\tilde{\nu}^- = S_r(0,s)\nu S_r(0,s)^*$$

is independent of the variable s. It writes

$$ilde{
u}^- = A(r) d\lambda_r \otimes ds \otimes \delta_{\sigma=0} \otimes \delta_{
ho=2},$$

Where A(r) is a family of hermician trace class operators and $d\lambda_r$ a non negative measure. Let $e_n(r) \in H$ a Hilbert Basis diagonalizing A(r) (eigenvalues λ_n). We get

$$\tilde{\nu}^{-} = \sum_{n} \lambda_{n} \langle \cdot, e_{n,r} \rangle_{H} e_{n,r} d\lambda_{r} \otimes ds \otimes \delta_{\sigma=0} \otimes \delta_{\rho=2},$$

$$\nu^{-} = \sum_{n} \lambda_{n} \langle \cdot, S(s,0) e_{n,r} \rangle_{H} S(s,0) e_{n,r} d\lambda_{r} \otimes ds \otimes \delta_{\sigma=0} \otimes \delta_{\rho=2},$$

Conclusion in the transversal LF regime

 $\nu^{-} = \sum \lambda_n \langle \cdot, S(s, 0)e_{n,r} \rangle_H S(s, 0)e_{n,r} d\lambda_r \otimes ds \otimes \delta_{\sigma=0} \otimes \delta_{\rho=2},$ If $\zeta(s, r)$ is supported in the region where the bicharacteristic grazes a part of Γ^r , then for any $\epsilon > 0$,

$$\lim_{n\to+\infty}\int_{s,r}\chi(s,r)\int_{z\in(0,\delta h^{-1/2})}|w_n|^2dzdrds=0,$$

We deduce that for any $\epsilon > 0$,

$$\langle \nu^-, \zeta(s,r) \mathbb{1}_{x>0} \chi(\epsilon(z,D_z)) \rangle = 0$$

$$\Rightarrow 0 = \int_{r,s} \zeta(r,s) \operatorname{Tr}(\sum_{n} \lambda_{n} \langle \cdot, S(s,0)e_{n,r} \rangle_{H} 1_{z>0} S(s,0)e_{n,r} d\lambda_{r} ds,$$
$$= \int_{r,s} \zeta(r,s) \sum_{n} \lambda_{n} \int_{z>0} |S(s,0)e_{n,r}|^{2} ds d\lambda_{r} \quad (6)$$

(we used that $S(s, o)e_n$ is a Hilbert basis of H). Hence implies $\forall n, \lambda_n = 0$ and $\nu^- = 0$.