# Persistence of periodic orbits in perturbed dissipative evolutionary equations <br> (Joint work with Jack Hale) 

## Geneviève Raugel

CNRS and Université Paris-Sud
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Dedicated to Jack Hale on the occasion of his 80th birthday

## Introduction

Let $X$ be a Banach space and $S_{0}(t): X \mapsto X$ be a dynamical system.
We assume that $S_{0}(t)$ admits a periodic orbit $p_{0}(t)$ of (least) period $T_{0}$ and that this periodic orbit is non-degenerate.
For $\varepsilon>0$ small, let $S_{\varepsilon}(t): X \mapsto X$ be a perturbed dynamical system, such that $S_{\varepsilon}(t)$ "converges" in some sense to $S_{0}(t)$ when $\varepsilon \rightarrow 0$.
Question Does $S_{\varepsilon}(t)$ admit a periodic orbit $p_{\varepsilon}(t)$ close to $p_{0}(t)$ of period $T_{\varepsilon}$ close to $T_{0}$ ? Is this periodic orbit unique?

## An abstract setting

$$
S_{0}(t): \quad u_{t}=A_{0} u(t)+G_{0}(u(t)), u(0)=u_{0} \in X
$$

and

$$
S_{\varepsilon}(t): \quad u_{t}=A_{\varepsilon} u(t)+G_{\varepsilon}(u(t)), u(0)=u_{\varepsilon} \in X
$$

where, for $\varepsilon \geq 0, A_{\varepsilon}$ is the generator of a $C^{0}$-semigroup and either $G_{\varepsilon}: X \rightarrow Y$, where $Y \hookrightarrow X$ or $G_{\varepsilon}: X \rightarrow Z$ where $X \hookrightarrow Z$. We assume that $A_{\varepsilon} \rightarrow A_{0}$ and $G_{\varepsilon} \rightarrow G_{0}$.
Definition: $p_{0}(t)$ is a non-degenerate periodic solution of (least) period $T_{0}$ if 1 is an isolated (algebraically) simple eigenvalue of the period map $U_{0}\left(T_{0}, 0\right)$ where $U_{0}(\sigma, s) w^{s}=w(\sigma)$ is the solution of

$$
\begin{equation*}
w_{t}=A_{0} w+D G_{0}\left(p_{0}(t)\right) w, w(s)=w^{s} \in X \tag{3}
\end{equation*}
$$

## Classical result

Classical theorem If (1) and (2) are finite-dimensional systems of ODE's, if $A_{\varepsilon}$ and $G_{\varepsilon}$ are continuous in $\varepsilon$ and if $p_{0}(t)$ is a non-degenerate periodic orbit of (1) of period $T_{0}$, then there exist $\varepsilon_{0}>0$ and $r>0$ such that, for $0<\varepsilon \leq \varepsilon_{0}$, (2) has a unique periodic solution $p_{\varepsilon}(t)$ of period $T_{\varepsilon}$ with $\left|T_{\varepsilon}-T_{0}\right| \leq r$ and $\left\|p_{0}(0)-p_{\varepsilon}(0)\right\| \leq r$. And $p_{\varepsilon}(t)$ is continuous in $\varepsilon$. Proof: Poincaré method or Lyapunov-Schmidt method .... The same theorem holds in the case of parabolic equations. Persistence of periodic orbits in equations of retarded type and neutral type (J. Hale and M. Weedermann, JDE, 2004)

## Problems arising in infinite dimensions

- In general, given $u_{0} \in X, t \in \mathbb{R} \mapsto S_{\varepsilon}(t) u_{0} \in X$ is not a Hölder-continuous map.
-The interesting perturbations are not regular (not continuous in $\varepsilon$ )
Goals:
-To generalize the previous theorem to dissipative systems (2 methods).
-To describe the local unstable (stable) manifolds of
$\Gamma=\left\{p_{0}(t) \mid t \in\left[0, T_{0}\right)\right\}$


## Part I: Examples

Example 1: A system of damped wave equations
$\Omega \subset \mathbf{R}^{n}$ is a bounded smooth domain.
$X=\left(H^{1}(\Omega) \times L^{2}(\Omega)\right)^{2}, f_{i}:(u, v) \in H^{1}(\Omega) \times H^{1}(\Omega) \mapsto H^{s}(\Omega)$,
$0<s<1, i=1,2$.
The dynamical system $S_{0}(t)$ is defined by

$$
\begin{align*}
u_{t t}+\beta u_{t}-\Delta u+\alpha u & =f_{1}(u, v) \\
v_{t t}+\beta v_{t}-\Delta v+\alpha u & =f_{2}(u, v) \\
\frac{\partial u}{\partial \nu}=\frac{\partial v}{\partial \nu} & =0 \text { on } \partial \Omega  \tag{4}\\
\left(u, u_{t}, v, v_{t}\right)(0, x) & =\left(u_{0}, u_{1}, v_{0}, v_{1}\right) \in X .
\end{align*}
$$

where $\alpha>0, \beta>0$. It can be written as a first order system

$$
w_{t}=A_{0} w(t)+G_{0}(w), w(0)=w_{0} \in X, \text { where }
$$

$G_{0}: w \in X \mapsto G_{0}(w)=\left(0, f_{1}, 0, f_{2}\right) \in Y=\left(H^{s+1}(\Omega) \times H^{s}(\Omega)\right)^{2}$
The embedding of $Y$ in $X$ is compact.

## Properties of Example 1

1. $S_{0}(t)$ is asymptotically smooth (a.s.): $\Rightarrow$ Every bounded invariant set $B$ is relatively compact
2. Regularity in space and time of bounded invariant sets If $w(t)$ is a bounded solution for $t \in(-\infty, T)$, then, for $t \leq T$,

$$
w(t)=\int_{-\infty}^{t} e^{A_{0}(t-s)} G_{0}(w(s)) d s \in\left(H^{s+1} \times H^{s}\right)^{2}
$$

Same regularity properties for linearized equations along smooth global bounded solutions and thus for the map $U\left(T_{0}, 0\right)$ if we have a periodic orbit.
Same regularity for the eigenfunctions of $U\left(T_{0}, 0\right)$ associated with eigenvalues $\lambda,|\lambda| \geq 1$.

## Examples of perturbations of $S_{0}(t)$

1. Regular perturbations: perturbations $f_{i}(\varepsilon, u, v)$ or $\gamma(\varepsilon)$
2. Non regular perturbations: perturbations of the domain $\Omega$ Thin domains: $Q_{\varepsilon}=\{(x, y) \mid x \in \Omega, 0<y<\varepsilon h(x)\}$ $h(x)=1$ : R. Johnson, M. Kamenski, P. Nistri, JDDE 1998 General $h, n=1$ : B. Abdelhedi, 2005

The thin domain perturbation is singular since we can only estimate $\left\|S_{\varepsilon}(t) w_{0}-S_{0}(t) w_{0}\right\|_{\tilde{X}}$ for $w_{0} \in \tilde{Y}$ where $\tilde{Y} \subsetneq \tilde{X}$.

Example 2: Hyperbolic perturbation of the Navier-Stokes equations
We consider the equations defined on $\Omega=\mathbf{T}^{2}$ :

$$
\begin{aligned}
\varepsilon u_{t t}^{\varepsilon}+u_{t}^{\varepsilon}-\nu \Delta u^{\varepsilon} & =-\nabla p^{\varepsilon}-u^{\varepsilon} \cdot \nabla u^{\varepsilon}+f \\
\operatorname{div} u^{\varepsilon} & =0 \\
\left(u^{\varepsilon}, u_{\tau}^{\varepsilon}\right)(0, y) & =\left(u_{0}(y), u_{1}(y)\right) .
\end{aligned}
$$

$S_{\varepsilon}(t)$ is a local dynamical system on $X=V^{s+1}(\Omega) \times V^{s}(\Omega)$, for $s \geq 0$, where $V^{s}=\left\{w \in H^{s}(\Omega)^{2} \mid\right.$ div $w=0, w$ is $2 \pi$-periodic $\}$. When $\varepsilon$ goes to 0 , we obtain the system $S_{0}(t)$ generated by the Navier-Stokes equations

$$
\begin{aligned}
v_{\tau}-\nu \Delta v & =-\nabla p-v \cdot \nabla v+f \\
\operatorname{div} v & =0 \\
v(0, y) & =u_{0} .
\end{aligned}
$$

Existence of periodic orbits: Chen and Price, Com. Math. Physics, 1999. Paicu, R.

## Other examples

Example 3: Second grade fluids Let $\Omega=\mathbf{T}^{2}$. For $\alpha>0, S_{\alpha}(t)$ is a global dynamical system on $V^{3}$,

$$
\begin{aligned}
& \partial_{t}(u-\alpha \Delta u)-\nu \Delta u+\operatorname{curl}(u-\alpha \Delta u) \times u+\nabla p=-\nabla p+f \\
& \operatorname{div} u=0 \\
& u(0, x)=u_{0}(x)
\end{aligned}
$$

where curl $u=\left(0,0, \partial_{1} u_{2}-\partial_{2} u_{1}\right)$ if $u=\left(u_{1}, u_{2}\right)$.

1. $S_{\alpha}(t)$ is not smoothing in finite time.
2. $S_{\alpha}(t)$ is a non regular perturbation of $S_{0}(t)$.
3. $S_{\alpha}(t)$ is a.s. and the bounded invariant sets are more regular.

Example 4: A system of weakly damped Schrödinger equations

## Summary

In these examples, $S_{\varepsilon}(t)$ has no Hölder-regularity in time and the convergence of $S_{\varepsilon}(t)$ to $S_{0}(t)$ is not good. BUT

1. $S_{\varepsilon}(t)$ is a.s. due to dissipation (second grad fluids or Schrödinger equations) or to compactness of the non-linearity (damped wave equation).
2. The bounded invariant sets of $S_{\varepsilon}(t)$ are more regular.
3. $S_{\varepsilon}(t)$ can be compared to $S_{0}(t)$ for smoother data,

$$
\left\|S_{\varepsilon}(t) u-S_{0}(t) u\right\|_{X} \leq C\left(t,\|u\|_{Y}\right) \varepsilon^{d}
$$

## Part II. A first method : An integral or functional method

$$
\begin{equation*}
S_{\varepsilon}(t): \quad u_{t}=A_{\varepsilon} u(t)+G(u), u(0)=u_{\varepsilon} \in X . \tag{5}
\end{equation*}
$$

We assume that

- (A.1) $G \in C^{2}(X, Y)$, where $Y \subsetneq X$ and $X, Y$ are Hilbert spaces.
- (A.2) $A_{\varepsilon}$ is the generator of a $C^{0}$-semigroup on $X$ and $Y$ and

$$
\left\|e^{A_{\varepsilon} t}\right\|_{L(X, X)}+\left\|e^{A_{\varepsilon} t}\right\|_{L(Y, Y)} \leq C e^{-\alpha t}, \alpha>0
$$

- (H.1) $S_{0}(t)$ admits a non-degenerate periodic solution $p_{0}(t)$ of period $T_{0}$ and, for $t \geq 0$,

$$
\left\|p_{0}(t)\right\|_{\tilde{Y}}+\left\|\dot{p}_{0}(t)\right\|_{\tilde{Y}} \leq R_{0},
$$

where $\tilde{Y} \subset D\left(\left(A_{0}\right)^{a}\right), a>0$.

- (H.2) If $U\left(T_{0}, 0\right)^{*} \varphi_{0}^{*}=\varphi_{0}^{*}$ and $\left\langle\dot{p}_{0}, \varphi_{0}^{*}\right\rangle=1$, then

$$
\left\|\varphi_{0}^{*}\right\|_{\tilde{Y}} \leq C_{0} .
$$

- (H.3) One has the estimate

$$
\left\|e^{A_{\varepsilon} t}-e^{A_{0} t}\right\|_{L(Y, X)} \leq C \varepsilon^{d}
$$

where $d>0$.

- (H.4) $t \mapsto e^{A_{0} t} D_{u} G\left(p_{0}\right)$ is locally Hölder-continuous from $\mathbf{R}^{+}$ into $L(X, X)$ and $t \mapsto e^{A_{0} t} G\left(p_{0}\right)$ is locally Lipschitz-continuous from $\mathbf{R}^{+}$into $X$.
- (H.5) $\left|\left\langle A_{\varepsilon} e^{A_{\varepsilon} t} G(w), \varphi_{0}^{*}\right\rangle\right| \leq C\|w\|_{X}$, for any $w \in X$.


## Theorem

Under the hypotheses (A.1), (A.2), (H.1) to (H.5), there exist $\varepsilon_{0}>0$ and $r>0$, such that, for $0<\varepsilon \leq \varepsilon_{0}$, there exists a unique periodic solution $p_{\varepsilon}(t)$ of Eq.(5) with period $T_{\varepsilon}$ such that $\left|T_{\varepsilon}-T_{0}\right| \leq r$ and $\left\|p_{0}(t)-p_{\varepsilon}(t)\right\| x \leq r$. Moreover, $p_{\varepsilon}$ converges to $p_{0}$ and $T_{\varepsilon}$ converges to $T_{0}$ when $\varepsilon \rightarrow 0$ (estimates).

Consequence: Existence (and uniqueness) of perturbed periodic orbits of period close to $T_{0}$ for systems of damped wave equations on thin domains.

## Proof

## Ingredients:

- Use of the variation of constants formula (Krasnoselskii, Zabreiko, Pustylnik and Sobolevskii, 1966; Gurova and Kamenskii, 1996, R. Johnson, M. Kamenski, P. Nistri, JDDE 1998, etc..)
- Lyapunov-Schmidt method
- Strict contraction fixed point theorem

Idea: If $p_{\varepsilon}(t)$ is a periodic orbit of $(5)$ of period $T_{\varepsilon}$, then,

$$
p_{\varepsilon}(0)=p_{\varepsilon}\left(T_{\varepsilon}\right)=e^{A_{\varepsilon} T_{\varepsilon}} p_{\varepsilon}(0)+\int_{0}^{T_{\varepsilon}} e^{A_{\varepsilon}\left(T_{\varepsilon}-s\right)} G\left(p_{\varepsilon}(s)\right) d s
$$

or
$p_{\varepsilon}(0)=\left(I-e^{A_{\varepsilon} T_{\varepsilon}}\right)^{-1} \int_{0}^{T_{\varepsilon}} e^{A_{\varepsilon}\left(T_{\varepsilon}-s\right)} G\left(p_{\varepsilon}(s)\right) d s$
$p_{\varepsilon}(t)=e^{A_{\varepsilon} t}\left(I-e^{A_{\varepsilon} T_{\varepsilon}}\right)^{-1} \int_{0}^{T_{\varepsilon}} e^{A_{\varepsilon}\left(T_{\varepsilon}-s\right)} G\left(p_{\varepsilon}(s)\right) d s+\int_{0}^{t} e^{A_{\varepsilon}(t-s)} G\left(p_{\varepsilon}(s)\right) d s$

Change of time scaling $t \rightarrow \frac{T}{T_{0}} t$ leads the equation

$$
\begin{equation*}
\tilde{u}_{t}=\frac{T}{T_{0}} A_{\varepsilon} \tilde{u}(t)+\frac{T}{T_{0}} G(\tilde{u}), \quad \tilde{u}(0)=\tilde{u}_{\varepsilon} \in X . \tag{6}
\end{equation*}
$$

We set:
$\Sigma_{\varepsilon, T}(t)=e^{\frac{T}{T_{0}} A_{\varepsilon} t}, C_{T_{0}}^{\text {per }}=\left\{w \in C^{0}(X, X) \mid w\right.$ is $T_{0}$-periodic $\}$.

$$
\begin{aligned}
\left(J_{\varepsilon}(T) w\right)(t)= & \Sigma_{\varepsilon, T}(t)\left(I-\Sigma_{\varepsilon, T}\left(T_{0}\right)\right)^{-1} \int_{0}^{T_{0}} \Sigma_{\varepsilon, T}\left(T_{0}-s\right) w(s) d s \\
& +\int_{0}^{t} \Sigma_{\varepsilon, T}(t-s) w(s) d s \\
F_{\varepsilon}(T, w)= & J_{\varepsilon}(T)\left(\frac{T}{T_{0}} G(w)\right)
\end{aligned}
$$

- $\varphi \in C_{T_{0}}^{\text {per }}$ is a fixed point of $F_{\varepsilon}(T,$.$) iff \varphi$ is a periodic solution of (6).
- $D_{u} F_{0}\left(T_{0}, p_{0}\right) \dot{p_{0}}=\dot{p_{0}}, 1$ is a (algebraically) simple eigenvalue, and $D_{u} F_{0}\left(T_{0}, p_{0}\right)-I \in L\left(C_{T_{0}}^{\text {per }}(X)\right)$ is Fredholm of index 0 .
- Good estimates of $F_{\varepsilon}-F_{0}$ and of $D_{u} F_{\varepsilon}-D_{u} F_{0}$.

Let $v_{0}^{*} \in\left(C_{T_{0}}^{p e r}(X)\right)^{*}$ such that $\left(D_{u} F_{0}\left(T_{0}, p_{0}\right)-I\right)^{*} v_{0}^{*}=0$ and $\left\langle\dot{p}_{0}, v_{0}^{*}\right\rangle=1$, then

$$
C_{T_{0}}^{p e r}(X)=\left\{\dot{p}_{0}\right\}+Z, \quad Z=\left\{w \in C_{T_{0}}^{p e r}(X) I\left\langle w, v_{0}^{*}\right\rangle=0\right\}
$$

Goal: For $r>0$ small, find $(T, \varphi) \in B_{\mathbf{R}}\left(T_{0}, r\right) \times B_{C_{T_{0}}^{\text {per }}(X)}(0, r)$,
such that

$$
F_{\varepsilon}\left(T, p_{0}+\varphi\right)-\left(p_{0}+\varphi\right)=0
$$

Use a Lyapunov-Schmidt method and solve
$L_{\varepsilon}(\varphi, T)=F_{\varepsilon}\left(T, p_{0}+\varphi\right)-\left(p_{0}+\varphi\right)-\left\langle F_{\varepsilon}\left(T, p_{0}+\varphi\right)-\left(p_{0}+\varphi\right), v_{0}^{*}\right\rangle=0$ $\left\langle F_{\varepsilon}\left(T, p_{0}+\varphi\right)-\left(p_{0}+\varphi\right), v_{0}^{*}\right\rangle=0$.

Step 1: One shows that

$$
\mathcal{L}_{\varepsilon}(\varphi, T)=\varphi-\left(\left(D_{u} F_{0}\left(T_{0}, p_{0}\right)-I\right)_{/ z}\right)^{-1} L_{\varepsilon}(\varphi, T)
$$

is a strict contraction in $B_{Z}(0, r)$, for $\varepsilon>0$ and $r$ small.
$\Rightarrow$ For $\varepsilon$ small and $T$ close to $T_{0}$, there exists a unique fixed point $\varphi(\varepsilon, T) \in B_{Z}(0, r)$ of $\mathcal{L}_{\varepsilon}(\varphi, T)$.

Step 2: One solves the equation

$$
\begin{equation*}
M_{\varepsilon}(T)=\left\langle F_{\varepsilon}\left(T, p_{0}+\varphi(\varepsilon, T)\right)-\left(p_{0}+\varphi(\varepsilon, T)\right), v_{0}^{*}\right\rangle=0 . \tag{7}
\end{equation*}
$$

Under the hypothesis (H.5), $M_{\varepsilon}$ is a strict contraction from $B_{\mathbf{R}}(0, r)$ into itself. Thus there exists a unique $T_{\varepsilon} \in B_{\mathbf{R}}(0, r)$ such that (7) holds. And $\varphi\left(\varepsilon, T_{\varepsilon}\right)$ is the $T_{0}$-periodic solution of (6).

## Remarks:

1. If $p_{0}(t)$ is of class $C^{2}$, we show, without the hypothesis (H.5) that (7) has a solution in $B_{\mathrm{R}}(0, r)$ by the Schauder fixed point theorem. If we assume that all periodic orbits of (5) are of class $C^{1+\delta}, \delta>0$, then one can prove the uniqueness.
2. In the case of thin product domains, R. Johnson, M. Kamenski and P. Nistri proved the existence of the perturbed periodic solution by using a topological degree argument.

## Part III. A modified Poincaré method

Appropriate method if the non-linearity is non-compact, but dissipative. Let $\tilde{Y} \hookrightarrow Y \hookrightarrow$ compact $X$ be Hilbert spaces. Let $S_{\varepsilon}(t), \varepsilon \geq 0$, be a dynamical system on $X, Y$ and $\tilde{Y}$. We assume (H.1), (H.2) and

- (H.3) $t \mapsto p_{0}(t)$ is of class $C^{2}(\mathbf{R}, X)$.
- (H.4) $r\left(\sigma_{\mathrm{ess}}\left(U\left(T_{0}, 0\right)\right)\right)<1$.
- (H.5) There exist $R_{1} \geq 2 R_{0}$ and $0<k_{1}<1$ s.t., for $T_{0} / 2 \leq t \leq 2 T_{0},\|u\|_{Y} \leq R_{1}$, for $\varepsilon \geq 0$,

$$
\left\|S_{\varepsilon}(t) u\right\|_{Y} \leq k_{1} R_{1} .
$$

- (H.6) There exists $\beta>0$ s.t., for $R>0,0<t_{0}<t \leq 2 T_{0}$, for $w \in B_{Y}(0, R)$,

$$
\begin{aligned}
\left\|S_{\varepsilon}(t) w-S_{0}(t) w\right\|_{X} & \leq \varepsilon^{\beta} K_{0}(R) \\
\left\|\left(D S_{\varepsilon}(t) w\right) u-\left(D S_{0}(t) w\right) u\right\|_{X} & \leq \varepsilon^{\beta} K_{0}(R)\|u\|_{Y}
\end{aligned}
$$

- (H.7) Besides $\lambda_{0}=1$, there exist $m$ distinct eigenvalues $\lambda_{i} \neq 1$ of $U\left(T_{0}, 0\right)$ with algebraic multiplicity $d_{i}$ such that $\left|\lambda_{i}\right| \geq 1$. The corresponding (generalized) eigenvectors are in $\tilde{Y}$.
There exists $0<k_{2}<1$ s.t., if $P_{i}$ is the spectral projection onto the eigenspace associated with $\lambda_{i}, i=0,1 \ldots, m$ and $Q=I-\sum_{i=0}^{m} P_{i}$, we have

$$
\left\|U\left(T_{0}, 0\right) \varphi\right\|_{X} \leq k_{2}\|\varphi\|_{X}, \quad \forall \varphi \in Q X
$$

- (H.8) There exists $K>0$ s.t., if $u_{\varepsilon}(t)$ is a bounded orbit of $S_{\varepsilon}(t)$, for $t \in \mathbf{R}$, then $u_{\varepsilon}(t)$ is bounded in $\tilde{Y}$ and

$$
\sup _{t}\left\|u_{\varepsilon}(t)\right\|_{\tilde{Y}} \leq K \sup _{t}\left\|u_{\varepsilon}(t)\right\|_{X}
$$

If $u_{\varepsilon}^{1}(t)$ and $u_{\varepsilon}^{2}(t)$ are bounded orbits of $S_{\varepsilon}(t)$, for $t \in \mathbf{R}$, then

$$
\sup _{t \in \mathbf{R}}\left\|u_{\varepsilon}^{1}(t)-u_{\varepsilon}^{2}(t)\right\|_{Y} \leq K \sup _{t \in \mathbf{R}}\left\|u_{\varepsilon}^{1}(t)-u_{\varepsilon}^{2}(t)\right\|_{X}
$$

- (H.9) For any $\varphi \in \tilde{Y}$, the map $t \mapsto S_{\varepsilon}(t) \varphi$ is locally Lipschitz-continuous from $\mathbf{R}$ into $X$.


## Second method: Results

Theorem
Under the hypotheses (H.1) to (H.9), there are $\varepsilon_{0}>0$ and $r>0$, such that, for $0<\varepsilon \leq \varepsilon_{0}$, there exists a unique periodic orbit $p_{\varepsilon}(t)$ of $S_{\varepsilon}(t)$ with period $T_{\varepsilon}$ such that $\left|T_{\varepsilon}-T_{0}\right| \leq r$ and $\left\|p_{0}(0)-p_{\varepsilon}(0)\right\|_{x} \leq r$. Moreover, $p_{\varepsilon}$ converges to $p_{0}$ and $T_{\varepsilon}$ converges to $T_{0}$ when $\varepsilon \rightarrow 0$ (estimates).

Consequence: Existence (and uniqueness) of perturbed periodic orbits of period close to $T_{0}$ in the examples 2 to 4 (hyperbolic Navier-Stokes; second grade fluids, weakly damped Schrödinger equations).
Ingredients of the proof:

1. Lyapunov-Schmidt method (modified Poincaré method)
2. Schauder fixed point theorem (topological degree also?)
3. Uniqueness

## Proof

We write $X=\left\{\dot{p}_{0}(0)\right\} \oplus Z_{1} \oplus Z_{2}$ where $Z_{1}=\oplus_{i=1}^{m} P_{i} X, Z_{2}=Q X$. We recall that $P_{0} w=\left\langle w, \varphi_{0}^{*}\right\rangle \dot{p}_{0}(0)$ and $\left\langle v, \varphi_{0}^{*}\right\rangle=0$ if $v \in Z_{1} \oplus Z_{2}$. Goal: For $r>0$ small, find $\left(T_{\varepsilon}, p_{\varepsilon}(0)\right) \in B_{\mathbf{R}}\left(T_{0}, r\right) \times B_{X}\left(p_{0}(0), r\right)$, such that

$$
S_{\varepsilon}\left(T_{\varepsilon}\right) p_{\varepsilon}(0)=p_{\varepsilon}(0)
$$

Step 1: We write: $p_{\varepsilon}(0)=p_{0}(0)+\varphi+\psi$, where $\psi=\sum_{j=1}^{d} \alpha_{j} \psi_{j} \in Z_{1}, \varphi \in Z_{2}$. For $r_{2}>0$ small, we set

$$
\mathcal{B}=\left\{\varphi \in B_{Z_{2}}\left(0, r_{2}\right) \mid \varphi+p_{0}(0) \in B_{Y}\left(0, R_{1}\right)\right\},
$$

and, for $\psi \in B_{Z_{1}}\left(0, r_{1}\right),\left|T-T_{0}\right|<\eta, 0<\varepsilon \leq \varepsilon_{0}$, where $\varepsilon_{0}, r_{1}, \eta$ are small, we define the map

$$
\mathcal{L}_{\varepsilon}(T, \varphi, \psi)=\left(I-P_{0}-\sum_{i=1}^{m} P_{i}\right)\left(S_{\varepsilon}(T)\left(p_{0}(0)+\varphi+\psi\right)-p_{0}(0)\right)
$$

$\mathcal{L}_{\varepsilon}(T, \varphi, \psi) \in \mathcal{B}$. By Schauder fixed point theorem, there exists a fixed point $\varphi_{\varepsilon}(T, \psi) \in \mathcal{B}$ of the $\operatorname{map} \mathcal{L}_{\varepsilon}(T, \cdot, \psi)$.

## Proof (continued)

Step 2: For $\left|T-T_{0}\right|<\eta, 0<\varepsilon \leq \varepsilon_{0}$, to find $\psi \in B_{Z_{1}}\left(0, r_{1}\right)$ s. t .

$$
\mathcal{M}_{\varepsilon}(T, \psi) \equiv \sum_{i=1}^{m} P_{i}\left(S_{\varepsilon}(T)\left(p_{0}(0)+\varphi_{\varepsilon}(T, \psi)+\psi\right)-p_{0}(0)\right)=\psi
$$

If $r_{1}>0, \eta>0$ and $\varepsilon_{0}>0$ are small, $\mathcal{M}_{\varepsilon}(T, \psi) \in B_{Z_{1}}\left(0, r_{1}\right)$.
By Schauder fixed point theorem, the map $\mathcal{M}_{\varepsilon}(T, \cdot)$ has a fixed point $\psi_{\varepsilon}(T) \in B_{Z_{1}}\left(0, r_{1}\right)$.
Step 3: For $0<\varepsilon \leq \varepsilon_{0}$, to find $T$ with $\left|T-T_{0}\right|<\eta$, s.t.

$$
\left.\left\langle S_{\varepsilon}\left(T_{\varepsilon}\right)\left(p_{0}(0)+\varphi_{\varepsilon}\left(T, \psi_{\varepsilon}(T)\right)+\psi_{\varepsilon}(T)\right)-p_{0}(0)\right), \varphi_{0}^{*}\right\rangle=0
$$

or to find a fixed point $\tau,|\tau|<\eta$ of the map

$$
\begin{gathered}
F_{\varepsilon}(\tau)=\tau \dot{p}_{0}(0)-\left\langleS _ { \varepsilon } ( T _ { 0 } + \tau ) \left( p_{0}(0)+\varphi_{\varepsilon}\left(T_{0}+\tau, \psi_{\varepsilon}\left(T_{0}+\tau\right)\right)\right.\right. \\
\left.\left.+\psi_{\varepsilon}\left(T_{0}+\tau\right)\right)-p_{0}(0), \varphi_{0}^{*}\right\rangle
\end{gathered}
$$

For $\varepsilon_{0}>0$ small enough, $F_{\varepsilon}$ has a fixed point $\tau_{\varepsilon}$, with $|\tau|<\eta$.

## Proof (end)

Set $T_{\varepsilon}=T_{0}+\tau_{\varepsilon}$. Thus, $S_{\varepsilon}(\cdot)$ has a periodic orbit

$$
p_{\varepsilon}(t)=S_{\varepsilon}(t)\left(p_{0}(0)+\varphi_{\varepsilon}\left(T_{\varepsilon}, \psi_{\varepsilon}\left(T_{\varepsilon}\right)\right)+\psi_{\varepsilon}\left(T_{\varepsilon}\right)\right)
$$

of period $T_{\varepsilon}$. Since $p_{\varepsilon}(t)$ is a periodic orbit of $S_{\varepsilon}(\cdot), p_{\varepsilon}(t)$ is uniformly bounded in $\tilde{Y}$.
Step 4: Since $p_{\varepsilon}(t)$ is uniformly bounded in $\tilde{Y}$, we can use the hypotheses (H.8) and (H.9) to show the uniqueness of the periodic solution with period $T_{\varepsilon}$ s.t. $\left|T_{\varepsilon}-T_{0}\right|<r$ and $\left\|p_{0}(0)-p_{\varepsilon}(0)\right\| x \leq r$ (Strict contraction argument).

## Part IV. Local unstable and stable manifolds

We take $\varepsilon=0$ in Equation (5),

$$
u_{t}=A_{0} u+G(u), \quad u(0)=u_{0} \in X
$$

and suppose:

- (A.1) and (A.2) hold,
- $p_{0}(t) \in C^{2}(\mathbf{R}, X)$,
- $r\left(\sigma_{\mathrm{ess}}\left(U\left(T_{0}, 0\right)\right)\right)<1$
- the periodic orbit $\Gamma_{0}=\left\{p_{0}(t) \mid t \in\left[0, T_{0}\right)\right\}$ is hyperbolic, i.e.

$$
\left(\sigma\left(U\left(T_{0}, 0\right)\right)-\{1\}\right) \cap S^{1}=\emptyset
$$

- $D_{u} G\left(p_{0}(t)\right) \in L(X, X)$ is a compact map, for $t \in\left[0, T_{0}\right)$.

Thus the index $i\left(\Gamma_{0}\right) \equiv i\left(U\left(T_{0}, 0\right)\right)$, which is the number of eigenvalues $\lambda$ of $U\left(T_{0}, 0\right)$ with $|\lambda|>1$, is finite.

Let $V$ be a neighbourhood of $\Gamma_{0}$, we define the local stable and unstable sets by

$$
\begin{aligned}
W_{\text {loc }}^{s}\left(\Gamma_{0}\right) \equiv W^{s}\left(\Gamma_{0}, V\right)=\left\{u \in X \mid S_{0}(t) u\right. & \in V, \forall t \geq 0 ; \\
& \left.\lim _{t \rightarrow+\infty} \delta X\left(S_{0}(t) u, \Gamma_{0}\right)=0\right\}, \\
W_{\text {loc }}^{u}\left(\Gamma_{0}\right) \equiv W^{u}\left(\Gamma_{0}, V\right)=\left\{u \in X \mid S_{0}(t) u\right. & \in V, \forall t \leq 0 ; \\
& \left.\lim _{t \rightarrow-\infty} \delta_{X}\left(S_{0}(t) u, \Gamma_{0}\right)=0\right\}
\end{aligned}
$$

Theorem (J. Hale, R)
$W_{\text {loc }}^{s}\left(\Gamma_{0}\right)$ (resp. $W_{\text {loc }}^{u}\left(\Gamma_{0}\right)$ ) is a $C^{1}$-submanifold of codimension $i\left(\Gamma_{0}\right)$ (respectively of dimension $\left.i\left(\Gamma_{0}\right)+1\right)$. Moreover, there exist positive constants $\alpha$ and $\beta$ such that, for any $u_{0}$ in $W_{\text {loc }}^{s}\left(\Gamma_{0}\right)$ (resp. in $\left.W_{\text {loc }}^{u}\left(\Gamma_{0}\right)\right)$, there is a number $\tau_{u_{0}}\left(\right.$ resp. $\left.\theta_{u_{0}}\right)$ s.t.

$$
\begin{aligned}
\left\|S_{0}(t) u_{0}-p_{0}\left(t+\tau_{u_{0}}\right)\right\|_{x} \leq C e^{-\alpha t}, \quad t \geq 0 \\
\left(\text { resp. }\left\|S_{0}(t) u_{0}-p_{0}\left(t+\theta_{u_{0}}\right)\right\|_{x} \leq C e^{\beta t}, \quad t \leq 0\right)
\end{aligned}
$$

Corollary Under the above hypotheses and the hypotheses of Part II, if $\Gamma_{0}$ is the periodic orbit of period $T_{0}$ of $S_{0}(t)$ and $\Gamma_{\varepsilon}$ the periodic orbit of period $T_{\varepsilon}$ of $S_{\varepsilon}(t)$, then

$$
\operatorname{dist}_{X}\left(W_{l o c}^{u}\left(\Gamma_{0}\right), W_{l o c}^{u}\left(\Gamma_{\varepsilon}\right)\right) \leq C \varepsilon^{d}
$$

where $d>0$.
(Particular case of thin domains has been proved by B. Abdelhedi (2005))

Ingredients of the proof:

- If $V$ is a neighbourhood of $\Gamma_{0}$, we define the local synchronized stable and local synchronized unstable sets of $p_{0}\left(\theta_{0}\right) \in \Gamma_{0}$ as

$$
\begin{aligned}
& \mathcal{W}_{\theta_{0}, l o c}^{s}\left(\Gamma_{0}\right) \equiv \mathcal{W}_{\theta_{0}}^{s}\left(\Gamma_{0}, V\right)=\{u=X \mid S_{0}(t) u \in V, \forall t \geq 0 ; \\
&\left.\lim _{t \rightarrow+\infty} S_{0}(t) u-p_{0}\left(t+\theta_{0}\right)=0\right\} \\
& \mathcal{W}_{\theta_{0}, l o c}^{u}\left(\Gamma_{0}\right) \equiv \mathcal{W}_{\theta_{0}}^{u}\left(\Gamma_{0}, V\right)=\left\{u \in X \mid S_{0}(t) u \in V, \forall t \leq 0 ;\right. \\
&\left.\lim _{t \rightarrow-\infty} S_{0}(t) u-p_{0}\left(t+\theta_{0}\right)=0\right\}
\end{aligned}
$$

We define the local synchronized stable and local synchronized unstable sets of $\Gamma_{0}$ as

$$
\begin{aligned}
& \mathcal{W}_{\text {loc }}^{s}\left(\Gamma_{0}\right)=\cup_{\theta_{0} \in\left[0, T_{0}\right)} \mathcal{W}_{\theta_{0}}^{s}\left(\Gamma_{0}, V\right) \\
& \mathcal{W}_{\text {loc }}^{u}\left(\Gamma_{0}\right)=\cup_{\theta_{0} \in\left[0, T_{0}\right)} \mathcal{W}_{\theta_{0}}^{u}\left(\Gamma_{0}, V\right)
\end{aligned}
$$

and we show that

$$
\mathcal{W}_{l o c}^{s}\left(\Gamma_{0}\right)=W_{l o c}^{s}\left(\Gamma_{0}\right), \quad \mathcal{W}_{l o c}^{u}\left(\Gamma_{0}\right)=W_{l o c}^{u}\left(\Gamma_{0}\right)
$$

- We introduce a new coordinates system around $\Gamma_{0}$. Every $u$ in a small tubular neighbourhood of $\Gamma_{0}$ is written as

$$
u=p_{0}(\theta)+Q(\theta) w,\|w\|_{x} \leq \delta
$$

where $\delta>0$ is small enough, $\theta \in R, w \in W, W$ is a linear subspace of $X$ of codimension $1, Q(\theta)$ is a continuous $T_{0}$-periodic map of $W$ into $X$ and $Q(\theta) W$ is transversal to $\dot{p}_{0}(\theta)$ for all $\theta$.

We obtain the new system

$$
\begin{aligned}
\dot{\theta} & =1+g(w) \\
\dot{w} & =\tilde{A} w+h(\theta, w)
\end{aligned}
$$



